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# Multi-functional L-histidine self-assembled monolayers on SnO<sub>2</sub> electron transport layer to boost photovoltaic performance of perovskite solar cells

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# ABSTRACT

The photovoltaic performance of perovskite solar cells (PSCs) is sensitive to the interface properties between perovskite and electron transport layer (ETL). An appropriate interface engineering on  $SnO_2$  ETL is highly desired for achieving efficient PSCs. Herein, we introduced L-histidine (L-His) self-assembly layer on  $SnO_2$  ETL to realize three following functions. The crystal quality and conductivity of  $SnO_2$  ETL have been improved since the carboxyl groups in L-His coordinate with the tin ion in  $SnO_2$  and fill the oxygen vacancy. The improved wettability caused by L-His deposition induces the formation of perovskite films with larger grain size and better crystallization quality. The hydrogen bond formed between perovskite and L-His proved by proton nuclear magnetic resonance (<sup>1</sup>H NMR) enhances the interface interaction, and the L-His also modulates the energy level structure of the ETL/perovskite interface, which not only enhances the charge transport properties but also leads to a higher open-circuit voltage. Based on these benign optimizations, the CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>-based PSCs with the optimized  $SnO_2/L$ -His ETL have achieved a power conversion efficiency of 21.04%. Such biological friendly modification technique is highly compatibility to the technique of large-scale production of PSCs.

#### 1. Introduction

Due to favorable optoelectronic properties including high absorption coefficients, long carrier diffusion lengths and high defect tolerance, organic-inorganic halide perovskites (OIHPs) have drawn remarkable attention in the field of perovskite light-emitting diodes and solar cell devices [1-4]. In recent years, the power conversion efficiency (PCE) of planar-type perovskite solar cells (PSCs) has been improved to 25.8% by optimizing the absorber materials, charge transporting layers and their interface [5]. The interface between electron transport layer (ETL) and perovskite plays a vital role in determining the overall photovoltaic performance of PSCs, since it positively influences the crystal quality of perovskite layer and the electron extraction/transport properties [6]. The SnO<sub>2</sub> has become a common ETL material for planar heterojunction cells due to its excellent advantages such as high transmittance, high carrier mobility and lower temperature preparation process [7]. However, pinholes and oxygen vacancy defects within SnO2 films and especially the energy level mismatch between SnO<sub>2</sub> and perovskite still weaken the charge transport efficiency and interfacial stability. Due to the mismatch of energy levels and defects, the charge extraction efficiency at the SnO<sub>2</sub>/perovskite interface is low and the carrier recombination is severe, resulting in low device efficiency and poor stability. The development of the additives or interlayers can not only improve the performance of the ETL, but also has a certain potential to affect the perovskite crystallization and reduce the recombination at the interface, which will help to improve the overall performance of the device. Therefore, optimizing the interface of the electron transport layer is an effective way to improve the photovoltaic performance of the device.

So far, various approaches have been used to modify the interface between ETL and perovskite, such as introducing multifunctional additives or inorganic/organic overlayers [8–15]. Researchers added some multifunctional additives into the SnO<sub>2</sub> precursor solution to realize simultaneous passivation of SnO<sub>2</sub>, perovskite and interface defects through synergistic effect of anion and cation toward efficient and stable planar PSCs [12–14]. For instance, Zang group added LiOH into the SnO<sub>2</sub> layer and Li<sup>+</sup> doped into both perovskite and SnO<sub>2</sub> layers to

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Received 19 February 2022; Received in revised form 18 July 2022; Accepted 28 July 2022 Available online 28 July 2022 0013-4686/© 2022 Elsevier Ltd. All rights reserved. improve the carrier transport and extraction [13]. Similarly, our group also used potassium sodium tartrate with chelating function and  $K^+/Na^+$ to synchronously optimize the SnO<sub>2</sub> layer, perovskite and interface with the help of  $K^+/Na^+$  diffusion [15]. As for the inorganic/organic overlayers, Zhu et al. covered the SnO<sub>2</sub> ETL with KCl to improve the energy level arrangement and reduce the number of trap states, which effectively improved the charge extraction/transport and resulted in a higher PCE [10]. Hu et al. inserted the MACl layer between the SnO<sub>2</sub> ETL and perovskite layer to successfully accelerate the carrier extraction and transport and suppress the interfacial recombination, so that a very high open circuit voltage of 1.19 V has been finally achieved [12]. Our group also used CdS quantum dot layer to cover SnO<sub>2</sub> ETL to accelerate charge extraction and transfer for achieving highly efficient PSCs [11]. However, these methods suffered from more complicated process, higher preparation cost and certain toxicity. Therefore, more efforts should be taken to explore simpler technique and bio-friendly materials to optimize the interface between perovskite and SnO<sub>2</sub> ETLs.

Self-assembled monolayers (SAMs) have been proved to be a kind of powerful technique to manipulate chemical and physical interfacial properties for improving the performance of photoelectric devices, which also owns advantages of low cost and compatibility with largescale preparation processes to satisfy the requirement of commercial application. So far, SAMs have been utilized to modify the interfaces of PSCs and the following positive effects have been exhibited [16,17]: (1) Appropriate SAMs covered on ETL can induce the growth of high-quality perovskite films with optimized morphology, better crystallinity, larger grain size and less defects. (2) Interfacial dipole moment can be well turned via modulating polarity of molecules, which helps to change the work function of ETL and build desired energy barrier at the interface to accelerate carrier extraction from the perovskite layer. (3) SAMs can passivate charge trap states on the interface and suppress the charge accumulation at the interface to reduce or eliminate device hysteresis.

In recent years, amino acids such as alanine (C<sub>3</sub>H<sub>7</sub>NO<sub>2</sub>), glycine (C2H5NO2) and L-cysteine (C3H7NO2S) with great advantages of lowcost and good biocompatibility have been used as SAMs to modify ETL/perovskite interface layers to improve the photovoltaic performance of PSCs [18-20]. The bifunctional groups of -COOH and -NH<sub>2</sub> within amino acids have been proved to have ability to simultaneously cross-link the ETL and perovskite layers, leading to the faster interfacial charge extraction and transport. However, there are many kinds of amino acids, so that it is particularly important to choose the simpler and more effective one to well realize the above functions. Among them, L-Histidine (C<sub>6</sub>H<sub>9</sub>N<sub>3</sub>O<sub>2</sub>, L-His) was firstly considered due to the following reasons. From point view of the molecular structure, the biggest difference between L-His and alanine, glycine and L-cysteine mentioned above is that L-His has an imidazole ring. Considering the chemical structure of L-His, the oxygen atoms on the carboxyl group could coordinate with Sn<sup>4+</sup> in SnO<sub>2</sub> nanoparticles to form a complex, which can promote the film quality of ETL to suppress the non-radiative recombination and improve the electron transfer and extraction efficiency at the ETL and interface [21]. The amino group can form hydrogen bonds with the organic cations of perovskite at the interface, which not only improves the compact interactions of the interface, but also improves the interfacial energy level matching, increases the charge extraction efficiency of the SnO2/perovskite interface, and reduces the carrier complex [22]. Imidazole ring is an electron deficient aromatic ring with non centrosymmetric structure, which can generate intramolecular charge transfer to improve the chelation ability of molecules with Sn ions or Pb related defects at the ETL/Perovskite interface [23]. Thus, the L-His is able to exert stronger chelation and defect passivation ability, which is more beneficial for the improvement of photovoltaic performance of PSCs.

Herein, L-His was selected as a SAM buffer layer to modify  $SnO_2$  ETL to enhance photovoltaic performance of PSCs. The crystal quality and photoelectric properties of both ETL and perovskite have been simultaneously improved. The interface charge transfer process has been

greatly improved. After optimizing the concentration of L-His SAMs, the PCE of  $MAPbI_3$  based PSCs can reach up to 21.04% with negligible hysteresis, enhanced stability and higher repeatability.

#### 2. Experimental

Preparation of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His: The glass/ITO substrates were successively cleaned with deionized (DI) water, detergent, ethanol, acetone and isopropanol under ultrasonic condition for 30 min. Afterwards, all substrates were dried in an oven at  $60^{\circ}$  for 2 h, and then further cleaning with UV/Ozone (UV-O3) for 20 min to remove surface functional groups before use. The original SnO<sub>2</sub> precursor solution was prepared by diluting the SnO<sub>2</sub> colloidal dispersion with deionized (DI) water (volume ratio of 1:1) and stirring for 3 h. Dissolve L-His in 1 mL DI water and stir for 3 h to obtain L-His solution with different concentration. The original SnO<sub>2</sub> precursor solution was spin-coated on the ITO substrate at a speed of 5000 rpm for 30 s, and then annealed on a hot plate at 150 °C for 30 min to obtain the ITO substrate at a speed of 5000 rpm for 30 s, and annealed on a hot plate at 150 °C for 20 min to obtain the buffer layer.

The fabrication process of PSCs and characterization conditions are similar with our previous reports [11,15], which can be found in the supporting information.

## 3. Results and discussion

To optimize the concentration of L-His, a series of PSCs with Glass/ ITO/SnO<sub>2</sub>/CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>/Spiro-OMeTAD/Ag structure (Fig. 1a) have been fabricated. Fig. S1 depicts the detailed preparation process of SnO<sub>2</sub>/L-His/Perovskite films. The L-His with different concentration is spun coated between the ETL and the perovskite layer. The photovoltaic performance parameters for the L-His/SnO2-based PSCs with different L-His contents are provided in Fig. S2. To check the reproducibility of PSCs, a batch of ten devices for each concentration of L-His was made. The average values in performances were obtained accordingly, which were then used to evaluate the experimental error. In the following analysis, the experimental error values were obtained by minus the average value and then divided by 2. According to Fig. S2, we can see that the optimal concentration of L-His is 0.06 mM, since the open circuit voltage ( $V_{oc}$ ), short circuit current ( $J_{sc}$ ) and fill factor (FF) of the devices are the highest at this concentration. In addition, forty devices were tested under standard light and the power conversion efficiency (PCE) of these PSCs at this concentration was generally higher. When the concentration of L-His is lower than 0.06 mM, L-His cannot completely and uniformly cover the rough SnO<sub>2</sub>, which cannot make SnO<sub>2</sub> achieve a better dispersion and lead to a decrease in the electron transport capacity [24]. When the concentration of L-His increased to 0.08 mM, the excessive L-His can damage the charge carrier mobility of SnO<sub>2</sub> layer due to the poor conductivity of L-His [25], so that the PCE of -PSCs was reduced. The optimal concentration of 0.06 mM was then used for the following preparation of perovskite films and devices, which was named SnO<sub>2</sub>/L-His based perovskite films and PSCs, respectively.

The effects of L-His on regulating the morphology and photoelectric properties of  $SnO_2$  ETLs were subsequently explored. Firstly, we used scanning electron microscopy (SEM) and atomic force microscopy (AFM) to investigate the surface morphology of  $SnO_2$  and  $SnO_2/L$ -His films. As shown in Fig. 1b and c, both films are compact without pinholes, but  $SnO_2/L$ -His film is flatter than  $SnO_2$  film. Accordingly, the root-mean-square roughness (RMS) of  $SnO_2$  and  $SnO_2/L$ -His films deduced from the AFM images in Fig. 1d and 1e is 1.4 nm and 1.1 nm, respectively. A slight decrease in roughness combined with SEM results prove that L-His effectively improve the uniformity and flatness of ETLs, preventing the direct contact between perovskite film and ITO electrode, which is beneficial for improving the charge transport capacity of  $SnO_2$  ETL [26].



Fig. 1. Characterization of SnO<sub>2</sub>/L-His ETL. a. Schematic diagram of the device structure with molecular structure of L-His. SEM images of b. SnO<sub>2</sub> and c. SnO<sub>2</sub>/L-His. AFM images of d. SnO<sub>2</sub> and e. SnO<sub>2</sub>/L-His.

Fig. S3 shows the XRD patterns of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His films deposited on soda lime glass at room temperature. The three diffraction peaks at 26.6°, 33.8° and 51.8° can be indexed to the (110), (101) and (211) crystal plane, respectively [27], indicating the formation of SnO<sub>2</sub> with tetragonal rutile structure according to PDF: No. 88–0287. The SnO<sub>2</sub>/L-His peaks are essentially unchanged compared to the crystal structure of SnO<sub>2</sub>. Meanwhile, the similarity of UV–Vis transmittance spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His films shown in Fig. S4 demonstrates that the L-His does not affect the light transmittance [28]. To illustrate the effect of L-His on the conductivity of SnO<sub>2</sub> ETL, the ETL only devices with architecture of ITO/SnO<sub>2</sub>/Ag were fabricated. The corresponding current-voltage (I-V) curves are shown in Fig. S5. Obviously, the I-V curve of the L-His-modified SnO<sub>2</sub> ETL is steeper than that of the SnO<sub>2</sub>

ETL. The conductivity of ETL film can be calculated using the following formula [29].

$$\sigma_0 = d/(AR) \tag{1}$$

where  $\sigma_0$  is the conductivity, d is the thickness of the ETL, A is the activity area (0.09 cm<sup>2</sup>), and R is the resistance [30]. Combining the cross-sectional SEM images in Fig. 3 and AFM height distribution measurements in Fig. S6, the thickness of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His film can be determined to be 23.2 nm and 24.3 nm, respectively. Thus, the calculated value of  $\sigma_0$  for SnO<sub>2</sub> and SnO<sub>2</sub>/L-His is  $2.97 \times 10^{-6}$  S cm<sup>-1</sup> and 4.14  $\times 10^{-6}$  S cm<sup>-1</sup>, respectively. Such obvious increase in conductivity indicates that the introduction of L-His can greatly improve the charge transport efficiency within the PSCs. Accordingly, the Hall Effect test (Fig. S7) also shows a decreasing tendency in resistivity. Therefore,



**Fig. 2.** The interaction between SnO<sub>2</sub> and L-His. a. FTIR spectra of SnO<sub>2</sub>, L-His and SnO<sub>2</sub>/L-His films. b. XPS survey spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His films. c. High-resolution N 1 s XPS spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His films. d. High-resolution Sn 3d XPS spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His films. Fitting curves for Sn 3d XPS spectra of e. SnO<sub>2</sub> and f. SnO<sub>2</sub>/L-His films. Fitting curves for O 1 s XPS spectra of g. SnO<sub>2</sub> and h. SnO<sub>2</sub>/L-His films.

using L-His to modify SnO<sub>2</sub> simultaneously can improve the surface quality and enhance the electrical property, which not only provides a desired platform for depositing high quality perovskite films, but also brings expectation to achieve higher efficiency of charge extraction and collection [31].

Fourier transform infrared spectroscopy (FTIR) and X-ray photoelectron spectra (XPS) techniques were utilized to unveil the relative influence mechanism of L-His on SnO<sub>2</sub> for resulting in the above improvement. Fig. 2a shows the FTIR spectra of pure L-His, SnO2 and SnO<sub>2</sub>/L-His samples. As shown in the FTIR spectrum of SnO<sub>2</sub> colloidal dispersion in deionized water, the peaks located at 588 cm<sup>-1</sup> and 948 cm<sup>-1</sup> can be assigned to the stretching vibration of O-Sn-O and asymmetrical Sn-O stretching vibration modes, respectively. Besides, a strong band near 1650 cm<sup>-1</sup> caused by O-H stretching vibration mode also appeared due to the existence of  $H_2O$  [32]. As shown in the FTIR spectrum of pure L-His powder, the characteristic vibrational peaks of C=O and C-O at 1633 and 1250 cm<sup>-1</sup> have been illustrated, which are attributed to the stretching vibration caused by the carboxyl group in L-His [33]. The peak at 1145 cm<sup>-1</sup> corresponds to the N-H vibrational stretching mode caused by the amino group [34]. Meanwhile, other two peaks located at 1590 and 1271 cm<sup>-1</sup> are corresponding to the characteristic C=N and C-N vibrational peaks modes, respectively, which are originated from the imidazole ring in L-His [34]. Compared to these spectra, from the FTIR spectrum of the SnO<sub>2</sub>/L-His sample, we can not only see the O-Sn-O and Sn-O peaks originated from SnO<sub>2</sub>, but also the C=O, C-O, N-H and C-N peaks originated from L-His, indicating that the L-His is successfully modified on the SnO<sub>2</sub> films. Moreover, we can see that the N-H peak in SnO<sub>2</sub>/L-His maintains at the initial position of L-His, indicating that the amino group of L-His does not interact with SnO<sub>2</sub> [32]. The C-N peak shifts to 1266 cm<sup>-1</sup>, which can be ascribed to the formation of hydrogen bonds with water molecules [35]. It is worth to notice that the peaks of C=O, C-O in SnO<sub>2</sub>/L-His shifted to 1621 and 1230 cm<sup>-1</sup>, respectively. Accordingly, the peak of O-Sn-O and Sn-O in SnO<sub>2</sub>/L-His shift to 613 and 951 cm<sup>-1</sup>, respectively. Such obvious interactive shifts indicate that the chemical interaction between -COOH in L-His and SnO<sub>2</sub> have occurred [34].

To further testify it, the XPS survey spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His films deposited on ITO substrates are shown in Fig. 2b. All spectra are all calibrated with the C 1s peak (284.8 eV). Besides Sn, O and C elements appeared in both samples, a strong N 1s peak at 400 eV also appears in the spectrum of SnO<sub>2</sub>/L-His film. To see it clearly, the high-resolute scan N 1s XPS spectra of both samples are presented in Fig. 2c. After fitting, three peaks located at 398.6 eV, 400.0 eV and 400.9 eV can be assigned to pyridine-like N, pyrrole-like N and -NH<sub>2</sub> group [36], respectively, further proving the successful modification of L-His on SnO<sub>2</sub> film. Since the L-His belongs to Lewis base due to the existence of carboxyl group, the O atom in the carboxyl group contains two lone pairs of electrons, which possesses strong electronegativity to facilitate the coordination with Sn ion in SnO<sub>2</sub> [30]. To analyze the possible reaction that happens in SnO<sub>2</sub>/L-His film, Sn 3d and O 1 s XPS spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His are shown in Fig. 2d-h. Due to the asymmetry of both spectra, we further fitted Sn 3d XPS spectra with Gaussian-Lorentzian functions and presented them in Fig. 2e,f. Two deconvolution fitting peaks located at 486.5 eV and 485.8 eV are originated from  $\text{Sn}^{4+}$  and  $\text{Sn}^{2+}$ , respectively, indicating the coexistence of  $\text{Sn}^{4+}$  and  $\text{Sn}^{2+}$  in  $\text{SnO}_2$  films [37]. The  $\text{Sn}^{2+}$ oxidation state is caused by the oxygen vacancies formed during spin coating, and its binding energy is higher than that of the Sn<sup>4+</sup> oxidation state [38]. We calculated the relative ratio of  $\text{Sn}^{4+}/\text{Sn}^{2+}$ , and the peak intensity ratio of  $\text{Sn}^{4+}/\text{Sn}^{2+}$  increased from 1.98 for  $\text{SnO}_2$  film to 3.71 for SnO<sub>2</sub>/L-His film. Meanwhile, after depositing L-His on SnO<sub>2</sub>, the Sn 3d peak position of  $Sn^{4+}$  and  $Sn^{2+}$  shifts to higher binding energy of 486.8 and 486.1 eV, respectively. Obviously, these results might because that the L-His brings the oxygen containing groups of -COOH into the SnO<sub>2</sub> films, which can coordinate with Sn<sup>4+</sup> to effectively reduce the surface oxygen vacancies of SnO<sub>2</sub> films [39].

films to unveil the detailed reasons for the reduced  $\text{Sn}^{2+}$  and Sn 3d peak shift. Fig. 2g shows a typical O 1 s XPS spectrum of pristine SnO<sub>2</sub>, which can be fitted with three independent peaks at 530.5 eV, 531.3 eV, and 532.7 eV, corresponding to lattice oxygen (OI), oxygen vacancy (OII), and chemisorbed oxygen (OIII), respectively [40]. The peak intensity ratio of I<sub>(OII)</sub>/I<sub>(OI)</sub> decreases from 0.62 of SnO<sub>2</sub> film to 0.54 of SnO<sub>2</sub>/L-His film. Meanwhile, the OI peak of SnO<sub>2</sub>/L-His film (Fig. 2h) exhibits an increase in the binding energy by about 0.2 eV. These results prove that the L-His modification effectively reduces the surface oxygen vacancies of SnO<sub>2</sub> films due to a chemical bonding reaction between Sn<sup>4+</sup> and -COOH, which is also consistent with the results with FTIR.

Combining the results of FTIR and XPS spectra, Fig. S8 shows a schematic diagram of L-His passivation mechanism. As we know, the L-His structure has a strong electron-withdrawing group composed by a carboxyl group connected to an aromatic ring. Therefore, compared to the control  $\text{SnO}_2$  film, the addition of L-His can increase the number of electrons around the Sn ions and prompt the chemical bonding reaction between  $\text{Sn}^{4+}$  and -COOH, which not only suppresses the oxygen vacancies, but also causes the binding energy of Sn ions move to the higher binding energy [41].

Effect on crystallization of perovskite films. In order to further study the effect of modified ETLs on perovskite films, Fig. 3a, b shows the top-view SEM images of perovskite films deposited on SnO<sub>2</sub> and SnO<sub>2</sub>/L-His ETLs. A few pinholes appear on the SnO<sub>2</sub>-based perovskite film. But the SnO<sub>2</sub>/ L-His based perovskite film contains continuous and regular perovskite grains with larger size. We further performed corresponding crosssectional SEM images and showed them in Fig. 3c,d. Unsurprisingly, the SnO<sub>2</sub> based perovskite film shows a stratified structure with many small grains existing in the films. In contrast, the SnO<sub>2</sub>/L-His based perovskite film shows columnar crystal grains from the bottom to the top. These columnar grains directly connect the ETL and hole transport layer (HTL), which facilitates the charge transport of perovskite layer [42]. The corresponding XRD patterns are shown in Fig. S9. Both samples show diffraction peaks at 14.2° and 28.4°, representing (100) and (220) lattice plane of perovskite with tetragonal structure, respectively [43]. The existence of L-His does not affect the crystal structure and growth orientation of CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> crystals. It is noteworthy that the SnO<sub>2</sub>/L-His based perovskite film owns stronger diffraction intensity, indicating a better crystallinity has been obtained [44], which is consistent with the variation tendency of SEM results. The mean crystals size in the primary crystal planes of the perovskite material is calculate by XRD analysis using Scherrer's formula:

$$D = K\lambda/(\beta \cos\theta) \tag{2}$$

where K is the Scherrer constant (0.89), D is the grain size (nm),  $\beta$  is the half-height width of the diffraction peak, which needs to be converted into radians (rad).  $\theta$  is the glancing angle,  $\lambda$  is the X-ray wavelength (0.154056 nm) [45,46]. In order to make the data more accurate as well as comparable, we chose the (110), (220) and (310) diffraction main peaks to do the calculation. The mean crystal sizes of control and SnO<sub>2</sub>/L-His based perovskite films is determined to be 30 nm and 34 nm, respectively. This particle size is also called the "subgrain size", which characterizes the "unit cell" size of the large grains to form the perovskite films. The increased subgrain size indicates that adding L-His at the interface can increase the grain size of the perovskite film and improve the crystalline quality [46].

Previous studies have shown that the nucleation process of perovskite is very sensitive to the surface energy of the substrate [47]. Then Fig. S10 shows the contact angle images of perovskite precursor solution on  $SnO_2$  and  $SnO_2/L$ -His ETLs. Compared to the pristine  $SnO_2$  ETL, the contact angle is reduced from 19.8° to 9.3° with introducing L-His layer. Such improved wettability facilitates the formation of the perovskite film [48,49]. In detail, the Gibbs free energy for heterogeneous nucleation can be expressed as Eq. (3):

(3)

We then further provide the core-level lines of O 1 s of both SnO<sub>2</sub>

$$\Delta G_{heterogeneous} = \Delta G_{homogeneous} \times f(\theta)$$



**Fig. 3.** Characterization of perovskite. Top-view scanning electron microscope (SEM) images of perovskite films coated on **a**. SnO<sub>2</sub> and **b**. SnO<sub>2</sub>/L-His substrates. The cross-sectional SEM images of **c**. SnO<sub>2</sub>-based and **d**. SnO<sub>2</sub>/L-His-based complete planar-type PSCs devices. Photographs of the perovskite films **e**. without and **f**. with L-His collected at the subsequent delay times after posing the substrates on a hot plate at 100 °C.

wherein  $f(\theta) = (2-3 \cos\theta + \cos^3\theta)/4$ , and  $\theta$  is the contact angle of the precursor solution [50]. Since the magnitude of  $\theta$  varies in the range of [0,  $\pi/2$ ], the larger the  $\theta$  is, the larger parameter  $f(\theta) \in [0, 1]$  is. In other words, the smaller contact angle makes the Gibbs free energy of heterogeneous nucleation lower, thereby resulting in the slower nucleation process, which is beneficial for obtaining larger grain size in perovskite films [51]. Obviously, the reduced Gibbs free energy caused by

introducing L-His in our case is apt to retard the crystallization process [52]. To further prove it, we carried out the experiment in the glove box and collected the photographs of crystallization process of perovskite films every 10 s under 100°C annealing conditions, which has been shown in Fig. 3e-f. It is clear that the  $SnO_2/L$ -His based perovskite films turn black slower than the control one. Therefore, a better crystallinity and enlarged grain size have been achieved in the perovskite deposited



**Fig. 4.** Charge transport between perovskite and different ETLs. **a.** Steady-state PL and **b.** TRPL spectra of perovskite films deposited on SnO<sub>2</sub> and SnO<sub>2</sub>/L-His ETLs. **c.** J-V characteristics of the SnO<sub>2</sub>-device and SnO<sub>2</sub>/L-His-device measured in the dark. **d.** Dark I-V curves of the electron-only devices. **e.** Transient photovoltage decay of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based PSCs **f.** EIS of planar-type PSCs with different ETLs.

on the SnO<sub>2</sub>/L-His [53]. Clearly, it exhibits the similar variation tendency as those under the atmospheric environment (Fig. S11), indicating that the L-His modification indeed delay the crystallization process of perovskite films.

Charge transfer dynamics. Fig. 4a shows the steady-state photoluminescence (PL) spectra of SnO<sub>2</sub>/perovskite and SnO<sub>2</sub>/L-His/perovskite. A significant PL quenching was observed in the SnO2/L-His/ perovskite compared to the control sample, which indicates that an efficient electron injection from perovskite into ETL has been achieved after introducing L-His on SnO2 ETL [54]. In addition, the blue shift of the emission peak indicates that the defect states at the ETL/perovskite interface and the perovskite grain boundary are reduced, which is consistent with XRD and SEM results [55]. Fig. 4b shows the corresponding normalized time-resolved photoluminescence (TRPL) spectra. Both spectra can be fitted by a bi-exponent decay function of  $I = A_1 \exp \left(\frac{1}{2}\right)$  $(-t/\tau_1)+A_2 \exp(-t/\tau_2)$ . The detailed fitting parameters are shown in Table S1. The average lifetime of SnO<sub>2</sub>/perovskite and SnO<sub>2</sub>/L-His/perovskite is 251.35 and 93.50 ns, respectively. Generally, the fast decay  $(\tau_1)$  can be attributed to the quenching process caused by the charge separation and the slow decay  $(\tau_2)$  can be attributed to the radiative recombination process of charge carriers [56]. Both  $\tau_1$  and  $\tau_2$  significantly decreased when the perovskite film is deposited on SnO<sub>2</sub>/L-His, indicating that the electrons have been efficiently extracted from the perovskite layer into SnO<sub>2</sub>/L-His and the charge separation and transport have been accelerated. Fig. 4c shows the dark J-V curves of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based PSCs. The leakage current of the SnO<sub>2</sub>/L-His based device is smaller than that of control device, indicating that the SnO<sub>2</sub>/L-His based device can prevent leakage and improve the photovoltaic performance.

The electron trap density (N<sub>t</sub>) within perovskite film is one of key factor to determine the charge transport process. Thus, the electron-only devices with the structure of ITO/ETL/perovskite/PCBM/Ag were fabricated to evaluate the N<sub>t</sub> of perovskite deposited on different substrates. Fig. 4d shows the corresponding dark current voltage (I-V) curves of both devices based on SnO<sub>2</sub> and SnO<sub>2</sub>/L-His ETLs. The linear correlation reveals an ohmic response at low bias voltages, and a nonlinear increase in current appears when the bias voltage is higher than the trap-filling limit voltage (V<sub>TFL</sub>, Fig. 4d), indicating that the traps are completely filled. The N<sub>t</sub> can be obtained from Eq. (4) [57]

$$N_t = 2\varepsilon_0 \varepsilon V/eL^2$$
(4)

where L is the thickness of the perovskite film, q  $(1.602 \times 10^{-19} C)$  is the elementary charge of the electron,  $\epsilon_r$  (32) is the relative dielectric and  $\epsilon_0$  (8.8542  $\times$   $10^{-14}$  F/cm) is the vacuum permittivity. From the cross-sectional SEM images (Fig. 3c,d), the SnO<sub>2</sub>-based perovskite film and SnO<sub>2</sub>/L-His based perovskite film thickness can be determined to be about 458 nm and 464 nm, respectively. The N<sub>t</sub> of the SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based perovskite film can be calculated to be  $2.631 \times 10^{15}$  cm<sup>-3</sup> and  $2.135 \times 10^{15}$  cm<sup>-3</sup>, respectively. The significantly reduced trap density greatly contributes to the faster charge separation and extraction [58].

Transient photovoltage decay (TPV) of PSCs under the open-circuit condition is another effective way to gain insight into the carrier recombination dynamics and free carrier lifetime within PSCs [59]. The TPV spectra of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based PSCs were shown in Fig. 4e and the corresponding parameters are shown in Table S2. The charge-recombination lifetime ( $\tau_v$ ) of the SnO<sub>2</sub>/L-His based device was substantially longer than that of the device based on pure SnO<sub>2</sub> (148 µs versus 76 µs), which is consistent with the slow charge recombination in the SnO<sub>2</sub>/L-His based perovskite films derived from TRPL measurement [60]. These results further prove that introducing L-His on SnO<sub>2</sub> ETLs indeed gradually improves the crystal quality of ETL and perovskite films by passivating the defects.

The electrical impedance spectroscopy (EIS) was employed to obtain transfer resistance in the PSCs. Fig. 4f shows the Nyquist plots of PSCs based on  $SnO_2$  and  $SnO_2/L$ -His ETLs measured under dark conditions

with the equivalent circuit shown in the inset. The corresponding fitting parameters including series resistance (R<sub>s</sub>), transfer resistance (R<sub>tr</sub>), and recombination resistance (Rrec) are listed in Table S3. The Rs mainly reflects the electron transport capability, the high-frequency component is the signature of the Rtr and the low frequency component is characteristic of the R<sub>rec</sub> [61]. It is worth noting that the SnO<sub>2</sub>/L-His based device possesses a smaller Rs value compared with the control device. Meanwhile, it also has the smallest  $R_{tr}$  of  $4.59{\times}10^3\,\Omega$  and the largest  $R_{rec}$ of  $4.29 \times 10^7 \Omega$  compared to the SnO<sub>2</sub>-based PSCs. The smaller R<sub>tr</sub> favors the extraction of electrons, while the large  $R_{\rm rec}$  effectively resists charge recombination, which is consistent with the above observations. In our case, the perovskite/HTL interface is identical for both devices and the only factor affecting these resistances is the perovskite/ETL interface. Thus, SnO<sub>2</sub>/L-His is further proved to be a desired ETL for planar PSC due to its great capacity to improve the charge transfer efficiency and suppress the recombination process.

The improvement of charge transport performance is usually related to the factors such as interface bonding and energy level matching. The L-His covered on the SnO<sub>2</sub> films is inevitably contact with perovskite precursor. To unveil the improvement of the charge transport performance of SnO<sub>2</sub>/L-His based perovskite film, the liquid-state <sup>1</sup>H NMR spectra of perovskite precursors with and without L-His in deuterated Dimethyl sulfoxide-d6 (DMSO-d6) are depicted in Fig. 5a-c. Two single peaks at  $\delta = 7.34$  ppm and  $\delta = 2.30$  ppm can be attributed to the NH<sub>3</sub><sup>+</sup> and CH<sub>3</sub> of methylammonium, respectively [62]. Fig. 5b shows a typical peak of perovskite (2.30 ppm), indicating the introduction of L-His does not affect the chemical structure of perovskite. However, the introduction of L-His shifted the secondary peak of perovskite slightly (Fig. 5c, 7.32 ppm), proving that hydrogen bond has been formed between the hydrogen of L-His and the halide of perovskite, thus promoting charge transport performance [63]. The hydrogen bond interaction between perovskite and L-His is shown in Fig. 5d.

The degree of energy band matching at the interface between ETL and perovskite films also determines the kinetics of charge extraction and combination [64]. The L-His deposited on  $SnO_2$  as SAM is inevitable to change the work function of ETLs. Therefore, the ultraviolet photoelectron spectroscopy (UPS) spectra of SnO2 and SnO2/L-His were performed and shown in Fig. 5e, f to evaluate their energy level. The cut-off binding energy ( $E_{cut off}$ ) of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His is estimated to be 16.70 and 17.10 eV, respectively. Then, the Fermi level (E<sub>F</sub>) of SnO<sub>2</sub> and  $SnO_2/L$ -His is calculated to be -4.52 and -4.12 eV using the equation  $E_F$ = $E_{cut\ off}$  -21.22 eV [65], respectively. The valence band maximum (VBM) of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His is 3.80 and 3.70 eV, respectively. Using the Tauc plot shown in Fig. 5g, the optical band gap  $(E_g)$  of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His is calculated to be 3.80 and 3.82 eV, respectively. Then the conduction band maximum (CBM) can be determined to be -4.52 and -4.00 eV. Based on these results, the possible energy level diagram of PSCs based on both ETLs has been illustrated in Fig. 5h. It can be seen that the increase of conduction band energy in SnO<sub>2</sub>/L-His ELTs can reduce the band offset between CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> and ETLs. Thus, the introduction of L-His layer successfully shifts the energy level of the SnO2 ETLs and the Fermi level of SnO2/L-His matched the conduction band of the perovskite well, which is another key factor to facilitate electron extraction and transport [66]. It also helps to achieve a higher open circuit voltage [67], which will be seen in the following discussion.

The photovoltaic performance of PSCs. With the superior optoelectronic properties discussed above, it is expected that the SnO<sub>2</sub>/L-His would be a better ETL in the PSCs than the SnO<sub>2</sub>. Therefore, planar PSCs were designed and fabricated based on different ETL. As shown in Fig. 6a, the J-V curves of planar PSCs based on different ETL with the key parameters including  $J_{sc}$ ,  $V_{oc}$ , FF, and PCE summarized in the inset of Fig. 6a. The PCE of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based PSCs is  $20.05\pm0.17\%$  and  $21.04\pm0.14\%$ , respectively. Fig. 6b shows the incident photon-to-charge conversion efficiency (IPCE) and the integrated current of the PSCs based on different ETLs. It is apparent that the device based on the SnO<sub>2</sub>/L-His shows significantly higher IPCE. The integrated current



**Fig. 5. Exploration of charge transport performance. a–c.** Proton nuclear magnetic resonance ( $^{1}$ H NMR) analysis.  $\checkmark$  represents the perovskite peak;  $\blacklozenge$  represents the L-His peak. **d.** Schematic diagram of the hydrogen-bond formation between perovskite and L-His. UPS spectra of **e.** SnO<sub>2</sub> film and **f.** SnO<sub>2</sub>/L-His film. **g.** Tauc plots of perovskite film deposited on SnO<sub>2</sub> and SnO<sub>2</sub>/L-His film. The inset shows the UV–vis absorption spectra of perovskite films deposited on different substrates. **h.** Energy level diagram of PSCs.



**Fig. 6.** The photovoltaic performance of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based planar PSCs devices. **a.** The typical J–V curves of SnO<sub>2</sub> and SnO<sub>2</sub>/L-His based PSCs. **b.** IPCE and integrated current density corresponding to planar PSCs with various ETLs. **c.** Variation of static current density and PCE measured with time. **d.** Long-term stability of PCE via storage time without any encapsulation under ambient condition.

value calculated by the IPCE spectra for the devices using  $SnO_2$  and  $SnO_2/L$ -His are 21.82 and 22.43 mA/cm<sup>2</sup>, respectively, which is very close to the J-V results [65].

Stability and hysteresis are two key characteristics of PSCs. Fig. S12 shows the J-V curves under reverse and forward scan. The hysteresis index (HI) calculated by HI = (PCE<sub>reverse</sub> - PCE<sub>forward</sub>)/PCE<sub>reverse</sub> is 5.2%

and 1% for  $SnO_2$  and  $SnO_2/L$ -His based devices, respectively [68]. The J-V curves of PSCs based on  $SnO_2$  ETLs showed a significant hysteresis. In general, the hysteresis of PSCs is attributed to the interfacial capacitance caused by the interfacial charge accumulation, which originates from ion migration, high trap density and unbalanced charge transport within the perovskite device [69]. According to the above results, both

the reduced trap density within devices and improved electron collection efficiency are reasonable to explain the hysteresis reduction [70].

The steady-state photocurrent density and PCE of the SnO<sub>2</sub> and SnO<sub>2</sub>/L-His devices were further measured at the maximum power point. As shown in Fig. 6c, the PCE of PSCs based on the SnO<sub>2</sub> and SnO<sub>2</sub>/L-His ETLs were stable at 19.86±0.16% and 20.83±0.14%, corresponding to photocurrent densities of 21.13 and 21.79 mA/cm<sup>2</sup>, respectively. The photocurrent densities and PCE of both devices exhibit the negligible change under continuous solar illumination for 100 s, which proves excellent irradiation stability of the cells [70]. Fig. 6d shows the PCE variation curve via a function of storage time for unpackaged devices at 35% humidity and room temperature, and more detailed J-V parameters summarized in Table S4. It is clear that the SnO<sub>2</sub>/L-His-based device maintains an initial efficiency of 82% in ambient air after 700 h, while the device of SnO<sub>2</sub>-based device provides only 56% initial efficiency under the same storage conditions.

#### 4. Conclusions

In summary, we propose a simple and effective method of L-His SAMs modification on SnO<sub>2</sub> ETL to comprehensively improve the photovoltaic performance of PSCs. First, the carboxyl group in L-His can coordinate with tin ions to fill the oxygen vacancies and improve the crystalline quality and conductivity of SnO<sub>2</sub> ETLs. Secondly, the introduction of L-His can improve the wettability of the interface and improve the crystallization of the perovskite film. Thirdly, the -NH2 in L-His can form hydrogen bond between perovskite and L-His to enhance the charge transport performance. In addition, the introduction of L-His can also optimizes the energy level structure of the ETL/perovskite interface to further accelerate charge transfer. As a result, a maximum PCE of 21.04% with almost no hysteresis and enhanced stability in CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>-based PSCs has been achieved on the L-His-modified SnO<sub>2</sub> ETL. Such biological friendly modification technique is highly compatibility to the technique of large-scale production of PSCs, and the elucidation of relevant mechanisms in this work is very beneficial for understanding the functions of other amino acids in the future.

## CRediT authorship contribution statement

Xiaoxu Sun: Formal analysis, Methodology, Investigation, Writing – original draft. Haipeng Jiang: Conceptualization, Investigation. Yansen Sun: Software. Zonghan Guo: Validation, Investigation. Zhenyu Pang: Software. Fengyou Wang: Conceptualization, Validation. Jinghai Yang: Resources, Supervision. Lili Yang: Resources, Supervision, Writing – review & editing.

#### **Declaration of Competing Interest**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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#### Supplementary materials

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