

Large-Area Monolayer MoS₂ Nanosheets on GaN Substrates for Light-Emitting Diodes and Valley-Spin Electronic Devices

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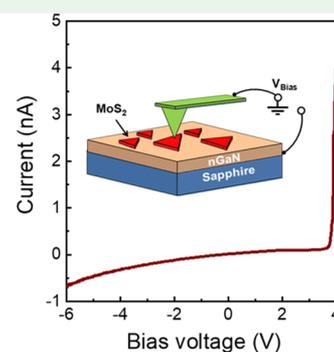
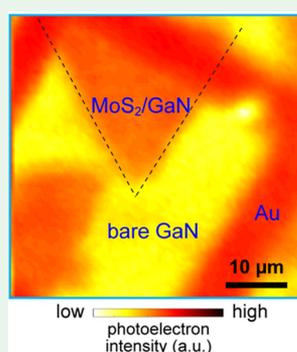
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ABSTRACT: Van der Waals heterostructures composed of atomically thin two-dimensional (2D) materials and three-dimensional (3D) materials provide a multidimensional material integration strategy, which combines materials with different characteristics leading to a wider degree of freedom than a single component, and offer a way for developing electronic and optoelectronic devices with multifunctionalities, such as high-frequency electronic devices, photodetectors, valley-spin electronic devices, and so on. This report demonstrates the direct growth of large-area monolayer MoS₂ single-crystal nanosheets with a side length of more than 100 μm on 3D GaN substrates by the perylene-3,4,9,10-tetracarboxylic acid tetrapotassium salt (PTAS) seed-assisted chemical vapor deposition (CVD) method. The seeding promoters changed the growth kinetics of MoS₂ on the GaN substrate, which is different from the previously reported epitaxial growth behavior. The size of our synthesized single-crystal MoS₂ nanosheets is 2 orders of magnitude larger than the reported epitaxially grown MoS₂ on the GaN substrate. Meanwhile, the as-synthesized MoS₂ by the seed-assisted CVD method has comparable crystal quality as that of the reported epitaxially grown MoS₂ on the GaN substrate. Moreover, detailed characterizations indicate that noticeable charge transfer occurs between MoS₂ and the GaN substrate, which suggests that the MoS₂/GaN heterostructure has great potential applications in the field of light-emitting diodes (LED) and valley-spin electronic devices.

KEYWORDS: MoS₂, chemical vapor deposition, seeding promoters, 2D/3D heterostructures, charge transfer



INTRODUCTION

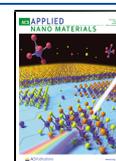
In recent years, van der Waals heterostructures composed of newly emerged two-dimensional (2D) materials and traditional three-dimensional (3D) materials have attracted extensive attention.^{1–6} The coupling of quasi-particles at the interface of heterostructures provides a new route for developing innovative device designs such as tunnel diodes, ultrabroad-band photodetectors, light-emitting diodes (LEDs), and high-power devices.^{7–13} As a typical III–V compound, GaN is known as a third-generation semiconductor after silicon and gallium arsenide. The crystal structure of GaN is a hexagonal wurtzite configuration, as the inner Ga and N atoms are all covalently bonded with each other, which is similar to silicon and diamond materials. Therefore, GaN can be considered a typical 3D material. Also, GaN is a wide direct-band-gap semiconductor with a bandgap of 3.4 eV, which is widely used in high-power devices, advanced wireless communication, and high-frequency microwave devices.^{14,15} Good chemical stability at high temperatures and compatibility with mature process semiconductor technology make GaN a potential substrate to grow 2D semiconductor materials and construct 2D/3D

heterostructures, meeting the demands of new electronic devices with multifunctionality and further broadening the application fields of both 2D materials and 3D GaN.^{9–12} Monolayer MoS₂ is a typical 2D semiconductor that has been intensively studied in the past decade. Both GaN and MoS₂ belong to the hexagonal crystal system with similar lattice constants (3.19 Å for GaN, 3.16 Å for MoS₂); the in-plane lattice mismatch is less than 1%. Furthermore, the thermal expansion coefficients of GaN and MoS₂ are very close ($3.95 \times 10^{-6} \text{ K}^{-1}$ for GaN, $4.92 \times 10^{-6} \text{ K}^{-1}$ for MoS₂).^{16–18} Thus, we can anticipate that the direct growth of MoS₂ on the GaN substrate will not generate in-plane stress during the growth process, which can guarantee the high crystal quality of MoS₂.

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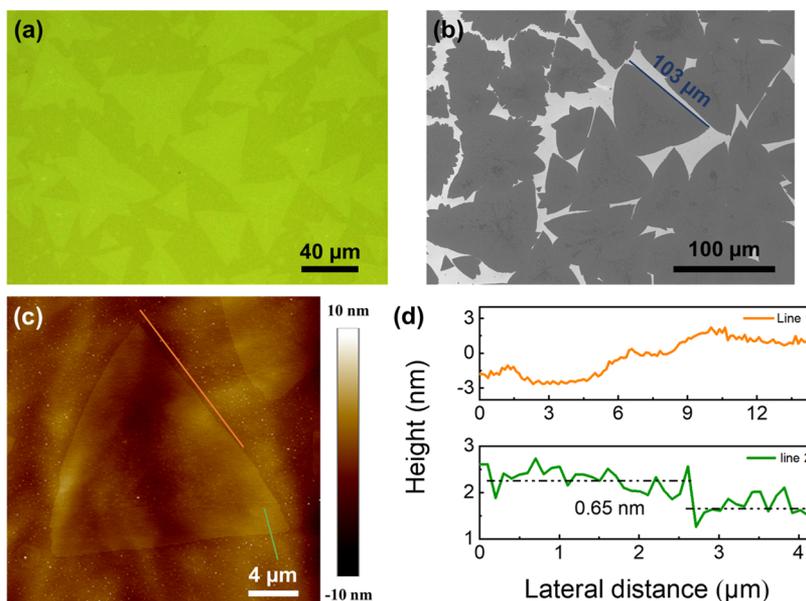


Figure 1. (a, b) Optical and SEM images of MoS₂ grown on GaN(0001)/sapphire substrates. (c) AFM image of the MoS₂/GaN heterostructure. (d) Height profile corresponding to the same color line marked in (c).

There are some interesting phenomena in the hybrid structure composed of the 3D GaN and the 2D MoS₂, such as the increased polarization of MoS₂ under circularly polarized light and the band bending at the interface of the heterostructures.^{11,19} At present, a transfer method was used to prepare 2D/3D heterostructures. For example, Moun et al. transferred the mechanically exfoliated MoS₂ to the GaN substrates and studied the photodetection and diode characteristics of the MoS₂/GaN heterostructure.²⁰ However, the exfoliated MoS₂ sample is not suitable for the preparation of large-scale integrated devices. Henck et al. transferred large-area chemical vapor deposition (CVD)-grown MoS₂ nanoflakes onto GaN substrates via the poly(methyl methacrylate) (PMMA)-assisted method. They found a significant charge-transfer phenomenon at the interface of the MoS₂/GaN heterostructure, whereas the wetting transfer method can cause damage to the crystal quality of MoS₂ during the chemical corrosion period and introduce some unwanted residual impurities.¹⁹ The CVD method has unique advantages in the preparation of large-area 2D/3D heterostructures due to its benefits of good controllability, relatively low cost, and good crystallinity of the obtained thin film. It is the most popular preparation technology currently for large-area 2D materials and their complex heterostructures.^{21–27} Ruzmetov and Wan et al. reported that CVD-grown MoS₂ nanosheets on GaN substrates exhibited apparent epitaxial growth behavior; the MoS₂ domains were aligned at either 0 or 60° orientations.^{9–11} However, the side length of the MoS₂ nanosheets obtained was generally below 10 μm, which limits its application in the field of large-scale devices.⁹

To overcome the bottleneck of size limit in CVD growth of monolayer MoS₂ epitaxially grown on traditional 3D GaN substrates, we used perylene-3,4,9,10-tetracarboxylic acid tetrapotassium salt (PTAS) seeds to promote CVD growth of large-area monolayer MoS₂ single crystals on GaN substrates. A triangle-shaped monolayer MoS₂ single-crystal nanosheet with a side length of more than 100 μm was synthesized on GaN substrates, which is 2 orders of magnitude larger than the size of MoS₂ crystal domains epitaxially grown

on the GaN substrates reported in the literature.^{9–11} Since the introduced PTAS catalyst changed the surface energy of the GaN growth substrate, the MoS₂ nanosheets do not exhibit epitaxial growth behavior on the substrate, that is, the MoS₂ nanosheets have no specific orientation with GaN substrates. A series of characterizations were conducted on the performance of the MoS₂/GaN heterostructure, including Raman spectroscopy, photoluminescence spectroscopy (PL), X-ray photoelectron spectroscopy (XPS), angle-resolved photoemission spectroscopy with submicron lateral resolution (micro-ARPES), Kelvin probe force microscopy (KPFM), and conductive atomic force microscopy (CAFM). The results indicate that the MoS₂ grown on GaN substrates has a good crystal quality, and there is a significant charge transfer between MoS₂ and GaN. The performance of our synthesized large-area MoS₂/GaN is comparable to that of the epitaxially grown MoS₂/GaN heterostructures. The direct growth of MoS₂/GaN heterostructures introduced in this paper promotes the large-scale preparation of 2D/3D heterostructures and their application in the field of optoelectronic devices.

EXPERIMENTAL SECTION

Growth of MoS₂. The MoS₂ sample was grown by the CVD method on n-type GaN(0001)/sapphire substrates in a two-temperature-zone furnace. The GaN(0001) film was epitaxially grown on Al₂O₃ with dislocation density less than $5 \times 10^8 \text{ cm}^{-2}$ and a thickness of 4.5 μm. Before growth of MoS₂, the GaN(0001)/sapphire substrates were pretreated by drop-casting 5 μL of 100 μM PTAS. Then, 150 mg of sulfur (Aladdin, 99.99%) and 3 mg of MoO₃ (Aladdin, 99.9%) were used as the growth precursors, and high-purity N₂ gas (99.999%) was used as the carrier gas. The two zones were separately heated to 170 and 680 °C within 38 min and maintained for 3 min for the growth with 20 sccm N₂ flow in an atmospheric pressure.

Characterization. The detailed morphologies of MoS₂ were characterized by optical microscopy (Keyence digital microscope VHX-600), scanning electron microscopy (SEM) (Zeiss Sigma HD), and atomic force microscope (AFM) (Bruker, Dimension FastScan) equipped with CAFM and KPFM test modules. The Raman and PL

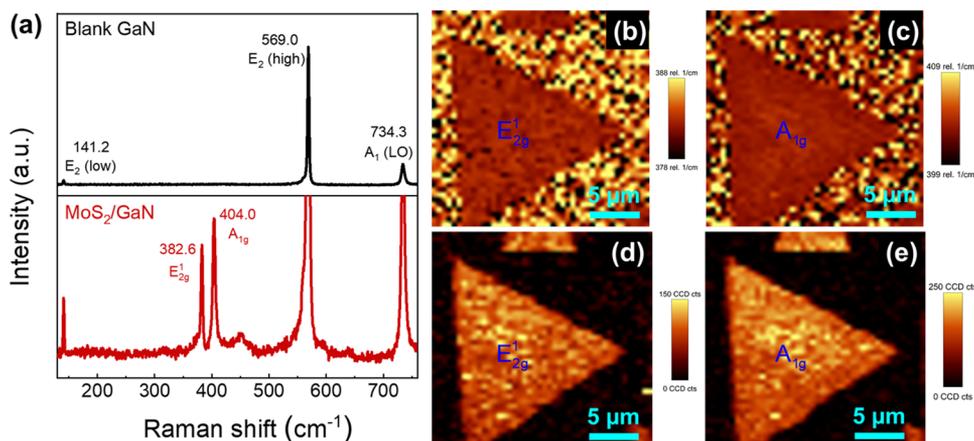


Figure 2. (a) Raman spectrum of blank GaN and MoS₂/GaN heterostructures. (b, c) Raman mapping of peak positions of E_{2g}¹ mode (b) and A_{1g} mode (c). (d, e) Raman mapping of peak intensities of E_{2g}¹ mode (d) and A_{1g} mode (e).

spectra were performed by WITec Alpha 300R with 532 or 633 nm laser. The laser power of 3 mW was chosen to obtain a Raman spectrum with a good signal-to-noise ratio. The XPS measurement was carried out by ThermoFisher ESCALAB 250Xi with an Al K α X-ray source. The electrodes were patterned by photolithography using a laser direct writing machine (Microwriter ML3) followed by deposition of 5 nm of Cr and 50 nm of Au. Micro-ARPES measurements were conducted at the SpectroMicroscopy beamline at the Elettra Synchrotron Light Source (Italy) utilizing a 74 eV photon beam.²⁸ The measurement temperature was 100 K. The overall energy, angle, and lateral resolutions were 50 meV, 1, and ~ 1 μm , respectively. Au mesh grids (40 \times 40 μm^2) were deposited on MoS₂/GaN by the laser direct writing machine and the lithography technology to avoid the charging effect in photoemission. Prior to ARPES measurements, MoS₂/GaN heterostructures were annealed at 573 K in a preparation chamber.

RESULTS AND DISCUSSION

The growth of 2D monolayer MoS₂ nanosheets in our experiment was conducted by the CVD method on 3D GaN substrates with special temperature control of the growth source as described in the **Experimental Section**. Figure 1a,b shows the morphology of MoS₂ grown on a GaN(0001)/sapphire substrate assisted-growth by PTAS seeding promoters under an optical microscope and SEM. It was found that the orientation of these MoS₂ domains is randomly distributed, and the maximum side length of triangular-shaped MoS₂ domains can reach up to 103 μm , which is 2 orders of magnitude larger than the MoS₂ domains epitaxially grown on the GaN substrate with a specific orientation by the CVD method in previous reports.^{9–11} This can be ascribed to the fact that the introduced seeds changed the surface energy of the GaN substrate and reduced the nucleation energy of MoS₂ during the growth period,^{29–31} making the growth of MoS₂ on GaN from thermodynamic equilibrium control with specific orientation become kinetic rate control with a random direction.³² Our growth strategy for the synthesis of MoS₂ single crystals with a side length of up to tens to hundreds of microns on GaN substrates has a lower growth temperature at 680 $^{\circ}\text{C}$ and a shorter growth time for 3 min compared with the epitaxial growth of MoS₂ on GaN at 800 $^{\circ}\text{C}$ for 15 min with side length less than 10 μm without seeding promoters as reported in the literature.⁹ The growth mechanism may be described as follows. Under the growth temperature of MoS₂, the PTAS seeds decompose into small fragments, which contain conjugate benzene rings as the parent molecule. The

plane structure of conjugate benzene rings can be absorbed on the surface of the GaN substrates and thus change the surface energy of GaN. At the same time, the conjugate benzene rings can also be highly reactive centers, which can anchor the growth source, that is, Mo and S atoms, in a plane and finally form a nucleation center, followed by the extended growth of MoS₂ in a planar way. With prolonged growth time, we can obtain a continuous monolayer MoS₂ film as seen in the following device sections. Further analysis of the morphology of MoS₂/GaN by AFM indicates that MoS₂ is a monolayer with a thickness of 0.65 nm, which is consistent with the thickness of the monolayer MoS₂ grown by CVD and the mechanical exfoliated sample.^{31,33} From the AFM morphology image, it can be seen that the surface of the MoS₂ flake is very clean and there are almost no impurity particles. On the other hand, we find that there are only a few incompletely reacted MoO_x molybdenum source particles in the blank GaN area as will be discussed by the XPS results below. Although the fluctuation of GaN topography exceeds 5 nm, the as-grown monolayer MoS₂ domains closely attach to the GaN surface. In view of the good chemical stability of GaN at high temperatures, GaN can also be a suitable substrate for growing other 2D transition-metal dichalcogenide (TMDC) compounds, thereby constructing various 2D/3D heterojunctions with more abundant properties.

The composition and crystal quality of the MoS₂/GaN heterostructure were studied with Raman spectroscopy. In general, Raman spectroscopy is a useful characterization tool that can quickly and nondestructively analyze the layer number and crystal quality of 2D materials.^{34,35} Figure 2 shows the Raman spectrum and mapping results of the MoS₂/GaN heterostructure. The upper panel of Figure 2a is the Raman spectrum of the blank GaN substrate after growing MoS₂. There are three typical Raman peaks located at 141.2 cm⁻¹ (E₂, low), 569.0 cm⁻¹ (E₂, high), and 734.3 cm⁻¹ (A₁ (LO)), consistent with the Raman features of GaN with a hexagonal wurtzite configuration as reported in the literature,^{36,37} indicating that the GaN substrate can be stable under the high-temperature growth environment of MoS₂. The lower panel of Figure 2a is the Raman spectrum of the MoS₂/GaN heterostructure. In addition to the Raman peaks of GaN, two other typical Raman peaks of MoS₂ can be detected at 382.6 cm⁻¹ and 404.0 cm⁻¹, which belong to the in-plane vibration mode (E_{2g}¹) and the out-of-plane vibration mode (A_{1g}),

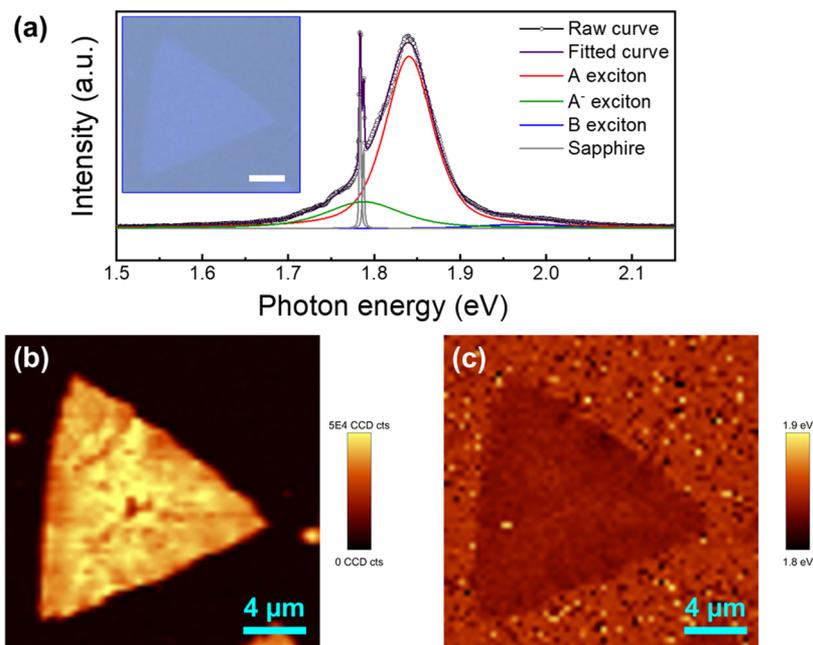


Figure 3. (a) PL spectrum of the monolayer MoS₂ domain grown on GaN; the inset is the optical image of the MoS₂/GaN heterostructure. The scale bar is 4 μm. (b, c) PL intensity mapping (b) and peak position mapping (c) of a typical monolayer MoS₂ domain on GaN corresponding to the inset in (a).

respectively. The difference between these two peaks is 21.4 cm⁻¹, which is slightly larger than that of the monolayer MoS₂ grown on SiO₂/Si wafers and sapphire substrates by CVD.^{29–31} Such a peak difference may be due to the thermal effect caused by the detection laser with a power of 3 mW.^{38,39} On the other hand, it may be caused by the charge transfer between MoS₂ and the GaN substrates, thus changing the doping level of MoS₂ as the peak position of the A_{1g} vibration mode is sensitive to doping.⁴⁰ Figure 2b,c shows the Raman mapping of the peak position of the two vibration modes (E_{2g}¹ and A_{1g}) of a typical MoS₂ domain. The peak positions of the two Raman modes have almost no change in the entire MoS₂ area, indicating that there is no local stress and local doping within the MoS₂ domain. Figure 2d,e shows the Raman mapping of the intensity of E_{2g}¹ and A_{1g} modes; the spatially uniform intensity of these two Raman modes indicates that the quality of the CVD-grown MoS₂ nanosheet on the GaN substrate is of good uniformity.

Then, we used PL spectroscopy to further evaluate the quality of the MoS₂ grown on the GaN substrate. Figure 3 shows the PL spectrum and mapping of the MoS₂/GaN heterostructure. Since monolayer MoS₂ is a direct-band-gap semiconductor, a strong PL peak is detected under 532 nm laser radiation, as shown in Figure 3a. The PL spectrum was fitted for better understanding the behavior of exciton in MoS₂; the peak located at 1.84 eV corresponds to the A exciton peak in MoS₂. The full width at half-maximum (FWHM) of A exciton is only 66 meV; such a narrow peak width indicates the high quality of the MoS₂ grown on GaN. The sharp peaks near 1.8 eV indicate the sapphire substrate. Due to the lack of inversion symmetry and the strong spin-orbital coupling effect in the monolayer MoS₂, leading to the splitting of the valence-band maximum,⁴¹ there are two direct exciton emission peaks, that is, A exciton and B exciton, as shown in the PL spectrum of MoS₂, but the B exciton peak is very weak. Recently, it has been reported that the intensity

ratio of A exciton versus B exciton (A/B) can be used to evaluate the crystal quality of MoS₂ qualitatively. The larger the ratio of A/B, the better the quality of MoS₂.^{42,43} The PL spectral shape of our synthesized MoS₂ on the GaN substrate is very similar to the suspended monolayer MoS₂ and the CVD-grown MoS₂ sample on h-BN,^{44,45} indicating that much less defects exist in MoS₂ samples, further illustrating the excellent quality of the MoS₂ grown on the GaN substrate. Figure 3b,c shows the PL mapping of the peak intensity and the peak position of the monolayer MoS₂ domain on the GaN substrate. It is also found that the peak intensity and the peak position of the entire crystal domain are relatively uniform, indicating that the quality of the MoS₂ crystal domain is uniform. However, the very small area with a darker intensity in the center of the mapping image can be due to the initial nucleation point of MoS₂.⁴⁶ The above optical properties further suggest that MoS₂ grown on GaN has a good crystal quality, which may benefit from the minor lattice mismatch ratio of MoS₂ and the GaN substrate, making GaN an excellent substrate for growing other high-quality 2D TMDC compounds that are lattice-matched with GaN. In addition, the circular polarization behavior of the MoS₂/GaN heterostructure is also studied. Figure S1a shows the circularly polarized PL spectrum of the MoS₂/GaN heterostructure under a 633 nm laser ($h\nu = 1.96$ eV) at room temperature. The use of an excitation wavelength of 633 nm here is due to the fact that resonant excitation can increase the polarization of 2D semiconductor materials under helicity circularly polarized laser because of the suppressed phonon-assisted intervalley scattering.¹¹ The following formula can calculate the polarization helicity

$$P = (I_{\sigma^+} - I_{\sigma^-}) / (I_{\sigma^+} + I_{\sigma^-}) \quad (1)$$

where I_{σ^+} and I_{σ^-} are the intensities of the PL spectra of MoS₂ under right-circularly and left-circularly light. It is found that MoS₂ on the GaN substrate has obvious valley polarization at

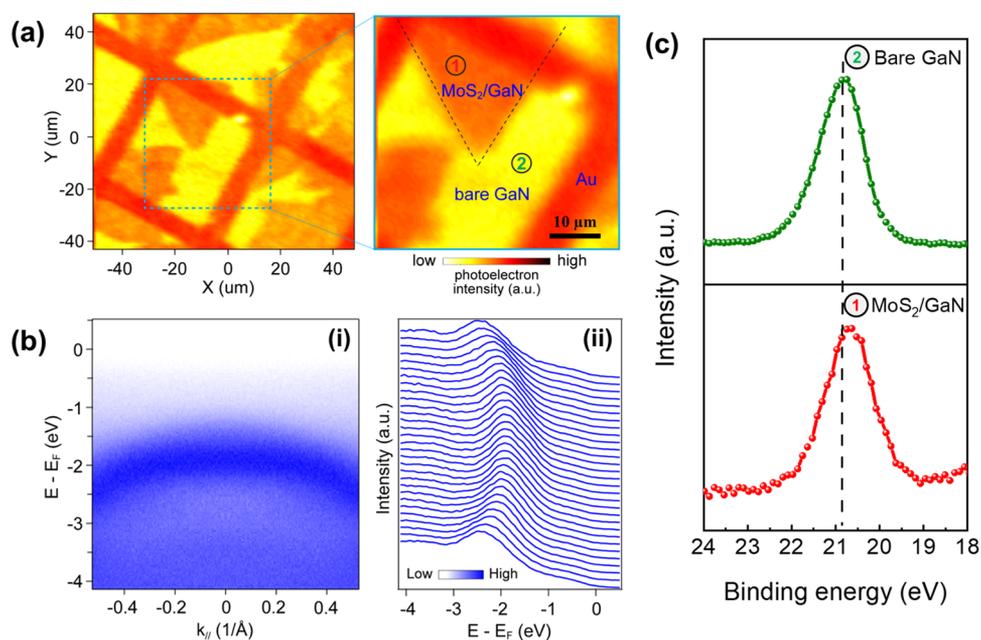


Figure 4. (a) Micro-ARPES photoemission real-space map of MoS₂/GaN measured with the photon energy of 74 eV. Photoelectrons with approximately the maximum kinetic energy are integrated. Regions of MoS₂/GaN, bare GaN, and Au mesh are marked. (b) ARPES photoemission spectra of the MoS₂/GaN heterostructure near the Γ point of the Brillouin zone (i) and its momentum–distribution curve plots (ii), both of which clearly show good holelike band dispersions of the valence band. (c) High-resolution photoelectron spectrum of the Ga 3d core-level electrons on bare GaN and MoS₂/GaN heterostructures corresponding to the marked dot in (a).

room temperature, which is much larger than that of MoS₂ grown on the sapphire substrate in the same condition as shown in Figure S1b, indicating that the interaction between the GaN substrate and MoS₂ can manipulate the valley helicity of MoS₂, which is due to the competition between the intervalley scattering rate and the faster exciton decay rate in MoS₂ caused by the GaN substrate.¹¹ The charge transfer in the interface of the MoS₂/GaN heterostructure increases the circular polarization of MoS₂, making the 2D/3D heterostructures composed of TMDCs and GaN have great promising application prospects in the field of valley electronics.

The composition of the MoS₂/GaN heterostructure was further analyzed by XPS as it is an important characterization technique for qualitative and quantitative analyses of the element component in samples. As shown in Figure S2a, the Mo 3d core-level spectra can be fitted with two pair peaks. The Mo 3d_{3/2} and 3d_{5/2} with band energies of 235.6 and 232.4 eV correspond to the incompletely reacted MoO_x particles as mentioned above measured on blank GaN by AFM. The other pair peaks of 3d_{3/2} at 232.6 eV and 3d_{5/2} at 229.5 eV correspond to Mo⁴⁺ in MoS₂. The spin–orbit splitting energy of Mo⁴⁺ 3d is 3.1 eV. These binding energies of Mo and S are consistent with the reported XPS peak values of MoS₂ in the literature.⁴⁷ At the same time, we also measured the XPS spectra of the Ga 3s core level as shown in Figure S2b, which is in line with the Ga 3s of GaN reported in the literature,⁴⁸ further indicating that GaN can be stable in the growth environment of MoS₂.

To investigate the band structure and electronic properties of the 2D/3D hybrid MoS₂/GaN heterostructure, we carried out state-of-the-art micro-ARPES measurements.^{28,49,50} After focusing the light to submicron size, we scanned the sample surface and acquired the photoemission 2D real-space map at the submicron lateral resolution as shown in Figure 4a, in

which regions of the triangle-shaped MoS₂/GaN heterostructure, bare GaN, and Au mesh can be apparently discerned. Then, we moved the light beam to the MoS₂/GaN heterostructure region and obtained the valence-band dispersion of MoS₂ at the Γ point of the Brillouin zone. Clearly, the valence band exhibits good holelike dispersions,¹⁹ as shown by the false-color display of the ARPES spectrum in Figure 4b(i) and massive plots of momentum–distribution curves in Figure 4b(ii). We note that the quality of our ARPES data is comparable to that measured on exfoliated MoS₂ nanoflakes.^{51,52} Interestingly, from the high-resolution photoelectron spectrum of the Ga 3d core-level electrons in the MoS₂/GaN heterostructure as shown in Figure 4c, we found that Ga 3d has 200 meV red-shifted in the MoS₂/GaN heterostructure compared to the blank GaN substrate. This shift at a lower binding energy suggests that band bending occurs at the interface of the heterostructure as a result of charge transfer from the 2D monolayer MoS₂ to GaN.¹⁹

To explain the charge-transfer phenomenon at the MoS₂/GaN interface, KPFM was used to characterize the surface potential mapping of the MoS₂/GaN heterostructure as shown in Figure S3. There is a significant potential difference between the MoS₂ domain and GaN, and the potential distribution is uniform throughout the whole MoS₂ domain. The contact potential difference (CPD) between the KPFM tip and the sample is given by

$$(\text{CPD})_{\text{MoS}_2} = \frac{\phi_{\text{tip}} - \phi_{\text{MoS}_2}}{-e} \quad \text{and} \quad (\text{CPD})_{\text{GaN}} = \frac{\phi_{\text{tip}} - \phi_{\text{GaN}}}{-e} \quad (2)$$

where ϕ_{tip} , ϕ_{GaN} , and ϕ_{MoS_2} are the work functions of the KPFM tip, GaN, and MoS₂, respectively. The CPD between MoS₂ and GaN is given by

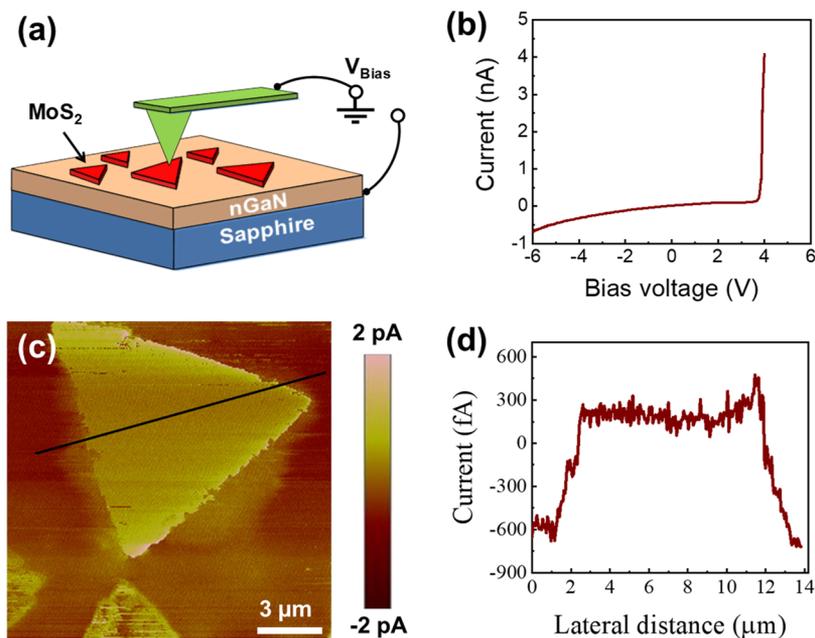


Figure 5. (a) Schematic of CAFM measurement of the MoS₂/GaN heterostructure. (b) *I*–*V* measurement of MoS₂/GaN by CAFM. (c) Current mapping of a typical MoS₂ domain on GaN. (d) Current difference with lateral distance across the MoS₂/GaN heterostructure corresponding to the dark line marked in (c).

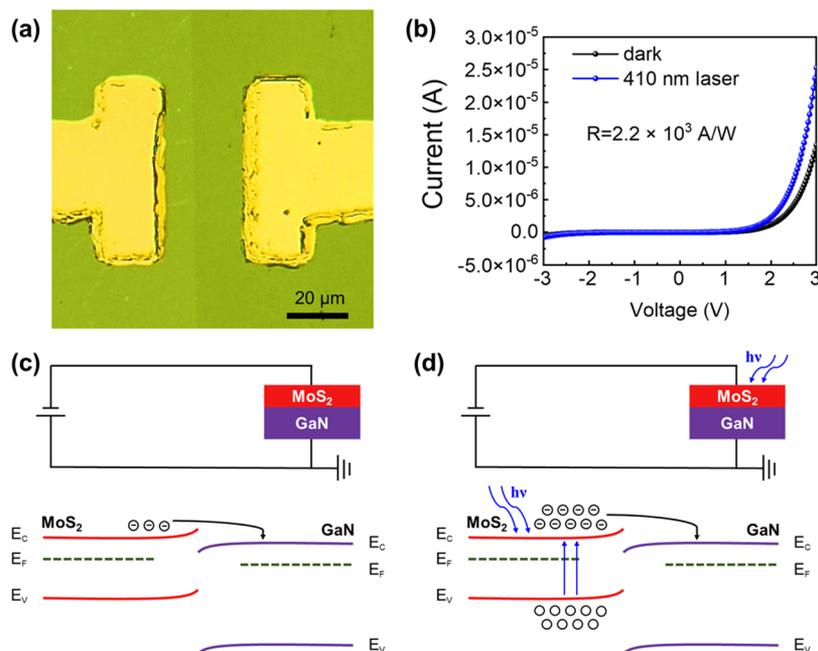


Figure 6. (a) Optical image of the continuous film of the MoS₂/GaN heterostructure device. (b) Current–voltage (*I*–*V*) curve of the MoS₂/GaN photodetector under dark conditions and after being illuminated by a 410 nm laser. (c, d) Band alignment of the MoS₂/GaN heterostructure before (c) and after illumination (d) under a positive bias voltage given to GaN with respect to MoS₂.

$$\Delta(\text{CPD}) = (\text{CPD})_{\text{GaN}} - (\text{CPD})_{\text{MoS}_2} = \frac{\phi_{\text{MoS}_2} - \phi_{\text{GaN}}}{-e} \quad (3)$$

It can be seen from Figure S3b that the surface potential of the GaN substrate is about 55 mV higher than that of the monolayer MoS₂, so the work function of the GaN substrate is 55 meV larger than that of the MoS₂ flake. Therefore, MoS₂ forms a heterojunction with the GaN substrate and has a built-

in potential difference, indicating that a significant charge transfer occurs at the interface between the two materials.²⁰

The charge transport at the vertical 2D/3D heterojunction interface is closely related to the properties of the 2D material. The current–voltage (*I*–*V*) relationship of the MoS₂/GaN heterojunction was carried out by CAFM, as illustrated by Figure 5a. The CAFM tip is contacted with the top surface of the MoS₂, and the other electrode is connected with blank GaN. The *I*–*V* measurement results in Figure 5b show that the current can vertically flow through the van der Waals gap at the

interface of the MoS₂/GaN heterojunction, resulting in obvious rectification characteristics, allowing the MoS₂/GaN heterostructure to be applied to the light-emitting diode. The current conduction through vertically stacked 2D/3D semiconductor heterojunctions brings the possibility of complex structures with specific electrical characteristics, such as utilizing MoS₂ as the epitaxial center for the growth of other 2D materials. It can be seen from Figure 5c,d that the electrical conductivity of the entire MoS₂ domain region in the MoS₂/GaN heterojunction is significantly larger than that of the blank GaN substrate, which is consistent with the results reported in the literature.⁹

Since MoS₂ can absorb visible light and generate photo-generated carriers, this 2D/3D heterojunction can be further applied to photodetectors. As shown in Figure 6a,b, the *I*–*V* curves present a rectifying behavior, which can be explained by the built-in potential between MoS₂ and GaN indicated by the KPFM results. Since the Fermi level of GaN is at a lower energy level than that of MoS₂, the built-in electric field is directed from MoS₂ to GaN. The band alignment under bias voltage was illustrated based on the KPFM analysis, as shown in Figure 6c.^{20,53,54} When a positive voltage is applied to GaN with respect to MoS₂, the built-in potential will be offset, leading to a decrease in the effective barrier and more electrons can transport from MoS₂ to GaN, resulting in a high value of the current. Similarly, a low value of the current is observed owing to the reinforcing of the built-in potential when reverse-biased. The photoresponse of the MoS₂/GaN heterostructure is illustrated by illuminating the device by a 410 nm laser as the energy is located between the bandgap of MoS₂ and GaN with a power density of 1 mW/cm²; the photoresponsivity (*R*) can illustrate the figure-of-merit of a photodetector by the following equation

$$R = \frac{I_{\text{illuminated}} - I_{\text{dark}}}{P_{\text{illuminated}}} \quad (4)$$

where *I*_{illuminated}, *I*_{dark}, and *P*_{illuminated} are the current after illumination, dark current, and the illuminated power of laser light falling on the active area of the device, respectively. The current was enhanced upon irradiation, which can be due to the increase of photogenerated carriers transferred from MoS₂ to GaN under a bias voltage. The calculated *R* of the MoS₂/GaN heterostructure is 2.2 × 10³ A/W at the bias voltage of 3 V; this value is higher than that of photodetectors based on monolayer MoS₂ field-effect transistors (FETs),⁵⁵ which might have contributed to the traps induced by the interface of the heterostructure. When the photogenerated holes are captured by the traps, hindering the recombination of excitons, then electrons are extracted from the source to maintain the charge neutrality of the channel, resulting in an ultrahigh photoresponsivity. However, the value of photoresponsivity is a little lower than the multilayer MoS₂/GaN heterostructure due to the stronger absorption in multilayer MoS₂ samples.^{20,54} We further analyze the external quantum efficiency (EQE) of the MoS₂/GaN heterostructure-based photodetector, which refers to the photoconversion efficiency, and can be expressed by the following formula

$$\text{EQE} = \frac{h \times c \times R}{e \times \lambda} \quad (5)$$

where *h* is the Planck constant, *c* is the speed of light, *R* is the photoresponsivity as calculated based on eq 4, *e* is the charge

of the electron, and *λ* is the wavelength of the illumination laser.⁵⁴ The EQE is calculated to be 6.6 × 10⁵% at the bias voltage of 3 V. A comparison of the performance parameters of the MoS₂/GaN photodetector with previous reports is outlined in Table S1,^{20,53–56} which means that the MoS₂/GaN heterostructure shows potential application prospects in photodetectors.

CONCLUSIONS

In summary, we demonstrated the utilizing of the PTAS seeding promoter-assisted CVD method to synthesize large-area monolayer MoS₂ single-crystal nanosheets with side length more than 100 μm on GaN substrates, which is 2 orders of magnitude larger than the reported epitaxially grown MoS₂ nanosheets on GaN substrates. The introduction of the PTAS seeds changed the surface energy of the GaN substrate and promoted the nucleation and rapid growth of MoS₂. A series of characterization results indicate that the seed-assisted CVD-grown large-area monolayer MoS₂ nanosheets on GaN have a good crystal quality. Their quality is comparable to the reported monolayer MoS₂ epitaxially grown on GaN. Our work provides a valuable guide for the large-area growth of 2D/3D heterostructures and their applications in LED, valley-spin electronic devices, and photodetectors, which are helpful for electrical designers and fundamental research in multifunctional devices.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsnm.1c02662>.

Circularly polarized PL characterizations of MoS₂ on GaN and on sapphire substrates, elemental analysis of the MoS₂/GaN heterostructure by XPS, surface potential distribution of MoS₂/GaN by KPFM, and photoresponse of the MoS₂/GaN heterostructure-based device (PDF)

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Author Contributions

C.C. and R.L. directed the research work. C.C. and P.Y. conceived and designed the experiments. P.Y. fabricated the MoS₂/GaN samples and performed the Raman and PL measurements. H.Y. conducted the ARPES test. P.Y., F.L., and X.G. fabricated the devices and performed the electrical performance measurements. P.Y. and Q.X. performed AFM, KPFM, and CAFM characterizations. C.C., P.Y., H.Y., H.W., and Z.-J.Q. analyzed the data. C.C., P.Y., and H.Y. cowrote the manuscript. All authors discussed the results and commented on the manuscript.

Notes

The authors declare no competing financial interest.

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