

# Laser-Modified Black Titanium Oxide Nanospheres and Their Photocatalytic Activities under Visible Light

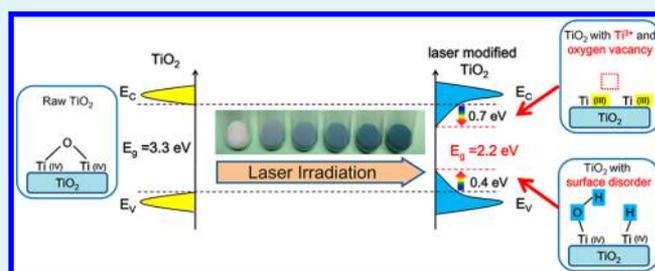
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## Supporting Information

**ABSTRACT:** A facile pulse laser ablation approach for preparing black titanium oxide nanospheres, which could be used as photocatalysts under visible light, is proposed. The black titanium oxide nanospheres are prepared by pulsed-laser irradiation of pure titanium oxide in suspended aqueous solution. The crystalline phases, morphology, and optical properties of the obtained nanospheres are characterized by means of X-ray diffraction (XRD), field-emission scanning electron microscopy (SEM), transmission electron microscopy (TEM), selected area electron diffraction (SAED), X-ray photoelectron spectroscopy (XPS), and UV–vis–NIR diffuse reflectance spectroscopy. It is shown that high-energy laser ablation of titanium oxide suspended solution benefited the formation of  $\text{Ti}^{3+}$  species and surface disorder on the surface of the titanium oxide nanospheres. The laser-modified black titanium oxide nanospheres could absorb the full spectrum of visible light, thus exhibiting good photocatalytic performance under visible light.

**KEYWORDS:** titanium oxides, laser, nanospheres, visible light, photocatalytic



## 1. INTRODUCTION

Titanium dioxide ( $\text{TiO}_2$ ) has attracted extensive interest because it was found to be an efficient photocatalyst by Honda and Fujishima.<sup>1</sup> At present,  $\text{TiO}_2$  has become the most promising candidate for photoelectrochemical water splitting,<sup>2–4</sup> organic pollutant degradation,<sup>5–7</sup> dye-sensitized solar cells,<sup>8,9</sup> and other photocatalytic applications<sup>10–14</sup> because of its low cost, good chemical stability, and nontoxicity.<sup>15</sup> However, pure  $\text{TiO}_2$  with its relatively large band gap (3.0 and 3.2 eV for rutile and anatase phases, respectively) can be activated only by ultraviolet (UV) light, which represents a small fraction of solar spectrum.

To overcome this drawback, pure  $\text{TiO}_2$  has been modified by impurity doping and dye sensitization. In impurity doping, the chemical composition of  $\text{TiO}_2$  is varied by adding transitional metals,<sup>16–18</sup> noble metals,<sup>19,20</sup> or nonmetals,<sup>21–24</sup> which could generate donor or acceptor states in the band gap. The doping of  $\text{TiO}_2$ , despite being the most common approach to improving visible light absorption characteristics of  $\text{TiO}_2$ , has many disadvantages. For example, these doped  $\text{TiO}_2$  catalysts have the drawbacks of secondary phase and the introduction of impurities strongly affecting the lifetime of the e–h pair.<sup>11</sup> In dye sensitization, organic molecules coated on the  $\text{TiO}_2$  surface were used to absorb visible light, and then an electron was injected from the excited state of the dye molecule into the conduction band of the semiconductor.<sup>25,26</sup> However, the dye degraded itself through the photocatalytic degradation of

organic pollutants. As an alternative to organic dyes, metallic nanostructures have been successfully used as photosensitizers,<sup>27</sup> but it was difficult to produce metallic nanostructure sensitizers with the advantages of finely tunable structure, mechanical durability, and low cost.

In recent years, an alternative approach has been developed to improve the visible and infrared optical absorption by engineering disorder states in  $\text{TiO}_2$ .<sup>28–35</sup> A porous  $\text{TiO}_2$  nanocrystal was prepared followed by the hydrogenation to create a disordered layer on the nanocrystal surface. The disorder-engineered  $\text{TiO}_2$  nanocrystals exhibited substantial solar-driven photocatalytic activities.<sup>28</sup> Several methods were developed to prepare hydrogenated black  $\text{TiO}_2$ , such as thermal treatments under a hydrogen atmosphere<sup>31–33</sup> and treatment by hydrogen plasma.<sup>34,35</sup> Fan et al. have developed a simple one-step combustion method to synthesize  $\text{Ti}^{3+}$ -doped  $\text{TiO}_2$ . The partially reduced  $\text{TiO}_2$ , which contains  $\text{Ti}^{3+}$  or oxygen vacancy, has been demonstrated to exhibit visible light absorption.<sup>29</sup>

Here, we proposed facile laser-modified  $\text{TiO}_2$  nanospheres that could be used as photocatalysts under visible light to degrade organic pollutants. The  $\text{TiO}_2$  nanoparticles were irradiated by a pulsed laser to form  $\text{TiO}_2$  nanosphere

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morphology. The  $\text{Ti}^{3+}$  species and disorder was introduced to the surface of the  $\text{TiO}_2$  nanospheres. The laser-modified  $\text{TiO}_2$  nanospheres could absorb visible and infrared light and had photocatalytic activities under visible light.

## 2. EXPERIMENTAL SECTION

**2.1. Preparation of  $\text{TiO}_2$  Nanospheres.** In a typical reaction, 20 mg of  $\text{TiO}_2$  (Aldrich, anatase, 100 nm) was added into 1 mL of distilled water followed by ultrasonic treatment for 10 min to form a suspension. The suspended solution was transferred into a cuvette. A Nd:YAG pulsed laser (Continuum PRII-8000,  $\lambda = 355$  nm, pulse duration = 8 ns, frequency = 10 Hz, power = 0.35 W, instantaneous power = 4.4 MW) was irradiated onto one side of the cuvette for several minutes, then the laser was irradiated onto the other side of the cuvette for the same time. The total duration of laser irradiation was  $n$  minutes. The laser beam gave a circular and homogeneous spot, which allowed a  $0.7 \text{ cm}^2$  area to be treated. The experimental setup is schematically shown in Figure 1.

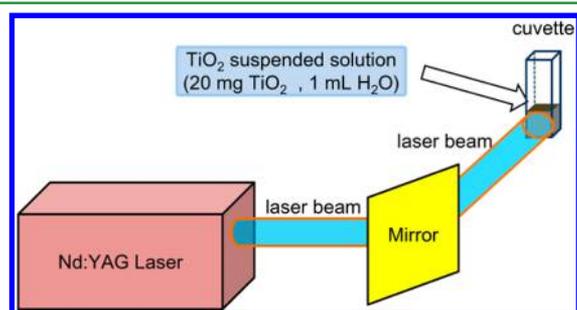


Figure 1. Schematic of the experimental setup.

After laser irradiation, the  $\text{TiO}_2$  was filtered and dried at  $80^\circ\text{C}$  for 12 h. The prepared samples were designated as  $\text{TiO}_2\text{-}n$  (the duration of laser irradiation was  $n$  minutes,  $n = 5, 15, 30, 60,$  and  $120$ ). All chemicals were analytical-grade and were used as received without further purification.

**2.2. Characterization.** To determine the crystal phase composition of samples, X-ray diffraction (XRD) measurements were carried out at room temperature using a Rigaku D/max-RA X-ray spectrometer with  $\text{Cu K}\alpha$  radiation ( $\lambda = 0.15418$  nm). An accelerating voltage of 40 kV and an emission current of 40 mA were used.

The morphology and structure properties of samples were investigated by field-emission scanning electron microscopy (FESEM, Hitachi S-4800). Samples for SEM observation were prepared by dispersing the prepared samples in alcohol, followed by ultrasonic treatment. A small portion of the suspended solution was dropped onto a smooth silicon slice. After being dried at  $110^\circ\text{C}$  for 4 h, the silicon slice was adhered on an aluminum mount by an adhesive conductive-carbon paper.

The microstructures of the samples were examined by transmission electron microscopy (TEM) on a JEOL JEM-2100F electron microscope. Samples for TEM observation were prepared by dispersing the prepared samples in alcohol, followed by ultrasonic treatment. A small portion of the suspended solution was dropped onto a copper grid, and then the copper grid was dried at  $60^\circ\text{C}$  for 12 h.

The core levels and valence band X-ray photoelectron spectra (XPS) were measured with an ESCALAB Mk II (Vacuum Generators) spectrometer using  $\text{Al K}\alpha$  X-rays (240 W). The binding energies were calibrated against the C 1s signal (284.6 eV) of adventitious carbon.

The UV-vis-NIR diffuse reflectance spectra were recorded on a Shimadzu UV-3101PC UV-vis-NIR spectrophotometer operating in the diffuse mode.

**2.3. Photocatalytic Reaction.** The photocatalytic activity of prepared  $\text{TiO}_2$  nanospheres was measured by monitoring the change in optical absorption of a rhodamine B solution at 553 nm during its photocatalytic decomposition process. In a typical experiment, 5 mg of

$\text{TiO}_2$  nanospheres was added to 15 mL of rhodamine B solution that had an optical absorption of approximately 2.0 at 553 nm. The suspended solution was magnetically stirred in the dark for 1 h and then exposed to green light from a green LED (voltage = 3.7 V, current = 25 mA, and center wavelength = 520 nm). We also tested the photocatalytic activity under irradiation of a red LED (voltage = 2.0 V, current = 45 mA, and center wavelength = 676 nm). The luminescent spectrum of the LED was shown in Figure S1 (Supporting Information). For a given time interval, the suspended solution was centrifuged, and then the filtrates were analyzed. The concentration of rhodamine B was monitored by the optical absorption peak by a Shimadzu UV3101PC UV-vis-NIR spectrophotometer.

## 3. RESULTS AND DISCUSSION

**3.1. Morphology and Structure Properties of  $\text{TiO}_2$  Nanospheres.** As shown in Figure 2, it could be seen that

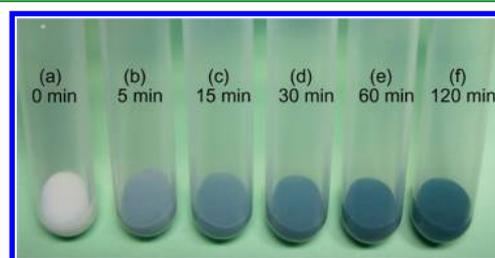


Figure 2. Suspended  $\text{TiO}_2$  solution before and after laser modification. (a) raw  $\text{TiO}_2$ , (b)  $\text{TiO}_2\text{-}5$ , (c)  $\text{TiO}_2\text{-}15$ , (d)  $\text{TiO}_2\text{-}30$ , (e)  $\text{TiO}_2\text{-}60$ , and (f)  $\text{TiO}_2\text{-}120$ .

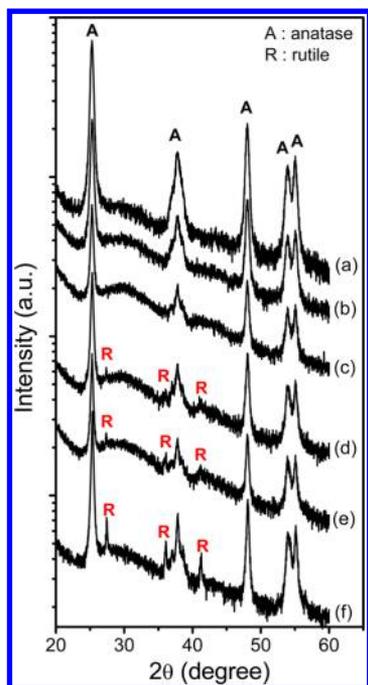
the color of the suspended  $\text{TiO}_2$  solution changed from white to gray after laser irradiation for only 5 min. As the duration of the laser modification increased, the color of the suspended  $\text{TiO}_2$  solution became deeper, turning totally black after laser irradiation for 120 min. After the laser irradiation, the  $\text{TiO}_2$  powder was filtered and dried. The photos of unmodified  $\text{TiO}_2$  and laser-modified black  $\text{TiO}_2$  powder are shown in Figure 3. The color of  $\text{TiO}_2$  powder changes from white to



Figure 3.  $\text{TiO}_2$  powder before and after laser modification. (a) Raw  $\text{TiO}_2$  and (b)  $\text{TiO}_2\text{-}120$ .

black after laser modification for 120 min, which suggests that the visible light could be totally absorbed by the laser-modified black  $\text{TiO}_2$ .

The structure of the  $\text{TiO}_2$  powders before and after the laser modification were explored by XRD. XRD revealed the structure of the mixture of the laser-induced  $\text{TiO}_2$  nanospheres and unreacted raw  $\text{TiO}_2$ . The powder XRD patterns of  $\text{TiO}_2$  synthesized under different durations of laser irradiation are shown in Figure 4. For the sample before laser irradiation, the XRD pattern is composed of lines at  $25.3, 36.9, 37.8, 38.6, 48.1, 53.8,$  and  $55.2^\circ$ , corresponding to (101), (103), (004), (112), (200), (105), and (211) lattice plane reflections of the anatase  $\text{TiO}_2$  phase (Powder Diffraction File (PDF) no. 21-1272, International Centre for Diffraction Data (ICDD), 1990), respectively. There are no other diffraction peaks except the



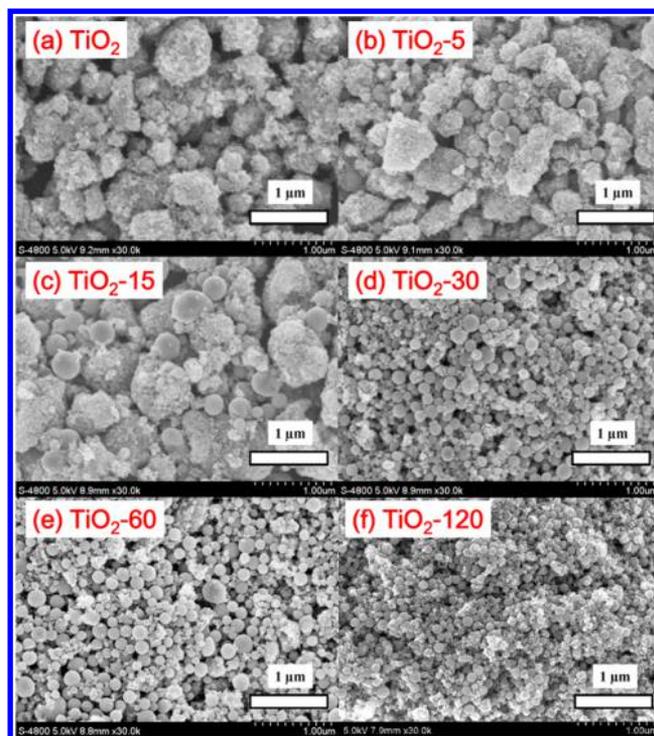
**Figure 4.** XRD patterns of  $\text{TiO}_2$  before and after laser modification ( $2\theta$  range from 20 to  $60^\circ$ ). (a) Raw  $\text{TiO}_2$ , (b)  $\text{TiO}_2$ -5, (c)  $\text{TiO}_2$ -15, (d)  $\text{TiO}_2$ -30, (e)  $\text{TiO}_2$ -60, and (f)  $\text{TiO}_2$ -120.

diffraction peaks of anatase  $\text{TiO}_2$  when the duration of laser irradiation was less than 15 min. When the laser irradiation duration lasts 30 min, three new diffraction peaks located at  $27.4$ ,  $36.1$ , and  $41.3^\circ$  appear, corresponding to (110), (101), and (111) lattice plane reflections of the rutile  $\text{TiO}_2$  phase (PDF no. 21-1276, ICDD, 1990). The diffraction peaks of rutile  $\text{TiO}_2$  phase increase as the duration of laser modification increases. The results of XRD suggest that the anatase phase  $\text{TiO}_2$  tends to transform to rutile phase under relatively long durations of laser modification because the rutile phase is more stable under high temperature and high pressure caused by pulsed-laser irradiation.

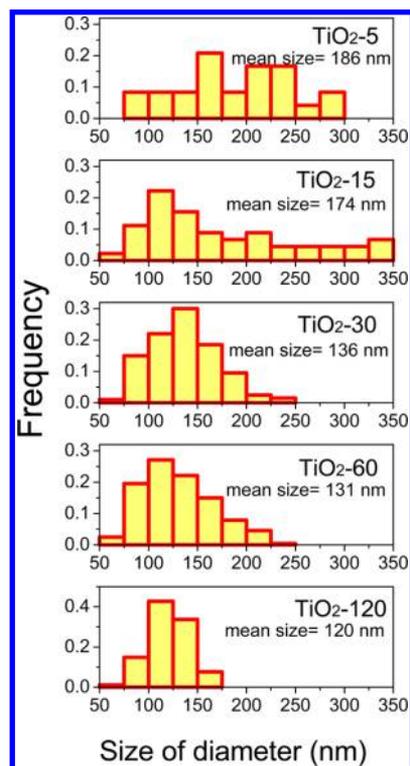
Figure 5 shows the SEM images of the  $\text{TiO}_2$  powders before and after the laser modification. It could be seen that  $\text{TiO}_2$  nanospheres appear after laser irradiation for only 5 min. By increasing the laser irradiated time, the amount of  $\text{TiO}_2$  nanospheres increases. A summary of statistics on the size of the obtained  $\text{TiO}_2$  nanospheres is shown in Figure 6. (The statistics on the size of the raw  $\text{TiO}_2$  are shown in Figure S2.)

The average size of  $\text{TiO}_2$  nanospheres is about 180 nm for the samples of  $\text{TiO}_2$ -5 and  $\text{TiO}_2$ -15. As the duration of laser irradiation increases, the average size decreases from 180 to approximately 130 nm. The size distribution of  $\text{TiO}_2$  nanospheres is relatively broad for the samples of  $\text{TiO}_2$ -5 and  $\text{TiO}_2$ -15, and the size distribution of  $\text{TiO}_2$  nanospheres becomes narrow for the samples of  $\text{TiO}_2$ -30,  $\text{TiO}_2$ -60, and  $\text{TiO}_2$ -120.

The results of SEM suggest that the  $\text{TiO}_2$  nanospheres are produced by the laser ablation of  $\text{TiO}_2$  nanoparticles. The initial  $\text{TiO}_2$  nanospheres have relatively large diameters and a broad size distribution. When irradiated by laser for a long duration, the initial  $\text{TiO}_2$  nanospheres continue to be ablated and to form relatively small  $\text{TiO}_2$  nanospheres, which have a narrow size distribution. Although the duration is prolonged from 30 to 120 min, the mean size of  $\text{TiO}_2$  nanospheres



**Figure 5.** SEM images of  $\text{TiO}_2$  before and after laser modification: (a) raw  $\text{TiO}_2$ , (b)  $\text{TiO}_2$ -5, (c)  $\text{TiO}_2$ -15, (d)  $\text{TiO}_2$ -30, (e)  $\text{TiO}_2$ -60, and (f)  $\text{TiO}_2$ -120.

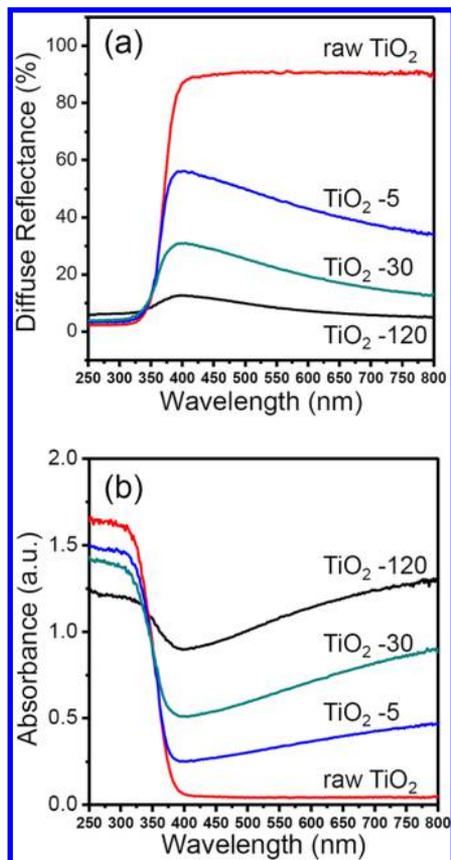


**Figure 6.** Size histogram of  $\text{TiO}_2$  nanospheres.

decreases only a little, which means the mean size around 120–130 nm is probably a stable condition under a relatively long duration of laser irradiation in our experimental conditions.

**3.2. Optical Properties of  $\text{TiO}_2$  Nanospheres.** To carry out a more complete characterization of light absorption,

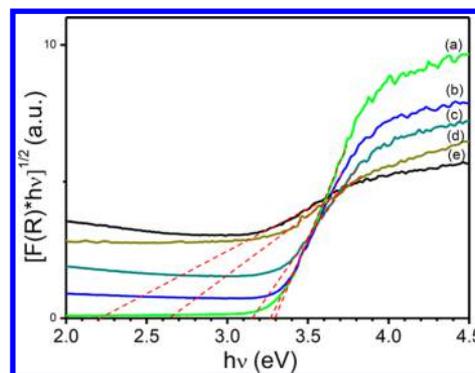
several UV–vis diffuse reflectance spectra of raw and laser irradiated TiO<sub>2</sub> were carried out. The results are shown in Figure 7. Figure 7a shows the diffuse reflectance data that is



**Figure 7.** UV–vis diffuse reflectance spectra of TiO<sub>2</sub> before and after laser modification: (a) diffuse reflectance data and (b) absorbance data ( $A = \log(1/R)$ ,  $A$  for absorbance and  $R$  for diffuse reflectance).

measured directly, and Figure 7b shows absorbance data that is calculated by the formula  $A = \log(1/R)$  ( $A$  for absorbance and  $R$  for diffuse reflectance). After laser modification, being compared with the raw white TiO<sub>2</sub> with an absorption edge of ca. 400 nm, the absorption of the laser-modified TiO<sub>2</sub> powders could extend to the full visible light spectrum and up to the near-infrared spectrum of 2000 nm. (UV–vis–NIR diffuse reflectance spectra are shown in Figure S3 in Supporting Information.) Meanwhile, the absorption above the wavelength of 400 nm increases as the duration of laser irradiation increasing.

The band gap values of samples are determined according to Kubelka–Munk function method as shown in Figure 8. The band gap of the laser irradiation TiO<sub>2</sub> shifts from 3.3 to 2.2 eV. Moreover, an enhanced absorption for photon energy below the 2.2 eV has been observed in our UV–visible spectra. It was reported that high concentration of oxygen vacancy could break the selection rule for indirect transitions of TiO<sub>2</sub> and enhanced absorption for photon energy below the direct band gap.<sup>29</sup> Furthermore, Ti<sup>3+</sup> and oxygen vacancy could be induced in TiO<sub>2</sub> by high-energy particle bombardment.<sup>36</sup> In our case, it is probably due to the laser irradiation that the high concentration of Ti<sup>3+</sup> and oxygen vacancy are generated on the surface of TiO<sub>2</sub> resulting in an enhanced absorption for photon energy

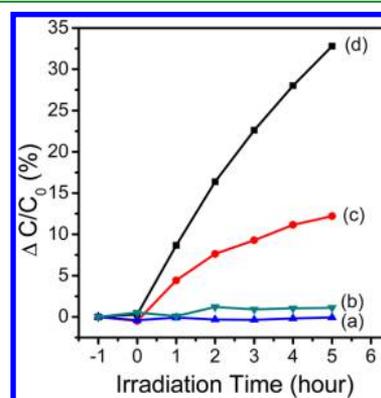


**Figure 8.** Plots of  $[F(R) \cdot hv]^{1/2}$  vs the energy of light for TiO<sub>2</sub> before and after laser modification: (a) raw TiO<sub>2</sub>, (b) TiO<sub>2</sub>-5, (c) TiO<sub>2</sub>-30, (d) TiO<sub>2</sub>-60, and (e) TiO<sub>2</sub>-120.  $F(R)$  is the Kubelka–Munk function, and  $F(R) = (1-R)^2/2R$ .

below the direct band gap. The exits of Ti<sup>3+</sup> and oxygen vacancy are confirmed by the characterization of XPS.

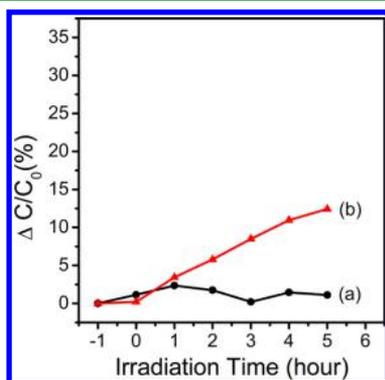
### 3.3. Photocatalytic Performance of TiO<sub>2</sub> Nanospheres under Visible Light.

The optical properties suggest that laser-modified TiO<sub>2</sub> nanospheres have the potential to act as a photocatalyst under visible light. Rhodamine B is a kind of organic dye, which is often used as a model pollutant to study the photocatalytic activity of photocatalysts. Here, we used rhodamine B as a model pollutant to study the photocatalytic activity of TiO<sub>2</sub> before and after laser modification. When the mixed solution of rhodamine B and TiO<sub>2</sub> was exposed to the green light irradiation from a LED at room temperature, it was found that the absorptive intensity of rhodamine B at 553 nm decreased as light irradiation time increased. We also found that the absorptive intensity of rhodamine B at 553 nm even decreased in the dark for 5 h. It was due to this that the adsorption equilibrium of rhodamine B with the catalyst cannot be completely reached in only 1 hour in a dark environment. To analyze the actual degradation ratio of rhodamine B, we calculated the  $\Delta C/C_0$  on the basis of the formula  $\Delta C/C_0 = (C/C_0)_{\text{dark}} - (C/C_0)_{\text{light}}$ . The results are shown in Figure 9. (The data of  $(C/C_0)_{\text{dark}}$  and  $(C/C_0)_{\text{light}}$  are shown in Figures S4 and S5 in the Supporting Information, respectively). The  $\Delta C/C_0$  is approximately constant at 0 when using P25 or pure TiO<sub>2</sub> as catalysts, indicating that P25 and pure TiO<sub>2</sub> cannot degraded rhodamine B under visible light. The  $\Delta C/C_0$  of laser-modified TiO<sub>2</sub> nanospheres increases as light irradiation time increases,



**Figure 9.** Photocatalytic degradation curves of rhodamine B under irradiation of green LED ( $\Delta C/C_0$  vs irradiation time): (a) P25, (b) raw TiO<sub>2</sub>, (c) TiO<sub>2</sub>-60, and (d) TiO<sub>2</sub>-120.

and the degradation ratio of rhodamine B could be up to 33% after irradiation with green light for 5 h when using TiO<sub>2</sub>-120 as catalyst. The sample of TiO<sub>2</sub>-120, which was irradiated by pulsed laser for 120 min, has strong absorption in the visible and infrared light region, which is in good agreement with the activity for photocatalytic degradation of organic pollutants in this region. Light absorption is a necessary condition for the photocatalyst to be functional. The photocatalytic activity of TiO<sub>2</sub>-120 is higher than that of TiO<sub>2</sub>-60 because of the higher visible light absorption ability of TiO<sub>2</sub>-120. We also tested the photocatalytic activity of the TiO<sub>2</sub>-120 under irradiation from a red LED. The results are shown in Figure 10. It is found that



**Figure 10.** Photocatalytic degradation curves of rhodamine B under irradiation of red LED ( $\Delta C/C_0$  vs irradiation time): (a) P25 and (b) TiO<sub>2</sub>-120.

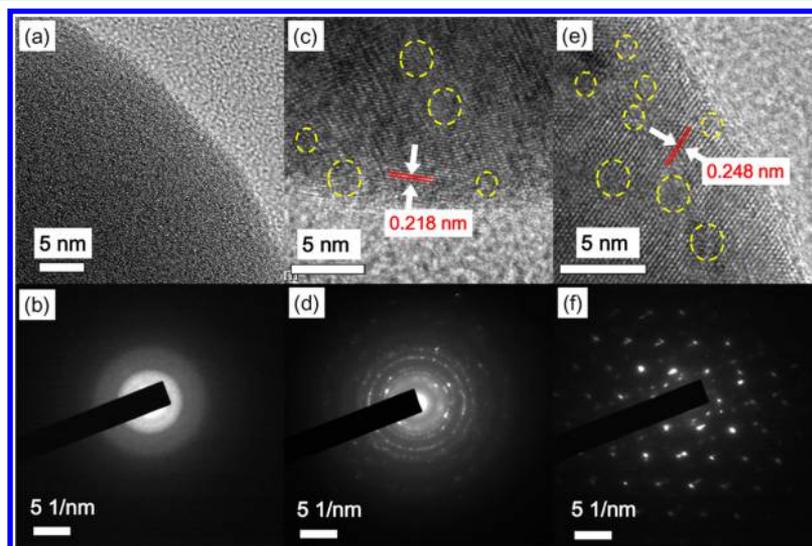
P25 cannot degrade rhodamine B under red light either. In contrast, the degradation ratio of rhodamine B could be up to 12% after irradiation by red light for 5 h when using TiO<sub>2</sub>-120 as catalyst, and the photocatalytic activity of TiO<sub>2</sub>-120 reveals a different intensity under green light and red light. This is probably due to the different light absorption ability. The energy of green light (center wavelength is 520 nm) is higher than the band gap of TiO<sub>2</sub>-120 (2.2 eV, corresponding to a light wavelength of 564 nm). Therefore, TiO<sub>2</sub>-120 could absorb green light intensely, and e<sup>-</sup>-h<sup>+</sup> pairs could be produced

effectively. For the red light (center wavelength is 676 nm), even though its energy is lower than the band gap, it could also be absorbed by TiO<sub>2</sub>-120, but it would be harder to produce e<sup>-</sup>-h<sup>+</sup> pairs effectively.

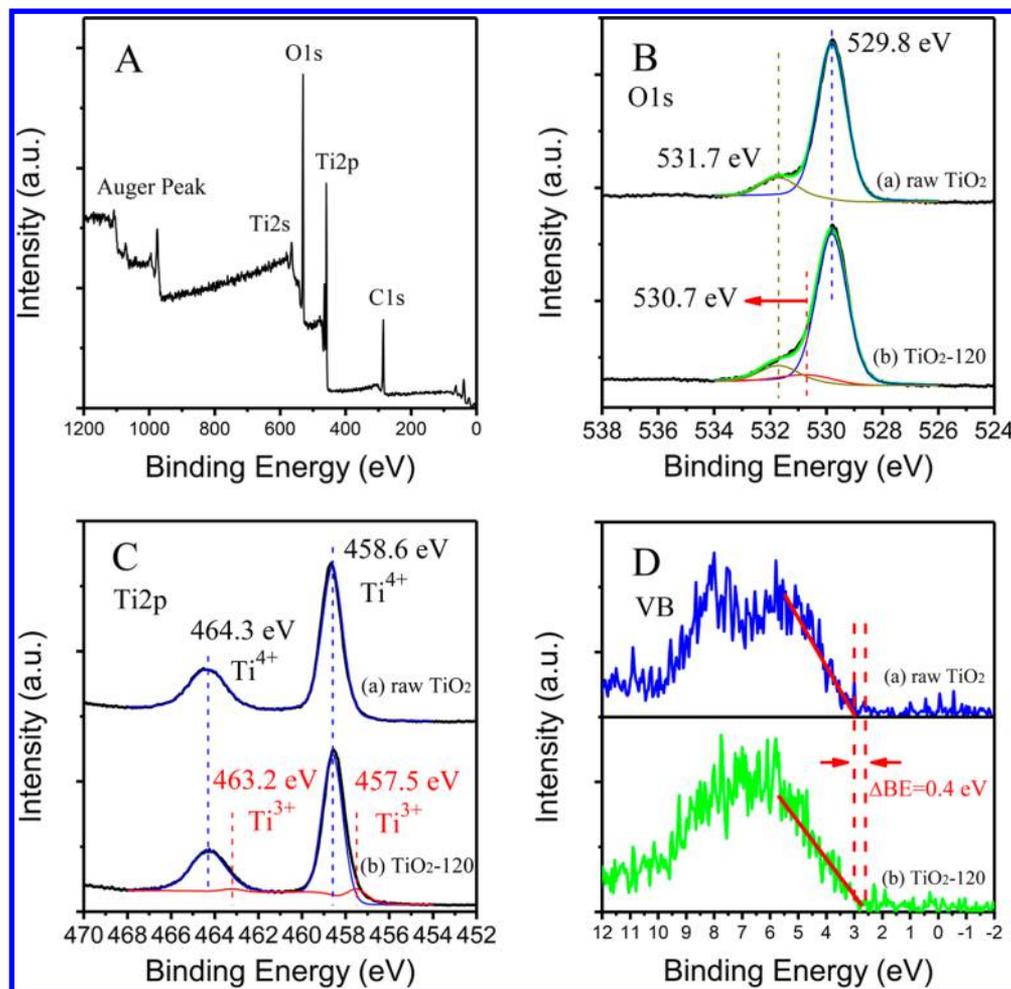
**3.4. Surface States of TiO<sub>2</sub> Nanospheres.** The TEM and selected-area electron-diffraction (SAED) patterns of laser-modified TiO<sub>2</sub> are shown in Figure 11. The TEM and SAED for TiO<sub>2</sub>-5 reveal an amorphous nature of particles. In the TEM image of TiO<sub>2</sub>-30, we could determine fringe separations of 0.218 nm, which corresponds to {111} rutile TiO<sub>2</sub> crystal planes. The SAED pattern of TiO<sub>2</sub>-30 is a group of concentric rings, indicating that the TiO<sub>2</sub> nanospheres were polycrystals. To investigate the crystalline phases of the TiO<sub>2</sub> nanospheres, a careful analysis of the SAED patterns was carried out (shown in Figure S6 in Supporting Information). Most interplanar distances could be assigned to various family planes of anatase and rutile phases. The distances of 0.351 and 0.162 nm are in agreement with {101} and {211} family planes of anatase TiO<sub>2</sub>, respectively. The distances of 0.310, 0.241, and 0.211 nm are in agreement with {110}, {101}, and {111} family planes of rutile TiO<sub>2</sub>, respectively. Thus, there are at least two crystalline phases coexisting in the ablated nanospheres. In the TEM image of TiO<sub>2</sub>-120, we could determine the fringe separations of 0.248 nm, which corresponds to {101} rutile TiO<sub>2</sub> crystal planes. The SAED for TiO<sub>2</sub>-120 reveals an almost single-crystal morphology.

We also find some defects and disorder in the TEM images of TiO<sub>2</sub>-30 and TiO<sub>2</sub>-120 that are outlined by dashed circles in Figure 11. The introduction of disorder at their surface would enhance visible and infrared absorption, with the additional benefit of carrier trapping.<sup>28</sup> This enhancement was probably the reason for the color change of the laser-modified TiO<sub>2</sub> nanospheres. Large amounts of lattice disorder in semiconductors could yield midgap states, which can form a continuum extending to and overlapping with the conduction band edge, known as band tail states. Similarly, large amounts of disorder can result in band tail states merging with the valence band.<sup>36,37</sup>

To investigate the change of the surface chemical bonding of TiO<sub>2</sub> nanocrystals induced by the laser irradiation, XPS analysis



**Figure 11.** TEM images and SAED patterns of laser-modified TiO<sub>2</sub>: (a and b) TiO<sub>2</sub>-5, (c and d) TiO<sub>2</sub>-30, and (e and f) TiO<sub>2</sub>-120. In c and e, dashed circles were applied to outline the disorder regions.



**Figure 12.** (A) XPS spectrum of  $\text{TiO}_2$ -120. (B) O 1s XPS spectra of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120. (C) Ti 2p XPS spectra of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120. (D) Valence band XPS spectra of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120.

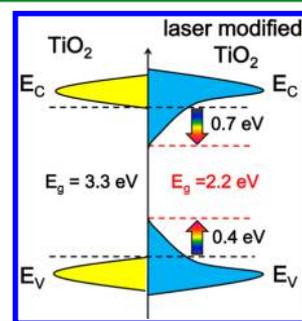
of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120 were carried out. On the basis of the survey XPS spectrum of  $\text{TiO}_2$ -120 (Figure 12A), the surface of  $\text{TiO}_2$ -120 is clean except for trace impurities of carbon. The O 1s spectra are shown in Figure 12B. The spectra are fitted with the nonlinear least-squares fit program using Gauss–Lorentz peak shapes. The spectra of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120 could be decomposed into a superposition of two and three peaks, respectively. For raw  $\text{TiO}_2$ , these peaks are attributed to lattice oxygen of  $\text{TiO}_2$  (529.8 eV) and adsorbed  $\text{O}_2$  (531.7 eV). In contrast, a new peak (530.7 eV) appears in the spectrum of  $\text{TiO}_2$ -120, which is attributed to Ti–OH species (surface disorder of  $\text{TiO}_2$ ). The details of the parameters of the O 1s peak are shown in Table S1 in the Supporting Information.

The XPS spectra of Ti 2p are shown in Figure 12C. The peaks located at 458.6 and 464.3 eV are assigned to  $\text{Ti}^{4+}$   $2p_{3/2}$  and  $\text{Ti}^{4+}$   $2p_{1/2}$ , respectively. For  $\text{TiO}_2$ -120, two new peaks appear. These peaks could be attributed to  $\text{Ti}^{3+}$   $2p_{3/2}$  (457.5 eV) and  $\text{Ti}^{3+}$   $2p_{1/2}$  (463.2 eV), respectively. The  $\text{Ti}^{3+}$  peak areas account for 12% of the total areas of Ti 2p, which means that 12% of Ti ions are  $\text{Ti}^{3+}$  ions. (The details of the parameters of the Ti 2p peak are shown in Table S2 in Supporting Information.) On the basis of the O 1s and Ti 2p XPS spectra, it is believed that Ti–OH and  $\text{Ti}^{3+}$  species are induced by the laser irradiation of  $\text{TiO}_2$ .

The density of states (DOS) of the valence band of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120 were also measured by valence band XPS, the

results which are shown in Figure 12D. The edges of the maximum energy of raw  $\text{TiO}_2$  and  $\text{TiO}_2$ -120 are 3.0 and 2.6 eV, respectively. It is suggested that a blueshift (0.4 eV) of the maximum energy is induced by the laser irradiation of  $\text{TiO}_2$ , and the valence band tail states extended above the valence band maximum for 0.4 eV. In consideration of the results of the UV–vis diffuse reflectance measurement, the conduction band tail states probably extended below the conduction band minimum for 0.7 eV.

The DOS of unmodified  $\text{TiO}_2$  and disorder-engineered  $\text{TiO}_2$  nanospheres is shown schematically in Figure 13. The change



**Figure 13.** Schematic illustration of electronic DOS of  $\text{TiO}_2$  before and after laser modification.

of DOS causes the enhancement of visible light absorption of laser-modified black TiO<sub>2</sub> nanospheres, which is probably the reason for the relatively good photocatalytic performance of laser-modified TiO<sub>2</sub> under visible light. The surface lattice disorder, including Ti–H and O–H bonds, could blue-shift the valence band maximum (VBM) of anatase TiO<sub>2</sub> by introducing midgap states while leaving its conduction band minimum (CBM) almost unchanged.<sup>38</sup> Furthermore, midgap states could be introduced closely below the CBM with the formation of Ti<sup>3+</sup> or oxygen vacancy.<sup>29</sup> In our case, band tail states were introduced around the VBM and CBM. The band tail states were probably due to the surface disorder and Ti<sup>3+</sup> induced by laser irradiation.

**3.5. Laser-Induced Reactive Quenching and Water-Mediated Growth of TiO<sub>2</sub> Nanospheres.** The formation process of TiO<sub>2</sub> nanoparticles by laser ablation of Ti plate target in a liquid medium has been intensively studied.<sup>39–42</sup> It was found that these TiO<sub>2</sub> had no change of the absorption of the visible light. It was probably due to the different pathway of the formation of TiO<sub>2</sub>. For the Ti plate target, according to previous studies,<sup>39</sup> the initial stage of interaction between pulsed-laser light and the target surface would generate a hot plasma plume over the laser spot on the metal plate, resulting in the ejection of titanium ions, atoms, and clusters from the melt surface. These species would strongly react with water molecules at the interfacial region. The chemical reaction products, such as titanium oxide or hydroxide molecules, would aggregate.

However, the formation of TiO<sub>2</sub> nanoparticles by laser ablation of a suspended TiO<sub>2</sub> solution should involve different processes. The TiO<sub>2</sub> nanoparticles were melting by the every single 8 ns high-energy pulse laser (instantaneous power = 4.4 MW). These initially melting TiO<sub>2</sub> clusters would then serve as nucleus and would rapidly be quenched in the liquid solution. These clusters must have a high fluidity so that they became spherical to reduce the surface energy. The crystalline structure of TiO<sub>2</sub> was broken by the laser beam at early stage. However, the anatase and rutile phase were reformed during the prolonged laser irradiation. This pathway probably benefited the formation of the disorder and the Ti<sup>3+</sup> species on the surface of the TiO<sub>2</sub>.

Ti<sup>3+</sup> and oxygen vacancy could be induced in TiO<sub>2</sub> by high-energy particle bombardment.<sup>36</sup> In our case, it was probably due to the laser irradiation that a high concentration of Ti<sup>3+</sup> and oxygen vacancy were induced in the surface of melting TiO<sub>2</sub> nanoparticles. The conduction and valence bands of TiO<sub>2</sub> were derived mainly from the Ti 3d and O 2p states, respectively. The high-energy laser beam ( $h\nu = 3.49$  eV) could make the electronic transition of the valence band electrons of the TiO<sub>2</sub>, which was composed of 2p electrons of O elements. On the surface of TiO<sub>2</sub>, O atoms donated the electrons to the Ti atoms. Therefore, Ti<sup>4+</sup> accepted the electron to form the Ti<sup>3+</sup>, and the O atoms could be peeled off from the surface of the TiO<sub>2</sub> to form the oxygen vacancies.

During the formation of the anatase and rutile phases, the oxygen vacancy and Ti<sup>3+</sup> were formed on the surface of the TiO<sub>2</sub> nanospheres. Because the laser irradiation was carried out in a suspended solution system, the Ti<sup>3+</sup> species tended to react with H<sub>2</sub>O, and surface disorder, including Ti–OH and Ti–H, were formed. The Ti<sup>3+</sup> species (oxygen vacancy) and surface disorder including Ti–OH were detected by XPS and TEM. These species could introduce band tail states in CBM and VBM and narrowed the band gap of TiO<sub>2</sub>. TiO<sub>2</sub> nanospheres

with absorption ability covering the full visible light spectrum were prepared by pulsed-laser ablation. The high-energy laser with a wavelength of 355 nm afforded a sufficient power to form a melting state and oxygen vacancies.

#### 4. CONCLUSIONS

We have demonstrated a new methodology for preparing black TiO<sub>2</sub> nanospheres, which could be used as photocatalysts under visible light effectively. Pulsed-laser ablation of pure TiO<sub>2</sub> suspended in aqueous solution was carried out, and the laser-induced local high temperature and high pressure led to the formation of TiO<sub>2</sub> nanospheres. The laser-modified TiO<sub>2</sub> nanospheres could absorb a full visible light spectrum, thus exhibiting good photocatalytic performance under visible light. We believe that this enhancement of photocatalytic performance is probably because of Ti<sup>3+</sup> and disorder on the surface of TiO<sub>2</sub> nanospheres which were introduced by pulsed-laser ablation.

#### ■ ASSOCIATED CONTENT

##### Supporting Information

Luminescent spectrum of LED. Size histogram of raw TiO<sub>2</sub>. UV–vis–NIR diffuse reflectance spectra of TiO<sub>2</sub>.  $(C/C_0)_{\text{dark}}$  versus time.  $(C/C_0)_{\text{light}}$  versus irradiation time. Analysis of the SAED patterns. Parameters of the O 1s peak of XPS. Parameters of the Ti 2p peak of XPS. The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsami.5b04568.

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##### Notes

The authors declare no competing financial interest.

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