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INFLUENCE OF HIGH PRESSURE ON THE EMISSION SPECTRA AND CRYSTAL FIELD PARAMETERS FOR TbP $_5$ 0 $_1$ 4

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The emission spectra and the energy levels of ${\rm Tb}^{3+}$ for ${\rm TbP}_5{\rm O}_{14}$ were measured at 81K up to 8 GPa. At various pressures between atm. and 8GPa, the crystal field parameters have been determined.

 ${
m TbP}_5{
m O}_{14}$ is an excellent fluorescent material, it has stimulated interest in its possible use in optical communication system. Recently, Tang Zhuoran has realized the fine structure of energy level of ${
m TbP}_5{
m O}_1$. The crystal field (c.f) parameters with ${
m C}_2{
m v}$ and ${
m C}_4{
m v}$ point site symmetry of ${
m Tb}$ in ${
m TbP}_5{
m O}_14$ at ambient pressure have been calculated respectively.

We measured the emission spectra of ${\rm TbP_5O_{14}}$ using a monochromator in the range of 12000- 20000 cm^{-1} at fixed temperature (81K) and various pressures up to 8GPa. A gasketed diamond anvil cell was used to produce high pressure, with 4:1 methanol-ethanol mixture for the hydrostatic pressure medium. Pressure was calibrated using the well-known Ruby fluorescence technique. The liquid nitrogen was spounted continously into the diamond cell to keep the constant temperature (81K). Thin plates of samples approximately $100\,\mu\text{m}$ in linear dimensions and $50\,\mu\text{m}$ thickness were used. The fluorescence was excited by the 4880 A° line of an argon laser. With the exception of the ${}^5\mathrm{D_4(A_1)}{}^-{}^7\mathrm{F_5(B_2)}$ transition, the observed 60 emission lines show red shift with increasing pressures. The energy levels corresponding to the spectra at various pressures are summarized in table 1. The assignment of the spectra was based on Tang's work[1]. From table 1 it can be seen that with increasing Pressure the most energy levels close to the

 $^7{\rm F}_0$, only a few levels are away from it. The shift rates of the $^7{\rm F}_{\rm J}$ are not larger than $10{\rm cm}^{-1}/{\rm GPa}$.

We have calculated the c.f. parameters for ${
m TbP}_5{
m O}_{14}$ from the measured energy levels. In this calculation we consider the luminescent center ${
m Tb}^{3+}$ occupying a site with distorted ${
m C}_{4v}$ symmetry ${
m [4]}$. The hamiltonian for equivalent 4f electrons of ${
m Tb}^{3+}$ can be written as a polynomical

$$H = H_0 + B_0^2 C_0^2 + B_0^4 C_0^4 + B_4^4 (C_{-4}^4 + C_4^4)$$
$$+ B_0^6 C_0^6 + B_4^6 (C_{-4}^6 + C_4^6)$$

where ${\rm H}_{\rm O}$ is the free ion Hamiltonian, the rest are the c.f. Hamiltonian, the tensor operators

$$C_q^k = \sum_i C_q^k (\theta_i, \varphi_i)$$

are related to the spherical harmonics Y_q^{k} by

$$C_{q}^{k} (\theta_{i}, \varphi_{i}) = (4\pi/(2k+1))^{\frac{1}{2}} Y_{q}^{k} (\theta_{i}, \varphi_{i})$$

Our calculating process and the definition of the rms deviations σ is similar to that described by reference [5]. At different pressures, the sets of the c.f. parameters and the rms deviations are listed in Table 2, the "free ion energies" of the crystal field energy levels in Table 3. From table 2 it can be seen that the absolute values of the c.f. parameters tend to increase with pressure. It implies that the c.f. tends to strenthen and therefore the split-

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Table 1 The observed energy levels (in cm $^{-1}$) at different pressure for ${\rm TbP_5O_{1\Delta}}$

at	GTITE	renr bres	Saic Loi	5 1	4	
		atm.	2GPa	4GPa	6GPa	8GPa
⁵ D ₄	E	14812	14789	14780	14770	14761
,	$^{A}_{1}$	14753	14748	14744	14743	14742
	A ₂	14741	14734	14729	14724	14719
	A 1	14722	14721	14713	14706	14698
	Ē	14708	14703	14697	14692	14686
	B_1	14698	14689	14684	14680	14675
	\mathbb{B}_{2}^{-}	14694	14682	14672	14667	14657
⁷ F ₀	A,	0	0	0	0	0
F,	Ε	-263	-255	-247	-237	-229
7 F $_{2}$	E	-573	-561	-550	-539	-528
	^B 1	-594	-589	-584	-577	-572
	A ₁	-739	-722	-705	-688	-671
	B_2	-842	-841	-840	-838	-831
⁷ F ₃	A ₂	-1364	-1364	-1364	-1360	-1360
J	Ē	-1401 ·	-1400	-1398	-1397	-1395
	E	-1425	-1425	-1425	-1423	-1423
	B_1	-1479	-1476	-1473	-1470	-1467
	B	-1500	-1506	-1511	-1515	-1521
⁷ F ₄	A ₁	-2221	-2218	-2214	-2208	-2202
	A ₂	-2272	-2267	-2262	-2255	-2251
	Ē	-2315	-2315	-2309	-2301	-2296
	$^{A}_1$	-2345	-2356	-2353	-2350	-2347
	$^{\mathrm{B}}_{2}$	-2388	-2376	-2366	-2354	-2344
	Ē	-2468	~2475	-2477	-2472	-2474
F_5	B_1	-3602	-3595	-3578	-3565	-3553
	E	-3614	-3602	-3591	-3576	-3564
	$^{\rm A}2$	-3673	-3664	-3656	-3648	-3640
	$^{A}_{1}$	-3698	-3699	-3701	-3704	-3706
	В2	-3723	-3729	-3735	-3738	-3745
	Ē	-3740	-3743	-3748	-3752	-3761

ting of each $^7\mathrm{F}_\mathrm{J}$ multiplet increases with pressure. The increasing of σ with pressure is pro-

bably caused by the increased deviation from ${\rm C_{4y}}$ site symmetry for the ${\rm Tb}^{3+}.$

Table 2. Crystal field parameters B_q^k and r_{ms} deviations (in cm $^{-1}$) at various pressures

	atm.	2GPa	4GPa	6GPa	8GPa
$_{0}^{2}$	-771	-816	-866	-986	-1062
B ₀ ⁴	-856	-861	-751	-702	-603
B ₄	578	598	628	664	685
B ₄ B ₆ B ₀ B ₄	843	902	1022	924	1031
в <mark>6</mark> В4	76	100	78	150	144
ø	17	19	21	22	24

Table 3. "Free ion energies" (in cm⁻¹) at different pressures

	atm.	2GPa	4GPa	6GPa	8GPa
7 F _O	-50	-56	-63	-79	-103
7 _F 0 7 _F 1 7 _{F2}	-242	-233	-224	-213	-205
$^{7}F_{2}$	-687	-678	-668	-658	-647
7 _{F3}	-1431	-1432	-1434	-1434	-1438
7 _{F4}	-2364	-2363	-2360	-2349	~2345
$7_{\text{F}_5}^{-7}$	-3636	-3630	-3625	-3617	-3612
7 _F 3 7 _F 3 7 _F 4 7 _F 5	-5704	-5699	-5697	-5677	-5671

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