

THE RECOMBINATION CROSS SECTION-ELECTRON ENERGY RELATION OF THE (e,A^o) EMISSION IN CdTe:Li

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By means of the effective mass approximation, the recombination cross section-electron energy relation of the free electron-acceptor emission in CdTe:Li has been studied. This relation has been compared with the relative theory.

1. INTRODUCTION

The photoluminescence investigation in CdTe crystal were made by several authors¹⁻³. In this paper, the recombination cross section (RCS) was introduced to character the probability. For the shallow acceptor A^o_{Li} (E_A=52meV), the effective mass approximation is valid. We have got RCS-electron energy relation by processing the photoluminescence line shape of the (e,A^o) emission in CdTe:Li.

2. EXPERIMENT

The impurity Li was diffused into CdTe crystal by two steps⁴. First, anneal of sample is made in a quartz tube at 350°C, under Te atmosphere for 320 hours. Second, Li diffusion is made from Li₂CO₃ solution deposited on a polish surface of CdTe crystal and in a horizontal furnace, under a continuous flow of H₂-N₂ mixture. The temperature is about 350°C and diffusion time is 1 hour.

For photoluminescence measurements were performed with Ar⁺ laser emitting about 100 mW power at 488 nm. The sample temperature was maintained at 11K by a cryogenic system.

3. RESULTS AND DISCUSSION

The energy position of the (e, A^o_{Li}) emission band is about at 12535 cm⁻¹. Both the (e,A^o_{Li}) emission band and the recombination band of the bound to the neutral acceptors (A^o_{Li},X) were well separated⁴. The (e, A^o) emission

line shape is given by the product⁵

$$I(\hbar\omega) \propto \rho(E_C) f(E_C) |M_{e,A^o}(k)|^2 \quad (1)$$

where $E_C = \hbar\omega - E_g + E_{A^o}$, $\rho(E_C)$ and $f(E_C)$ are the electron density of states and the electron distribution function respectively. The transition probability-energy relation has been considered in (1). But we can not obtain RCS-electron energy relation from (1) directly. However, we know that $I(\hbar\omega)$ are proportional to the recombination rate R . The recombination rate R in the (e,A^o) emission is given by⁶

$$R = \sigma(\hbar\omega) V n_e N_A^o \quad (2)$$

where $\sigma(\hbar\omega)$ is RCS, $V = (2E_C/m_c)^{1/2}$ is the electron velocity before recombination, $n_e = \rho(E_C) f(E_C)$ is the electron distribution in the conduction band, and N_A^o is the neutral acceptor concentration. RCS-electron energy relation, $\sigma(\hbar\omega) \propto E_C^{-1/2}$, has been determined by fitting (2) to the curve a which is the photoexcited emission spectrum of the sample in Fig. 1.

By the Ridley's theory, RCS of the (e,A^o) transition can be represented by⁷

$$\sigma(\hbar\omega) \propto \frac{P_{cv}^2}{[1 + (E_C/R_{HV})]^4 E_C^{1/2}} \quad (3)$$

where R_{HV} is the effective Rydberg energy, P_{cv} is the interband matrix element of the momentum operator. Using the formula (3) and the effective mass approximation, we have got the theoretical value of RCS (curve a in Fig. 2).

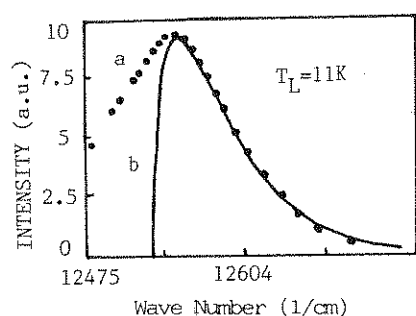


FIGURE 1

The point curve a is the (e, A^O) emission at 11K. The solid curve b has been obtained by fitting formula (2) to the spectrum data.

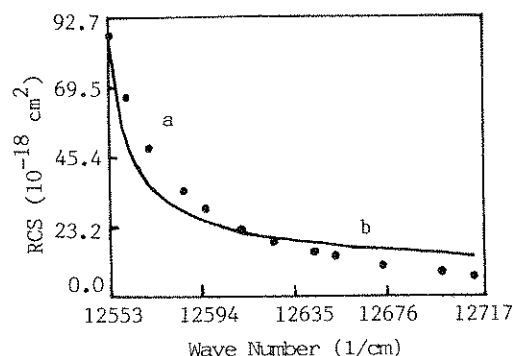


FIGURE 2

The curve b, the relation $\sigma(\hbar\omega) \propto E_C^{-1/2}$ has been compared with the theoretical value calculated by (3), curve a.

For the comparison $\sigma(\hbar\omega) \propto E_C^{-1/2}$ is also given as curve b in Fig. 2. Obviously, both curves a and b decrease with the increase of the electron energy. This means that the larger electron energy is, the smaller RCS is. The relation $\sigma(\hbar\omega) \propto E_C^{-1/2}$ is simpler for describing RCS than formula (3). The maximum of RCS in our experiment is $\sigma(\hbar\omega)_{\max} = 8.73 \times 10^{-17} \text{ cm}^2$ which is smaller than the geometric area of the acceptor A_{Li}^O , $\pi a^2 = 7.25 \times 10^{-16} \text{ cm}^2$.

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