

Section 7. Holeburning

**FLUORESCENCE LINE NARROWING AND INHOMOGENEOUS BROADENING OF  $\text{Sm}^{2+}$  IN  $\text{BaFCl}_x\text{Br}_{1-x}$**

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The variation of inhomogeneous and homogeneous line width with  $x$  for  $\text{Sm}^{2+}$  ions in  $\text{BaFCl}_x\text{Br}_{1-x}$  was studied by fluorescence line narrowing (FLN) and spectral hole-burning (HB). The ratios  $N$  of inhomogeneous width to homogeneous width approach a maximum for  $x = 0.8, 0.5$ .  $N$  is about 80 for transition from  ${}^7\text{F}_0$  to  ${}^5\text{D}_2$  of  $\text{BaFCl}_{0.5}\text{Br}_{0.5}$  at 77 K. The inhomogeneous lines for transitions between  ${}^5\text{D}_1$  and  ${}^7\text{F}_0$  of  $\text{Sm}^{2+}$  ions in  $\text{BaFCl}_x\text{Br}_{1-x}$  ( $x = 0.8, 0.5, 0.2$ ) have novel double-peaked shape. According to the FLN study, these novel shapes are explained by different probabilities of Cl and Br ligand configurations around  $\text{Sm}^{2+}$ . These novel shapes significantly broaden the inhomogeneous line width of  $\text{Sm}^{2+}$  and make hole burning available at 77 K. This might have an important implication on frequency domain optical storage.

1. Introduction

The first observation of photon gated hole burning in  $\text{BaFCl}:\text{Sm}^{2+}$  at 2 K was reported in 1985 by R.M. Macfarlane et al. [1]. In previous work [2], we have succeeded in raising the hole burning (HB) temperature to 77 K by adding Br ions to form  $\text{BaFCl}_{0.5}\text{Br}_{0.5}:\text{Sm}^{2+}$ , which might be important for the application of HB to frequency domain optical storage.

In this paper, we report results on  $\text{BaFCl}_x\text{Br}_{1-x}$ . The origin of a novel double-peaked inhomogeneous line shape of  $\text{BaFCl}_x\text{Br}_{1-x}$  ( $x = 0.8, 0.5, 0.2$ ) was clarified by fluorescence line narrowing (FLN) study. The variation of inhomogeneous line width ( $\Gamma_i$ ) and homogeneous line width ( $\Gamma_h$ ) with  $x$  was studied by FLN and HB.

2. Results

$\text{Sm}^{2+}$  ions have  $4f^6$  configuration. Fluorescence is due to transition from three metastable  ${}^5\text{D}_J$  ( $J = 0, 1, 2$ ) excited states to the  ${}^7\text{F}_J$  ( $J = 0, 1, 2, \dots, 6$ ) ground states.

The emission spectra of the  ${}^5\text{D}_0-{}^7\text{F}_0$  transition of  $\text{Sm}^{2+}$  in  $\text{BaFCl}$  under nonselective excitation at 337.1 nm of a nitrogen laser is peaked at  $14537.5 \text{ cm}^{-1}$  with a  $\Gamma_i$  of  $2.4 \text{ cm}^{-1}$ , while that of  $\text{BaFBr}:\text{Sm}^{2+}$  is peaked at  $14551 \text{ cm}^{-1}$  with a  $\Gamma_i$  of  $2.6 \text{ cm}^{-1}$ . However the emission lines of  $\text{BaFCl}_x\text{Br}_{1-x}:\text{Sm}^{2+}$  ( $x = 0.8, 0.5, 0.2$ ) have a novel double-peaked shape as shown in fig. 1. Their  $\Gamma_i$  are much broader than that for  $x = 1$  and  $x = 0$  ( $25 \text{ cm}^{-1}$  for  $x = 0.8$ ,  $26 \text{ cm}^{-1}$  for  $x = 0.5$ ). The line widths are almost the same at 77 K and 300 K. This

implies that the main contribution comes from the inhomogeneous broadening for  $x = 0.8, 0.5, 0.2$ .

FLN spectra of  ${}^5\text{D}_0-{}^7\text{F}_0$  of  $\text{BaFCl}_{0.8}\text{Br}_{0.2}:\text{Sm}^{2+}$  at 77 K excited at different frequencies within the inhomogeneous line width of the  ${}^7\text{F}_0-{}^5\text{D}_2$  transition, are shown in fig 2. The lines are much narrower and they shift with the excitation wavelength. The intensities of the FLN spectra ( $x = 0.8$ ) reach their maximum at  $14523 \text{ cm}^{-1}$  and  $14538 \text{ cm}^{-1}$ .  $\text{BaFCl}_{0.5}\text{Br}_{0.5}$  has similar FLN spectra. Based on the FLN data, we have plotted the inhomogeneous line shapes in fig. 1.

Homogeneous line widths of a transition from  ${}^5\text{D}_0, {}^5\text{D}_1$  to  ${}^7\text{F}_0$  were determined by the FLN method, exciting  ${}^5\text{D}_2$ . The homogeneous line width of a transition from  ${}^7\text{K}_0$  to  ${}^5\text{D}_2$  was determined by the HB method, exciting  ${}^5\text{D}_2$  and then monitoring emission  ${}^5\text{D}_0-{}^7\text{F}_0$  while scanning the persistent spectral hole in  ${}^5\text{D}_2$ . The results are shown in table 1. We can see that  $\Gamma_i$  and  $\Gamma_h$

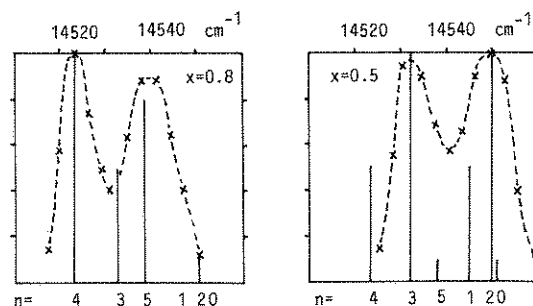


Fig. 1. Measured inhomogeneous line shapes of emission from  ${}^5\text{D}_0$  to  ${}^7\text{F}_0$  (broken line) and probabilities of ligand configuration  $\text{Cl}_n\text{Br}_{5-n}$  around the  $\text{Sm}^{2+}$  ion in  $\text{BaFCl}_x\text{Br}_{1-x}$  (vertical line).

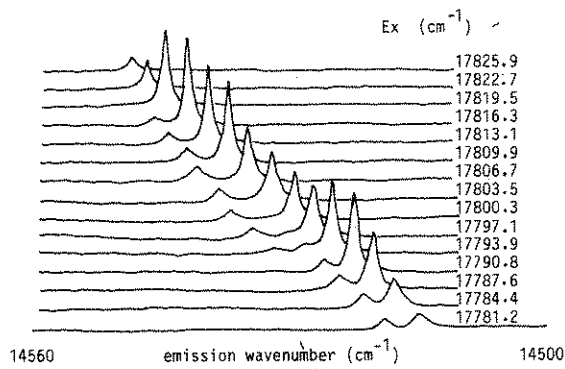


Fig. 2. FLN spectra of a transition from  ${}^5D_0$  to  ${}^7F_0$  of the  $\text{Sm}^{2+}$  ion in  $\text{BaFCl}_{0.8}\text{Br}_{0.2}$  at 77 K,  ${}^5D_2$  is excited.

reach their maximum at  $x = 0.5$ .  $N (= \Gamma_i/\Gamma_h)$  reaches its maximum at  $x = 0.8$  or  $0.5$ . The maximum  $N$  is as large as 80 for  ${}^5D_2$  with  $x = 0.5$ .

### 3. Discussion

In  $\text{BaFCl}:\text{Sm}^{2+}$ , the  $\text{Sm}^{2+}$  ions substitute at the  $\text{Ba}^{2+}$  ion sites. The coordination polyhedron of  $\text{Sm}^{2+}$  consists of four fluorine ions located in a plane perpendicular to the  $C$  axis, four chlorine ions in a plane parallel to the fluorine plane, and one extra chlorine ion on the  $C$  axis beyond the four chlorine ion planes. X-ray diffraction proves  $\text{BaFCl}_x\text{Br}_{1-x}:\text{Sm}^{2+}$  to be a solid solution with the same structure of  $\text{BaFCl}:\text{Sm}^{2+}$ . Chlorine and bromine ions occupy the chlorine ion places of  $\text{BaFCl}$  randomly. This will affect local crystal field strength and symmetry at the  $\text{Sm}^{2+}$  site. It is expected that an inhomogeneous line width of  $\text{BaFCl}_x\text{Br}_{1-x}:\text{Sm}^{2+}$  with  $x = 0.8, 0.5, 0.2$  is much broader than that for  $x = 1$  and  $0$ . However a simple statistical factor should make a single peaked Gaussian line shape. The double-peaked shape as shown in fig. 1

Table 1

Inhomogeneous line width ( $\Gamma_i$ ), homogeneous line width ( $\Gamma_h$ ) and ratio  $N (= \Gamma_i/\Gamma_h)$  of  $\text{Sm}^{2+}$  in  $\text{BaFCl}_x\text{Br}_{1-x}$  at 77 K

Transition	Parameter	$\text{BaFCl}_x\text{Br}_{1-x}:\text{Sm}^{2+}$ $x =$					Measured by
		1	0.8	0.5	0.2	0	
${}^5D_0-{}^7F_0$	$\Gamma_i$ ( $\text{cm}^{-1}$ )	2.4	25	26	10	2.6	FLN
	$\Gamma_h$ ( $\text{cm}^{-1}$ )	0.23	0.66	0.78	0.36	0.29	
	$N = \Gamma_i/\Gamma_h$	10	38	33	28	9	
${}^5D_0-{}^7F_1$	$\Gamma_i$	2.3	25	40	10	5	FLN
	$\Gamma_h$	0.42	0.78	1.3	1.0	0.36	
	$N$	5.5	38	30	10	14	
${}^7F_0-{}^5D_2$	$\Gamma_i$	3.3	35	40	16	5	HB
	$\Gamma_h$		0.48	0.5	0.35		
	$N$		73	80	46		

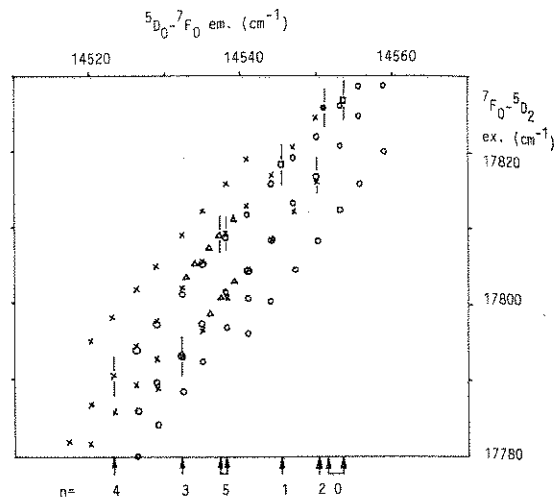


Fig. 3. Variations of the FLN peak of the transition from  ${}^5D_0$  to  ${}^7F_0$  with an excitation wavenumber of transition from  ${}^7F_0$  to  ${}^5D_2$  of  $\text{Sm}^{2+}$  in  $\text{BaFCl}_x\text{Br}_{1-x}$  at 77 K.  $\Delta$ ,  $x = 1$ ;  $\times$ ,  $x = 0.8$ ;  $\circ$ ,  $x = 0.5$ ;  $\square$ ,  $x = 0.2$ ;  $*$ ,  $x = 0$ .

is novel. The origin of this novel shape is discussed as follows.

The FLN data for different  $x$  are shown in fig. 3. The bars indicate the maximum positions of emission intensity when the excitation wavelength varies.

From fig. 3, we see that there are two excitation lines for each emission wavelength of  $\text{BaFCl}:\text{Sm}^{2+}$ . This is easily understood.  $\text{Sm}^{2+}$  ions in  $\text{BaFCl}$  replace the  $\text{Ba}^{2+}$  ion surrounded by five Cl ions and has  $C_{4v}$  symmetry. The  ${}^5D_2$  state splits into four sublevels in the  $C_{4v}$  crystal field, but only has two allowed transitions to the ground state  ${}^7F_0$ . The FLN spectra of  $\text{BaFCl}:\text{Sm}^{2+}$  have their maximum at  $14537.5 \text{ cm}^{-1}$ . Naturally,  $14537.5 \text{ cm}^{-1}$  is assigned to an  $\text{Sm}^{2+}$  ion surrounded by five chlorine ions, we indicate this by an arrow with  $n = 5$  in fig. 3.

For  $x \neq 1$  or  $0$ , the ligands Cl and Br around  $\text{Sm}^{2+}$  have different configurations with different probabili-

Table 2  
Probability of ligand  
 $\text{BaFCl}_x\text{Br}_{1-x}$

	$x = 1$
$\text{Cl}_5$	1
$\text{Cl}_4\text{Br}$	0
$\text{Cl}_3\text{Br}_2$	0
$\text{Cl}_2\text{Br}_3$	0
$\text{ClBr}_4$	0
$\text{Br}_5$	0

ties as shown in t most probable liga So we assign a  $\text{BaFCl}_{0.8}\text{Br}_{0.2}:\text{Sm}^{2+}$  transition  $\text{Cl}_5$ , since it  $\text{cm}^{-1}$  maximum of  $14523 \text{ cm}^{-1}$  is as  $\text{Cl}_4\text{Br}$ . For  $x = 0.5$  tions are  $\text{Cl}_3\text{Br}_2$   $14532 \text{ cm}^{-1}$  and  $\text{Cl}_2\text{Br}_3$  respec of  $\text{BaFBr}:\text{Sm}^{2+}$  i uration. Hence,  $\text{BaFCl}_{0.2}\text{Br}_{0.8}:\text{Sm}^{2+}$  maximum at  $1454$  assignments, the p urations  $\text{Cl}_n\text{Br}_{5-n}$ , lines with differer inhomogeneous sh different ligand co ligand configurati neous line shape, mogeneous shapes contributions from

From fig. 1, v

Table 2  
Probability of ligand configuration  $\text{Cl}_n\text{Br}_{5-n}$  around  $\text{Sm}^{2+}$  in  $\text{BaFCl}_x\text{Br}_{1-x}$

	$x = 1$	$x = 0.8$	$x = 0.5$
$\text{Cl}_5$	1	0.32768	0.03125
$\text{Cl}_4\text{Br}$	0	0.4096	0.15625
$\text{Cl}_3\text{Br}_2$	0	0.2048	0.3125
$\text{Cl}_2\text{Br}_3$	0	0.0512	0.3125
$\text{ClBr}_4$	0	0.0064	0.15625
$\text{Br}_5$	0	0.00032	0.03125

ties as shown in table 2.  $\text{BaFCl}_{0.8}\text{Br}_{0.2}:\text{Sm}^{2+}$  has the most probable ligand configuration of  $\text{Cl}_5$ ,  $\text{Cl}_4\text{Br}$  next. So we assign a maximum at  $14538\text{ cm}^{-1}$  for  $\text{BaFCl}_{0.8}\text{Br}_{0.2}:\text{Sm}^{2+}$  in fig. 3 to the ligand configuration  $\text{Cl}_5$ , since it nearly coincides with the  $14537.5\text{ cm}^{-1}$  maximum of  $\text{BaFCl}:\text{Sm}^{2+}$ . Another maximum at  $14523\text{ cm}^{-1}$  is assigned to the ligand configuration of  $\text{Cl}_4\text{Br}$ . For  $x = 0.5$ , the most probable ligand configurations are  $\text{Cl}_3\text{Br}_2$  and  $\text{Cl}_2\text{Br}_3$ , so the maximums at  $14532\text{ cm}^{-1}$  and  $14550\text{ cm}^{-1}$  are assigned to  $\text{Cl}_3\text{Br}_2$  and  $\text{Cl}_2\text{Br}_3$  respectively. The maximum at  $14551\text{ cm}^{-1}$  of  $\text{BaFBr}:\text{Sm}^{2+}$  is assigned to the  $\text{Br}_5$  ligand configuration. Hence, the maximum at  $14554\text{ cm}^{-1}$  of  $\text{BaFCl}_{0.2}\text{Br}_{0.8}:\text{Sm}^{2+}$  is assigned to  $\text{Br}_5$ , and another maximum at  $14545\text{ cm}^{-1}$  to  $\text{ClBr}_4$ . According to these assignments, the probabilities of different ligand configurations  $\text{Cl}_n\text{Br}_{5-n}$  are shown in fig. 1 using vertical lines with different heights. Comparing the measured inhomogeneous shape with envelope of probabilities of different ligand configurations and considering that each ligand configuration has its own Gaussian inhomogeneous line shape, we find that the double-peaked inhomogeneous shapes can be explained by envelopes over contributions from different ligand configurations.

From fig. 1, we estimate the inhomogeneous line

width of each ligand configuration to be about  $8\text{ cm}^{-1}$ , which are caused by strain in the solid solution  $\text{BaFCl}_x\text{Br}_{1-x}$ .

#### 4. Conclusion

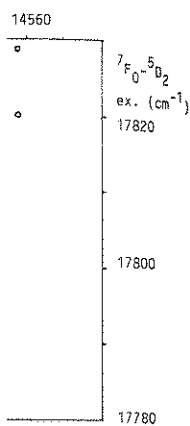
It is found that  $\text{BaFCl}_x\text{Br}_{1-x}:\text{Sm}^{2+}$  ( $x = 0.8, 0.5, 0.2$ ) have novel double-peaked inhomogeneous line shapes for transitions between  $^5\text{D}_j$  and  $^7\text{F}_0$  of  $\text{Sm}^{2+}$ . According to the FLN study, these novel shapes are explained by different probabilities of ligand configuration  $\text{Cl}_n\text{Br}_{5-n}$  around  $\text{Sm}^{2+}$ . These novel shapes significantly broaden the inhomogeneous line widths of  $\text{Sm}^{2+}$  and make hole burning available at 77 K. The variation of inhomogeneous and homogeneous line widths with  $x$  was studied. The ratio  $N$  of inhomogeneous width to homogeneous width approached a maximum at  $x = 0.8, 0.5$ .  $N$  is as large as 80 for the transition  $^7\text{F}_0-^5\text{D}_2$  of  $\text{BaFCl}_{0.5}\text{Br}_{0.5}$  at 77 K.

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#### References

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- [2] Changjiang Wei, Shihua Huang and Jiaqi Yu, *J. Lumin.* 43 (1989) 161.



transition from  $^5\text{D}_0$  transition from  $^7\text{F}_0$  to  $\Delta$ ,  $x = 1$ ;  $\times$ ,  $x = 0.8$ ;  $\circ$ ,  $x = 0$ .

shape is discussed as

are shown in fig. 3. Positions of emission length varies.

two excitation lines  $\text{BaFCl}:\text{Sm}^{2+}$ . This is  $\text{Cl}$  replace the  $\text{Ba}^{2+}$  has  $\text{C}_{4v}$  symmetry.  $\text{Cl}$  in the  $\text{C}_{4v}$  crystal positions to the ground  $\text{Cl}:\text{Sm}^{2+}$  have their  $14537.5\text{ cm}^{-1}$  is  $14532\text{ cm}^{-1}$  and  $14550\text{ cm}^{-1}$  are assigned to  $\text{Cl}_3\text{Br}_2$  and  $\text{Cl}_2\text{Br}_3$  respectively. The maximum at  $14551\text{ cm}^{-1}$  of  $\text{BaFBr}:\text{Sm}^{2+}$  is assigned to the  $\text{Br}_5$  ligand configuration. Hence, the maximum at  $14554\text{ cm}^{-1}$  of  $\text{BaFCl}_{0.2}\text{Br}_{0.8}:\text{Sm}^{2+}$  is assigned to  $\text{Br}_5$ , and another maximum at  $14545\text{ cm}^{-1}$  to  $\text{ClBr}_4$ . According to these assignments, the probabilities of different ligand configurations  $\text{Cl}_n\text{Br}_{5-n}$  are shown in fig. 1 using vertical lines with different heights. Comparing the measured inhomogeneous shape with envelope of probabilities of different ligand configurations and considering that each ligand configuration has its own Gaussian inhomogeneous line shape, we find that the double-peaked inhomogeneous shapes can be explained by envelopes over contributions from different ligand configurations.

From fig. 1, we estimate the inhomogeneous line

$\text{Cl}_x$  at 77 K.

Measured by
FLN
FLN
HB