



$$|x_{+n}(p, k)\rangle = -\sin\left(\frac{\theta_n}{2}\right)|g_{p-\hbar k/2}, n+1\rangle + \cos\left(\frac{\theta_n}{2}\right)|e_{p+\hbar k/2}, n\rangle, \quad (3)$$

where

$$\sin\theta_n = \frac{g(n+1)}{\sqrt{\Delta^2(p, k)/4 + g^2(n+1)^2}}. \quad (4)$$

It can be seen from Eq. (3) that the eigenfunction of Hamiltonian of the system is the linear combination of two different eigenstates of mass center momentum  $|g_{p-\hbar k/2}, n+1\rangle$  and  $|e_{p+\hbar k/2}, n\rangle$ . The Rabi vibration resulting from the interaction will inevitably cause the transition of momentum  $\hbar k$ . This is why this model can be used to explain the momentum and energy exchanges between the photon and atomic mass center motion.

It can also be seen from Eq. (3) that the spontaneous radiation probability from  $|e_p, 0\rangle$  to  $|g_{p-\hbar k}, 1\rangle$  is

$$P = P(k) = \frac{4g^2}{(\omega_a + pk/M - \omega)^2 + 4g^2} \times \cos\left[\left(\frac{\Delta^2(p, k)}{4} + 4g^2\right)^{1/2} t\right]. \quad (5)$$

This shows that the maximum value of spontaneous radiation probability  $P_{\max} = \cos(2gt)$  for momentum translation  $\hbar k$  where  $pk/M$  represents the atomic Doppler effect. Therefore, with Doppler effect taken into consideration, the condition for resonance is  $\omega = \omega_a + pk/M$ .

Let us deal with energy translation now. When  $t = 0$ , the atom in the ground state with momentum  $p$  enters the single-mode optical field in interference state  $|a\rangle$ . And the initial energy of the atom is  $E(0) = \varepsilon_p - \hbar\omega_a/2$ . The wave function for the optical-atomic system under the resonance condition is

$$|\psi(t)\rangle = e^{-i[E(p+\hbar k/2) - \hbar\omega/2]t/\hbar - |a|^2/2} \times \left\{ \sum_{n=0}^{\infty} \frac{\alpha^{n+1}}{\sqrt{(n+1)!}} \{ \cos[(n+1)gt] |g_p, n+1\rangle - i \sin[(n+1)gt] |e_{p+\hbar k}, n\rangle \} + |g_p, 0\rangle \right\}. \quad (6)$$

The exact value of atomic energy at moment  $t$  is

$$E(t) = \varepsilon_p + \frac{1}{2}\varepsilon_k + \frac{\hbar kp}{2M} - \frac{\hbar}{2}\left(\omega_a + \frac{kp}{M} + \frac{\varepsilon_k}{\eta}\right)$$

## References

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$$\times e^{-2|a|^2 \sin^2(gt)} \cos[(\sin 2gt)]. \quad (7)$$

The physical meaning of formula (7) is very obvious, that is, when the initial energy of the atom is  $E(0) = \varepsilon_p - \hbar\omega_a/2$ , the atom starts absorbing energy first after entering the cavity due to the photon recoil and Doppler effect, and the energy absorbed during  $\Delta t$  is

$$\Delta E = |a|^2 (g\Delta t)^2 \frac{\hbar}{2} \left( \omega_a + \frac{kp}{M} + \frac{\varepsilon_k}{\hbar} \right). \quad (8)$$

The release of the absorbed energy back to the system starts at  $t = \pi/2g$ , and the transfer of energy is an oscillation process with a period of  $\pi/g$ , as shown in Fig. 2 below.

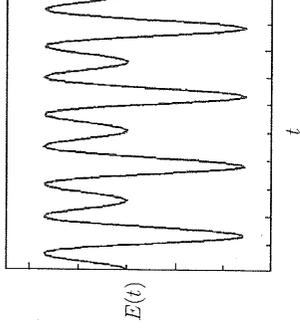


Fig. 2 Energy  $E(t)$  versus time.

The oscillation process with collapse and restore is right the characteristics of Jaynes-Cummings (JC) model, and this shows that the exchanges of energy and momentum under the resonance condition does not change this characteristics. However, the energy of atomic mass center motion must be taken into consideration in this case. On this occasion, the resonance condition can be expressed as follows:  $\omega = \omega_a + pk/2M + \hbar k^2/2M$ . The results indicate that the resonance condition must include the effect of Doppler effect and photon recoil.

It should be noted that a complete state with momentum  $p$  cannot be established through experiments. The momentum diffuse will damage the condition of resonance. However, we have presented in this paper a comparable result with the experiment at a certain approximation.

## References

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