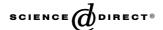


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# Thermal stability and spectral properties of Er<sup>3+</sup>-doped gadolinium borosilicate glasses

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## Abstract

A new glass of  $Gd_2O_3$ – $B_2O_3$ – $SiO_2$ – $Na_2O$ – $Er_2O_3$  system is prepared and the formation range of the glass is experimentally obtained. Excellent thermal stability of the glass has been demonstrated using differential thermal analysis. Effects of the  $B_2O_3$  content on the infrared spectral properties such as the 1.5  $\mu$ m emission bandwidth and the Judd–Ofelt parameter  $\Omega_t$  (t = 2, 4, 6) of  $Er^{3+}$  ions in the glass has been observed. The glass exhibits a large integrated emission cross sections for the 1.5  $\mu$ m emission band. The  $Er^{3+}$ -doped gadolinium borosilicate glass is a great candidate for broadband erbium doped fiber amplifiers.

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# 1. Introduction

In recent years, Er<sup>3+</sup>-doped glasses with a broad 1.5 μm emission band have been extensively investigated in searching erbium-doped fiber amplifiers (EDFA) with a wide and flat gain spectrum that is required for dense wavelength division multiplexing (DWDM) optical network systems in the telecommunication [1–7]. Among the reported Er<sup>3+</sup>-doped glasses, Er<sup>3+</sup>-doped tellurite glasses have shown wide 1.5 μm emission bandwidths and a large stimulated emission section. Mori reported excellent performance of a tellurite based EDFA, which gave 80 nm-wide gain around 1.5–1.6 μm [5]. However, tellurite glasses have relatively poor thermal stability in comparison with silicate glasses. Tanabe et al. re-

ported on bismuth-based glasses 43Bi<sub>2</sub>O<sub>3</sub>-xB<sub>2</sub>O<sub>3</sub>-(57 - x)SiO<sub>2</sub> with a bandwidth about 75 nm [7], but the description on thermal stability was unknown. In addition, thermal stability of Bi<sub>2</sub>O<sub>3</sub>-B<sub>2</sub>O<sub>3</sub>-Na<sub>2</sub>O glass reported by Yang [8], in which the glass transition temperature  $(T_g)$  and the temperature difference  $(\Delta T)$  only in excess of 415 and 123 °C, respectively. In this article, we report the preparation and the formation range of glass and the thermal stability of Er3+-doped gadolinium borosilicate glasses, Gd<sub>2</sub>O<sub>3</sub>–B<sub>2</sub>O<sub>3</sub>–SiO<sub>2</sub>–Na<sub>2</sub>O–Er<sub>2</sub>O<sub>3</sub> (GBSN:Er). A wide formation range of glass is experimentally obtained. Special attention was paid on  $10\text{Gd}_2\text{O}_3 - (80 - x)\text{B}_2\text{O}_3 - x\text{SiO}_2 - 10\text{Na}_2\text{O} - 0.5\text{Er}_2\text{O}_3$  $(G_{10}BSN:Er \text{ and } x = 0, 10, 20, 30, \text{ and } 40 \text{ mol}\%)$ , which shows high thermal stability. Spectral properties of the glass are also characterized. To our knowledge, there have been no research reports published yet concerning this glass system.

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# 2. Experimental

GBSN:Er glasses, in which the content of compositional Na<sub>2</sub>O is always 10 mol%, were prepared from anhydrous powders of  $Gd_2O_3$ ,  $B_2O_3$ ,  $SiO_2$ ,  $Na_2O$ , and  $Er_2O_3$  with more than 99.9% purity. After melting the powders in a platinum crucible at 1400–1450 °C for 1.5 h, they are poured into pre-heated brass molds. After the formation range of glass was confirmed, the  $G_{10}BSN$ :Er glasses were prepared. The quenched samples are annealed around 500 °C for 12 h and then cooled inside the furnace down to room temperature. The obtained  $G_{10}BSN$ :Er glasses have been named  $GS_0$ ,  $GS_{10}$ ,  $GS_{20}$ ,  $GS_{30}$ , and  $GS_{40}$  glasses, corresponding to the  $SiO_2$  content (x) of 0, 10, 20, 30, and 40 mol%, respectively.

Emission spectra were recorded using a TRIAX 550 spectrometer with a spectral resolution of 1 nm under excitation of 980 nm laser diode. Absorption spectra were obtained using a Perkin–Elmer 9 Lambda spectrometer. The differential thermal analysis (DTA) data was collected using a Perkin–Elmer 7 Series Thermal Analysis System.

### 3. Results

Fig. 1 shows the formation range of the GBSN:Er glass system. The symbols '●', '○', and '△' denote glass, non-glass and half glass state, respectively. The directions of arrowhead show the increase of the host content. The glass phase can be formed only when the content of SiO<sub>2</sub> is 0–50 mol%, Gd<sub>2</sub>O<sub>3</sub> is 0–30 mol% and B<sub>2</sub>O<sub>3</sub> is above 20 mol% in this glass system. The glass can also be obtained but becomes translucent at the contents of 60 mol% SiO<sub>2</sub> and 30 mol% Gd<sub>2</sub>O<sub>3</sub>, or at the contents of 60 mol% SiO<sub>2</sub> and 30 mol% B<sub>2</sub>O<sub>3</sub>. There is no glass phase formed in other glass components. In addition, Na<sub>2</sub>O is added into the material for the diminution of the melting temperature of the glass.

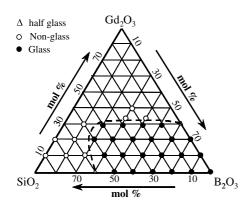


Fig. 1. Formation range of GBSN:Er glass system.

Fig. 2 shows the compositional dependence of glass transition temperature  $T_g$  ( $\blacksquare$ ), crystallization onset temperature  $T_x$  ( $\triangle$ ), and the temperature difference  $\Delta T = T_x - T_g$  ( $\blacksquare$ ) in the  $G_{10}BSN$ :Er glasses. It is found that both  $T_g$  and  $T_x$  increase as the SiO<sub>2</sub> content increases (up to x = 30) and the difference  $\Delta T$  reaches a maximum of 193 °C at x = 20. The values of  $T_g$  (>590 °C) and  $\Delta T$  (133–193 °C) in the  $G_{10}BSN$ :Er glasses are both higher than the corresponding temperatures in bismuth borate glass ( $T_g > 415$  °C; 170 °C >  $\Delta T > 123$  °C) [8] and tellurite glass ( $T_g > 290$  °C; 130 °C >  $\Delta T > 90$  °C) [9]. The  $G_{10}BSN$ :Er glass exhibits excellent thermal stability.

The spectral properties of glasses including 1.5  $\mu$ m bandwidth are directly related to Judd–Ofelt (J–O) parameters, which can be obtained by analyzing absorption spectra. Fig. 3 gives the absorption spectra of the  $G_{10}BSN$ :Er glasses recorded at room temperature. The spectra consist of seven absorption bands peaking at 1527, 975, 796, 652, 541, 520 and 488 nm, corresponding to the transitions from the ground state  $^4I_{15/2}$  to the

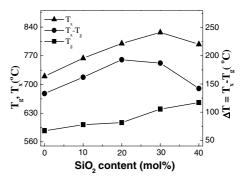


Fig. 2. Compositional dependence of glass transition temperature  $T_{\rm g}$  onset crystallization temperature  $T_{\rm x}$  and the difference  $T=T_{\rm x}-T_{\rm g}$  in the  $G_{10}$ BSN:Er glasses. The SiO<sub>2</sub> content 0, 10, 20, 30, and 40 mol% are responsible to the sample  $GS_0$ ,  $GS_{10}$ ,  $GS_{20}$ ,  $GS_{30}$ , and  $GS_{40}$ , respectively.

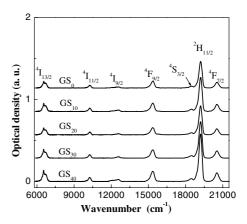


Fig. 3. The absorption spectra of the  $G_{10}BSN$ :Er glasses recorded at room temperature.

excited states  ${}^4I_{13/2}$ ,  ${}^4I_{11/2}$ ,  ${}^4I_{9/2}$ ,  ${}^4F_{9/2}$ ,  ${}^4S_{3/2}$ ,  ${}^2H_{11/2}$  and  ${}^4F_{7/2}$  of Er<sup>3+</sup> ion, respectively.

The radiative transition within the 4f'' configuration of  $Er^{3+}$  can be analyzed by use of the J–O theory [10,11]. The electric dipole transition line strength  $S_{\rm ed}$  between two J states can be expressed as [12]

$$S_{\text{ed}}(J;J') = \sum_{t=2,4,6} \Omega_t \langle \alpha SL, J \| U^{(t)} \| \alpha' S'L', J' \rangle^2$$
 (1)

where  $\Omega_t$  (t=2,4,6) is the Judd–Ofelt parameter,  $\alpha SL$  the other quantum numbers of the states, and the  $\langle \parallel U^{(t)} \parallel \rangle$  the doubly reduced matrix elements of unit tensor operators calculated in the intermediate coupling approximation. The  $\Omega_t$  parameters can be obtained from a least-squares fit of Eq. (1) to the line strengths of above-mentioned seven absorption bands. J–O parameters are the indicators of the local structure of the Er<sup>3+</sup>-doped host. The  $\Omega_2$  parameter is affected by the covalent chemical bonding and the  $\Omega_4$  and  $\Omega_6$  parameters are related to the rigidity of the medium where the Er<sup>3+</sup> ions are situated.

Fig. 4 depicts the relationship between the  $\Omega_t$  parameters of  $Er^{3+}$  and the  $B_2O_3$  content in the  $G_{10}BSN$ :Er glasses. It can be seen that  $\Omega_2$  decreases with the increase of  $B_2O_3$  content, while  $\Omega_6$  increases monotonically. There is a weaker degree of covalency of Er–O bond at the  $GS_0$  glass, because the  $\Omega_2$  reaches minimum value at the  $GS_0$  glass. In addition, comparing  $\Omega_4$  and  $\Omega_6$ , the rigidity and viscosity of the  $GS_{30}$  glass matrix are larger than that of all other glasses, because both the  $\Omega_4$  and  $\Omega_6$  reach minimum value at the  $GS_{30}$  glass.

Fig. 5 illustrates the emission spectra of the  $^4I_{13/2}$  to  $^4I_{15/2}$  transition of Er<sup>3+</sup> ions in the  $G_{10}BSN$ :Er glasses, where the area under the emission spectral profile is normalized. Because of the differences of the emission spectra in different glass hosts, the full width at half maximum (FWHM) is often used as a semiquantitative indication of the bandwidth. As shown in Fig. 5, the FWHM of the emission increases from 60 to 70 nm as the content of  $B_2O_3$  increases from 0 to 40 mol%.

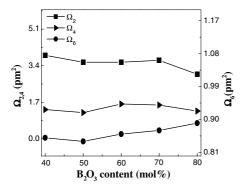


Fig. 4. The relationship between the  $\Omega_{t(2,4,6)}$  parameters of Er<sup>3+</sup> and the B<sub>2</sub>O<sub>3</sub> content in the G<sub>10</sub>BSN:Er glasses. The B<sub>2</sub>O<sub>3</sub> content 40, 50, 60, 70, and 80 mol% are responsible to the sample GS<sub>40</sub>, GS<sub>30</sub>, GS<sub>20</sub>, GS<sub>10</sub>, and GS<sub>0</sub>, respectively.

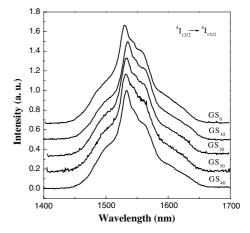


Fig. 5. The normalized  $1.5 \,\mu m$  emission spectra of  $Er^{3+}$  in the  $G_{10}BSN$ :Er glasses with different  $B_2O_3$  compositions.

Table 1 The values of the excitation energy E, the absorption and the calculated emission cross sections of  $\mathrm{Er}^{3+}$  at the peak place

Sample name	λ <sub>abp</sub> (nm)	λ <sub>emp</sub> (nm)	E (cm <sup>-1</sup> )	$\sigma_a^{\text{peak}} $ $(10^{-20} \text{ cm}^2)$	$\begin{array}{c} \sigma_{\rm e}^{\rm peak} \\ (10^{-20}~{\rm cm}^2) \end{array}$
$GS_0$	1527	1533	6531	0.3318	0.3469
$GS_{10}$	1527	1533	6556	0.3915	0.4592
$GS_{20}$	1526	1534	6526	0.3876	0.3994
$GS_{30}$	1529	1535	6546	0.4475	0.5188
$GS_{40}$	1528	1535	6551	0.3956	0.4665

The absorption cross section is obtained by using absorption spectra of the sample. According to the theory of McCumber [13], the stimulated emission cross section  $\sigma_e(\lambda)$  can be calculated using the following expression:

$$\sigma_{\rm e}(\lambda) = \sigma_{\rm a}(\lambda) \exp\left[\left(E - \frac{hc}{\lambda}\right) / kT\right]$$
 (2)

where  $\sigma_{\rm a}(\lambda)$  is the absorption cross section of the transition from the  $^4{\rm I}_{15/2}$  state to the  $^4{\rm I}_{13/2}$  state, E the temperature-dependent excitation energy, k the Boltzmann constant and T the temperature of the sample. The values of the excitation energy E, the absorption and the calculated emission cross sections of  ${\rm Er}^{3+}$  ions at the peak place are listed in Table 1. Both the absorption and emission cross sections of  ${\rm Er}^{3+}$  reach a maximum of  $0.45\times 10^{-20}$  and  $0.52\times 10^{-20}$  cm² at the GS<sub>30</sub> glass, respectively, as shown in Table 1.

# 4. Discussion

According to Eq. (1), the  $S_{\rm ed}$  of the  ${\rm Er}^{3+}$  ions is expressed as

$$S_{\rm ed}[^4I_{13/2}; ^4I_{15/2}] = 0.019 \times \Omega_2 + 0.118 \times \Omega_4 + 1.462 \times \Omega_6$$
(3)

Evidently, the  $\Omega_6$  plays a dominant role in broadening the bandwidth of the 1.5  $\mu$ m emission among the three  $\Omega_f$ s. Thus increasing the  $\Omega_6$  parameter will result in the increase of the 1.5  $\mu$ m emission bandwidth. Hence, both of the  $\Omega_6$  and the 1.5  $\mu$ m emission bandwidth increase with the increase of  $B_2O_3$  content in the  $G_{10}BSN$ :Er glasses, as shown in Fig. 6.

The gain bandwidth of optical amplifier can be evaluated by integrated emission cross section, estimated by the product of FWHM  $\times$   $\sigma_{\rm e}^{\rm peak}$  where  $\sigma_{\rm e}^{\rm peak}$  is the emission cross section at the peak of 1.5 µm band. The bigger the product is, the better the gain characteristic of EDFA yields. Fig. 7 represents the FWHM  $\times$   $\sigma_{\rm e}^{\rm peak}$  product of Er<sup>3+</sup> in different glass hosts. It can be seen that the value of the FWHM  $\times$   $\sigma_{\rm e}^{\rm peak}$  product in the prepared glasses are more than those of germanate [14], silicate [15] and phosphate [16] glasses. Furthermore, the maximum value of the product among these glasses reported in this work is smaller than that of oxyfluoride

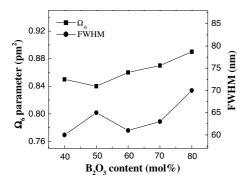


Fig. 6. The dependence of both of the  $\Omega_6$  value and the 1.5  $\mu m$  emission bandwidth on the content of  $B_2O_3$  in the  $G_{10}BSN$ :Er glass of  $Er^{3+}$  ions. The content of  $B_2O_3$  is 40, 50, 60, 70, and 80 mol% are responsible to the sample  $GS_{40}$ ,  $GS_{30}$ ,  $GS_{20}$ ,  $GS_{10}$ , and  $GS_0$ , respectively.

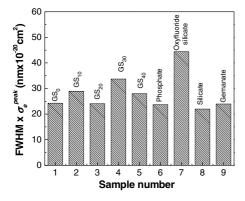


Fig. 7. The FWHM  $\times \sigma_e^{\text{peak}}$  product of Er<sup>3+</sup> in different glass hosts.

silicate glass [17], which have a bigger integrated emission cross section in the previous reported glasses.

### 5. Conclusions

A new glass of GBSN:Er glasses system is prepared, and formation range of the glass is experimentally obtained. On the analysis of DTA data, the  $\rm Er^{3^+}$ -doped  $\rm G_{10}BSN$ :Er glasses exhibit high thermal stability with glass transition temperatures higher than 590 °C. The glasses exhibit broad 1.5  $\mu m$  emission band around 70 nm, which are much larger than that of silica glasses with a narrow emission bandwidth around 30 nm presenting use. In addition, the glasses also have larger area of integrated emission cross sections than germanate, silicate and phosphate glasses. Therefore, the  $\rm Er^{3^+}$ -doped glass GBSN:Er system is a promising host for the EDFA in the 1.5  $\mu m$  window telecommunication system.

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