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### Point defects in active layers of TFEL devices based on ZnS

# LOU Zhidong<sup>1</sup>, A. N. Georgobiani<sup>2</sup>, XU Zheng<sup>3</sup>, XU Chunxiang<sup>4</sup>, TENG Feng<sup>4</sup>, YU Lei<sup>4</sup> and XU Xurong<sup>3,5</sup>

1. Applied Physics Department, Tianjin University, Tianjin 300072, China; 2. P.N. Lebedev Physical Institute, Russian Academy of Sciences, 117924 Moscow, Russian; 3. Physics Department, Northern Jiaotong University, Beijing 100044, China; 4. Institute of Material Physics, Tianjin Institute of Technology, Tianjin 300191, China; 5. Laboratory of Excited State Processes, Changchun Institute of Physics, Chinese Academy of Sciences, Changchun 130021, China

Abstract Point defects in the active layer of a layered optimization thin film electroluminescent device of ZnS:  $Er^{3+}$  were studied. The results indicate that besides  $Er^{3+}$  substituting for  $Zn^{2+}$  as luminescent centers, the dominant point defects are sulfur vacancies, zinc vacancies, shallow donors and deep acceptors. Their influence on electroluminescence is discussed.

#### Keywords: spectra of photoluminescence, point defects, thin film electroluminescence.

THE outstanding issue in the field of thin film electroluminescence (TFEL) is how to improve the blue brightness. Zinc sulfide (ZnS) is an important commercial TFEL material. Rare earth such as Er, Ce and Tb are usually doped into ZnS to obtain green or blue TFEL. The mechanism of electroluminescence indicates that the dominant factor affecting blue emission is hot electron energy. Xu et al. [1,2] put forward the structure of the layered optimization TFEL in order to solve this problem. According to his structure, the region of excitation of luminescent centers is separated from that of acceleration of hot electrons, and SiO<sub>2</sub> is used as the accelerating layer. It has been proved by some experiments that a large proportion of the intensity of the short wavelength emission can be achieved with this structure.

There exist many point defects in the active layers of TFEL devices due to some inevitable physical and chemical factors during thin film depositing, which influence the properties of electroluminescence. The dominant point defects in the active layers of TFEL devices based on ZnS are vacancies and impurities<sup>[3]</sup>. In the present note the influence of the point defects in ZnS: Er<sup>3+</sup> on TFEL is discussed.

#### 1 Sample preparation

There are two kinds of samples: powder ZnS: Er<sup>3+</sup> in the form of pellet and thin film ZnS: Er<sup>3+</sup>. A small amount of NH<sub>4</sub>Cl was doped into pure powder ZnS or powder ZnS with 0.5 mol % ErF<sub>3</sub> as flux. After the mixture was ground completely, pellets for deposition were formed by a compressor. Then the pellets were baked for 1.5 h at 1100 °C. Thin film samples were prepared with the electron beam evapora-

tion method.

#### 2 Experimental results and discussion

According to the method of sample preparation, there may exist the following point defects in ZnS:  $Er^{3+}$ : substitutional anions  $Er_{Zn}$ ,  $Cl^{-1}$  and  $F^{-1}$ , sulfur vacancies  $V_S$ , zinc vacancies  $V_{Zn}$  and a small amount of other uncontrollable impurities.

#### 2.1 Substitutional anion Erzn

Rare earth ions on zinc sites in ZnS normally act as luminescent centers, such as Er<sub>Zn</sub>. They are dis-

crete luminescent centers which have almost the same energy level structures as free ones. Fig. 1 is the electroluminescent spectrum of a ZnS:  $\rm Er^{3+}$  TFEL device arising from the dipole transitions between 4f energy levels. The main peaks are located at 528.5 nm and 549.5 nm corresponding to the  $^2H_{11/2} \rightarrow ^4I_{15/2}$  and  $^4S_{3/2} \rightarrow ^4I_{15/2}$  transitions, respectively. Usually the intensity of the former is stronger than the latter. There is also a red peak located at 665.0 nm corresponding to  $^4F_{9/2} \rightarrow ^4I_{15/2}$ . Besides, there are some weaker blue peaks.

We can draw a conclusion that Er<sup>3+</sup> ions are mainly introduced into ZnS lattice during evaporation deposition from the photoluminescent spectra of powder ZnS:Er<sup>3+</sup>, ZnS(Cl)

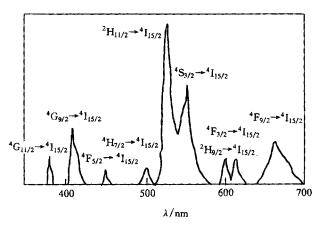


Fig. 1. EL spectrum of ZnS: Er3+.

and thin film ZnS:  $Er^{3+}$  at room temperature (figs. 2 and 3) under excitation of  $N_2$  laser (337.1 nm,  $3 \times 10^{-8}$  s). In powder ZnS;  $Er^{3+}$ , there only exists the same blue band centered at 449.0 nm as powder ZnS (Cl), but no emitting peaks of  $Er^{3+}$  exist. This blue band is a typical self-activated emitting band which is attributable to zinc vacancies<sup>[4]</sup>. We observed three sets of peaks of  $Er^{3+}$  in about 4- $\mu$ m-thick thin film ZnS:  $Er^{3+}$ . These peaks are added to a very wide background band which may be caused by dust. The first set of peaks due to the  ${}^4F_{7/2} \rightarrow {}^4I_{15/2}$  transition are located at 485.1, 491.6 and 494.8 nm, respectively. The second set due to  ${}^2H_{11/2} \rightarrow {}^4I_{15/2}$  are centered at 518.0, 522.4 and 527.3 nm, respectively. The third due to  ${}^4S_{3/2} \rightarrow {}^4I_{15/2}$  are located at 559.7, 564.5 and 567.8 nm, respectively. There are three peaks corresponding to every transition, which indicates that the energy levels of  $Er^{3+}$  are split by the crystal field. This phenomenon cannot be observed in electroluminescence. Perhaps it is the electric field

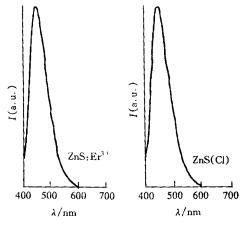


Fig. 2. PL spectra of powder ZnS: Er3+ and ZnS(Cl).

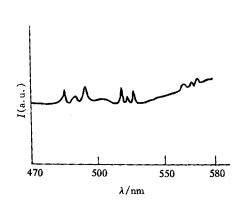


Fig. 3. PL spectrum of thin film ZnS: Er3+.

that offsets the influence of the crystal field.

There are two factors responsible for no light emitting of  $Er^{3+}$  in powder  $ZnS:Er^{3+}$ . First, there are few  $Er^{3+}$  ions that can enter powder ZnS lattice. Even if there are some, their light emission are very weak and covered by the strong recombination emission of ZnS. Second, the distribution of  $Er^{3+}$  ions in powder  $ZnS:Er^{3+}$  is not uniform since they are introduced into ZnS by the method of solid state mixing. It is possible that concentration quenching of  $Er^{3+}$  occurs, for some  $Er^{3+}$  clusters form on the surface of ZnS. However, the  $Er^{3+}$  clusters are broken up during the preparation of thin film  $ZnS:Er^{3+}$ , and the distribution of  $Er^{3+}$  ions becomes uniform. Therefore,  $Er^{3+}$  ions are introduced into the lattice during evaporation deposition.

#### 2.2 Vacancies

Figure 4 illustrates the photoluminescent spectrum of thin film ZnS(Cl) under N2 laser excitation. There is a blue band centered at 463.9 nm and a green shoulder centered at 524.6 nm at both room temperature and 77 K, but the intensities at 77K are higher. They are selfactivated recombination luminescence of ZnS attributable to the intrinsic defects in ZnS. Georgobiani<sup>[4]</sup> attributed the blue band to zinc vacancies, which comes from the transition of electrons from the conduction band to the singly-charged zinc vacancies  $V_{Zn}^{1-}$ . From the maximum position of the blue band in fig. 4, the optical activation energy of  $V_{Zn}^{\ 1-}$  was calculated to be at about 2.67 eV below the conduction band minimum, and 1.00 eV above the valence band maximum (the width of the forbidden band of ZnS is about 3.70 eV). The self-activated green shoulder relates to sulfur vacancies, which arises from the transition of electrons from the conduction band to the doubly-charged sulfur

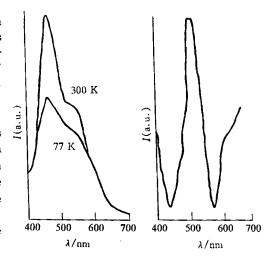


Fig. 4. PL spectra of Fig. 5. PL spectrum of thin film ZnS (Cl). Cu<sup>3+</sup> in powder Zn (Cl).

vacancies  $V_S^{2+}$ .  $V_S^{2+}$  was calculated to be at 2.36 eV below the conduction band minimum. Deng<sup>[5]</sup> indicated that there is a deep level at 2.14 eV below the conduction band minimum, which is very close to our estimation of  $V_S^{2+}$ . Besides, our estimation of  $V_S^{2+}$  approaches Pecheur's theoretical value<sup>[6]</sup> (at 2.48 eV below the conduction band) very much.

From fig. 3 we know that photoluminescence in thin film ZnS:  $Er^{3+}$  is mainly the emission of  $Er^{3+}$  ions. There are not any self-activated blue band and green shoulder of ZnS because the electron energy in the conduction band is transferred to the excited energy levels of  $Er^{3+}$ . However, there also exist S vacancies and Zn vacancies in thin film ZnS:  $Er^{3+}$  due to self-compensation.

(1) Sulfur vacancies  $V_S$ . A sulfur vacancy is a double donor<sup>[3]</sup> which has two levels in the forbidden band of ZnS. One is a singly-charged S vacancy  $V_S^{1+}$ , and the other is a doubly-charged S vacancy  $V_S^{2+}$ . Simple electrostatic repulsive energy considerations suggest that the separation between the first and second ionization of  $V_S$  is normally  $1.00~\text{eV}^{[7]}$ . We have already known that  $V_S^{2+}$  is at about 2.36~eV below the conduction band maximum, therefore  $V_S^{1+}$  is at about 1.36~eV below the conduction band maximum, which approaches Madelung's estimation of  $V_S^{1+}$  (at 1.4~eV below the conduction band maximum)<sup>[8]</sup> very much. There are two factors responsible for the generation of sulfur vacancies. First, in order to get thin film with good crystallinity, the substrate temperature is always kept at 250% or so. S ions would rapidly reevaporate from the hot substrate due to the high vapor pressure. Secondly, S vacancies may be considered as a consequence of the self-compensation of some uncontrollable shallow acceptors.

The Auger type-nonradiative energy transfer via sulfur vacancies  $^{[9,10]}$  occurs in ZnS: Ce<sup>3+</sup> and ZnS: Tb<sup>3+</sup> TFEL devices. This model is suitable to ZnS: Er<sup>3+</sup>, too. The relative intensities of Er<sup>3+</sup> peaks mainly depend on the impact excitation cross-sections, life times and state densities of Er<sup>3+</sup> energy levels,

and are also affected by the nonradiative energy transfer via sulfur vacancies. We have already known that the two levels  $V_S^{1+}$  and  $V_S^{2+}$  of S vacancies are located at 1.36 and 2.36 eV below the conduction band maximum.  $V_S^{1+}$  is known to have a strong Gaussian absorption F band centered at 2.27 eV and a weaker asymmetric absorption K band at 2.88 eV<sup>[11]</sup>.  $V_S^{2+}$  captures an electron from the conduction band or the surrounding  $Z_S^{1+}$  ions initially and becomes  $V_S^{1+}$ . A part of the dipole  ${}^2H_{11/2} \rightarrow {}^4I_{15/2}$  (528.5 nm, 2.35 eV) and  ${}^4S_{3/2} \rightarrow {}^4I_{15/2}$  (549.5 nm, 2.26 eV) transition energies is transferred to the absorption F band of  $V_S^{1+}$  centered at 2.27 eV. The probability of the energy transfer is proportional to the spectral overlapping between the emission and the absorption, and is in inverse proportion to the sixth power of the distance between  $E_S^{1+}$  and  $V_S$ . Therefore, the  ${}^4S_{3/2} \rightarrow {}^4I_{15/2}$  transition transfers more energy to the F band of  $V_S^{1+}$  than the  ${}^2H_{11/2} \rightarrow {}^4I_{15/2}$ . This may be one of the reasons why the intensity of the  ${}^4S_{3/2} \rightarrow {}^4I_{15/2}$  transition is weaker. The emitting intensities of the  ${}^4F_{5/2} \rightarrow {}^4I_{15/2}$  (458.0 nm, 2.71 eV) and  ${}^4F_{7/2} \rightarrow {}^4I_{15/2}$  (494.0 nm, 2.51 eV) transitions are lowered to a certain extent by the nonradiative energy transfer via  $V_S^{1+}$ .

So long as sulfur vacancies exist in the active layer of a ZnS TFEL device, it is inevitable that a part of the transition energy of short wavelength (blue or green) is lost. However, the luminance of short wavelength could be increased by restricting the amount of sulfur vacancies. Our study of  $ZnS_xO_{1-x}$ :  $Ce^{3+}$  TFEL devices<sup>[12]</sup> indicated that when oxygen substitutes for sulfur vacancies, the nonradiative energy transfer via  $V_S^{1+}$  is prevented, and the luminance of  $Ce^{3+}$  increases by about ten times.

(2) Zinc vacancies  $V_{Z_n}$ . A zinc vacancy is a double acceptor with two states in the forbidden band: singly-charged Zn vacancy  $V_{Z_n}^{1-}$  and doubly-charged Zn vacancy  $V_{Z_n}^{2-}$ .  $V_{Z_n}^{1-}$  in ZnS; Er<sup>3+</sup> thin film has been estimated to be at 2.67 eV below the conduction band maximum or 1.00 eV above the valence band minimum, which exactly equals the value of  $V_{Z_n}^{1-}$  Douglas<sup>[13]</sup> estimated.  $V_{Z_n}^{2-}$  was roughly investigated to be at 1.50 eV above the valence band maximum<sup>[7]</sup>. Zn vacancies caused by the self-compensation of Cl and F could influence thin film electroluminescence. Douglas *et al*. [13] and Shih *et al*. [14] attributed the origin of space charge to impact ionization of zinc vacancies. Losing an electron,  $V_{Z_n}^{2-}$  becomes  $V_{Z_n}^{1-}$ . There are two advantages of space charge generation: improving aging stability and reducing threshold voltage.

#### 3 Shallow donors and some other uncontrollable anion impurities

 $Cl^{-1}$  ions were introduced into ZnS because of flux NH<sub>4</sub>Cl. The ionic radii of  $Cl^{-1}$  and  $S^{2-}$  are 0.181 and 0.184 nm, respectively. The close size match of Cl and S suggests that Cl will invariably be incorporated as a substitutional donor  $Cl_S$ .  $Er^{3+}$  ions were introduced into ZnS in the form of a fluoride. The ionic radius of  $F^{-1}$  is 0.136 nm, much smaller than that of  $S^{2-}$ . Electro-negative ions  $F^{-1}$  also act as substitutional donors  $F_S$ . Shallow donors  $Cl_S$  and  $F_S$  with their level depths being several tenth eV are electron traps. Space charge can be generated by their field emission at electric fields.

Some anion impurities could be introduced into ZnS inadvertently. We observed the photoluminescence of Cu<sup>+</sup> in powder ZnS. Fig. 5 is its spectrum under xenon lamp (369 nm) as an excitation source. We found a green band centered at 510 nm which had afterglow after the xenon lamp was turned off, suggesting that this would be the luminescence of Cu<sup>+</sup>. Cu<sup>+</sup> substituting for Zn<sup>2+</sup> is a deep level at 1.26 eV above the valence band maximum. The influence of Cu<sup>+</sup> remains to be solved.

#### 4 Conclusions

In this note we studied the energy level positions in the forbidden band of the point defects in ZnS:  $Er^{3+}$  thin film with the photoluminescence measurements. The results indicate that besides  $Er^{3+}$  substituting for  $Zn^{2+}$  as luminescent centers, the dominant point defects are sulfur vacancies, zinc vacancies, shallow donors  $Cl_S$  and  $F_S$ , deep acceptors  $Cu^+$  and some other uncontrollable shallow acceptors. Zinc vacancies are a double acceptor and sulfur vacancies are a double donor. According to our investigations,  $V_S^{1+}$  and  $V_S^{2+}$  are at 1.36 and 2.36 below the conduction band minimum, respectively;  $V_{Zn}^{1-}$  and  $V_{Zn}^{2-}$  are at 1.00 eV and 1.50 eV above the valence band, respectively; and  $Cu^+$  is at 1.26 eV above the valence band.

In the meantime we discussed the influence of the above-mentioned point defects on electrolumi-

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nescence. The blue and green emission of  $\mathrm{Er}^{3+}$  is decreased due to the Auger type-nonradiative energy transfer via sulfur vacancies. Space charge generation occurs through the field emission of  $\mathrm{Cl}_S$  and the impact ionization of  $\mathrm{V}_{Z_n}^{2-}$ .

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## Identification of a DNase activated in *Xenopus* egg extracts undergoing apoptosis

## ZHU Shan, JIANG Zhengfan, ZHANG Bo, LU Zhigang and ZHAI Zhonghe\*

College of Life Sciences, Peking University, Beijing 100871, China. \* Corresponding author, E-mail address; swzb@ibmstone.pku.edu.cn.

Abstract A cell-free apoptosis system was established by adding dATP and cytochrome c to *Xenopus laevis* egg extracts S-150. Accompanied by an incubation process, an apoptosis-specific DNase was activated in egg extracts which depended on  $Mg^{2+}$  and inhibited by  $Zn^{2+}$ . Two nucleases existing in egg extracts were revealed by in-gel nuclease assay. Further experiments showed that 27 ku nuclease which was different from other  $Ca^{2+}/Mg^{2+}$ -dependent nucleases was a possible candidate involved in apoptosis.

Keywords: apoptosis, DNase, cell-free system, Xenopus egg extracts.

APOPTOSIS is a mode of cell death characterized by distinct morphological features, such as membrane blebbing, cytoplasmic and nuclear condensation, chromatin aggregation and formation of apoptotic bodies<sup>[1]</sup>. Multiple evidence shows that the activation of cystine protease interleukin  $1\beta$  converting-enzyme