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Infrared-to-ultraviolet up-conversion luminescence from AlF_3 : $0.2\%\text{Tm}^{3+}$, $10\%\text{Yb}^{3+}$ particles prepared by pulsed laser ablation

Guanshi Qin^{a,b}, Weiping Qin^{a,b,*}, Changfeng Wu^{a,b}, Shihua Huang^{a,b}, Dan Zhao^{a,b},
Jisen Zhang^{a,b}, Shaozhe Lu^{a,b}

^aLaboratory of excited states processes, Chinese Academy of Sciences, Changchun 130021, People's Republic of China

^bChangchun Institute of Optics, Fine Mechanics and Physics, Chinese Academy of Sciences, Changchun 130021, People's Republic of China

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Abstract

AlF_3 : $0.2\%\text{Tm}^{3+}$, $10\%\text{Yb}^{3+}$ particles were prepared by pulsed laser ablation (PLA). Under a 978 nm laser diode excitation, intense ultraviolet and weak blue up-conversion emissions were observed in the AlF_3 : $0.2\%\text{Tm}^{3+}$, $10\%\text{Yb}^{3+}$ particles. Intensity dependence of the up-conversion emissions on the pump power was measured. The results show that the population of the states $^1\text{I}_6$, $^1\text{D}_2$ and $^1\text{G}_4$ may come from a five-photon, four-photon and three-photon energy transfer up-conversion process.

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1. Introduction

Recently, short-wavelength solid-state lasers in the ultraviolet (UV) to green spectral range have attracted much attention due to a wide range of applications including high-density optical data storage, color display, and infrared sensor. The phenomenon of frequency up-conversion by using the intrinsic energy level matching of certain rare-earth (RE) ions, as one of the available approaches exploring short-wavelength solid-state lasers, has been investigated widely during the past two decades [1–5]. Excited state absorption (ESA) and energy transfer (ET) can be efficient up-conversion mechanisms in RE doped materials. Shihua Huang et al. [6] reported up-conversion in LaF_3 : Tm^{3+} excited at 647.1 nm by ESA. Co-doping of Yb^{3+} as sensitizer has yielded a substantial improvement on the up-conversion efficiency in Tm^{3+} , Pr^{3+} , Ho^{3+} and Er^{3+}

doped systems due to the efficient ET between the sensitizer and the pair or triads of RE ions [7–9]. In Yb^{3+} -sensitized RE doped materials, efficient infrared-to-UV up-conversion emissions at room temperature have been rarely investigated [9,10]. Studies on efficient UV luminescence materials are very necessary and valuable due to the need of developing short-wavelength solid-state lasers.

In this paper, we reported intense UV up-conversion properties of the AlF_3 : $0.2\%\text{Tm}^{3+}$, $10\%\text{Yb}^{3+}$ particles prepared by pulsed laser ablation (PLA). With the excitation of a 978 nm-laser-diode (LD), intense UV and weak blue up-conversion emissions were observed. The mechanism responsible for the UV emission was discussed.

2. Results and discussion

1064 nm from a pulsed Nd:YAG laser was used as the light source for PLA. The laser beam was focused onto the target (AlF_3 : $10\%\text{Yb}^{3+}$, $0.2\%\text{Tm}^{3+}$) prepared by solid state reaction (the raw materials are high purity (99.95%) AlF_3 , YbF_3 and TmF_3) at 1000 °C for 1 h dipped in a beaker filled

* Corresponding author. Address: Key Laboratory of excited states processes, Changchun Institute of Optics, Fine Mechanics and Physics, Chinese Academy of Sciences, Yanan Road 1, Jilin 130021, P. R. China Tel.: +86-4315937564; fax: +86 4315937614.

E-mail address: gsqin@yahoo.com (W. Qin).

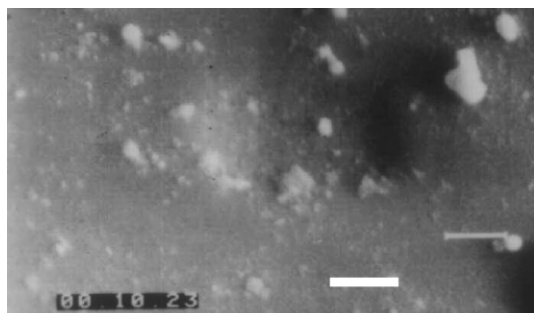


Fig. 1. The morphology of the particles measured with a transmission electron microscope.

with ethanol at an incidence angle of about 45° . The particles were collected in this beaker. The size of the particles is about $0.1 \sim 10 \mu\text{m}$ and the shape is not regular measured with a transmission electron microscope (JEM-2010), as shown in Fig. 1. X-ray analysis data show that the particles are polycrystalline and mixed phase of $\alpha\text{-AlF}_3$ and $\beta\text{-AlF}_3$. X-ray photoelectron spectroscopy (XPS) data show that the components of the particles are identical to that of the target. A 978 nm laser diode was focused on a pile of particles by a object lens ($\times 20$), the up-conversion emission spectra were measured with a fluorescence spectrophotometer (Hitachi F-4500) with the spectral resolution of 2 nm. The pump power of the 978 nm LD was numerate on the controller, which was used to modulate the electric current of the LD.

Fig. 2 is the room temperature up-conversion emission spectrum of the particles excited at 978 nm. Emissions in the UV and visible comes from the following transitions: $^1\text{I}_6 \rightarrow ^3\text{H}_6$ ($\sim 290 \text{ nm}$), $^1\text{I}_6 \rightarrow ^3\text{F}_4$ ($\sim 347 \text{ nm}$), $^1\text{D}_2 \rightarrow ^3\text{H}_6$ ($\sim 361 \text{ nm}$), $^1\text{D}_2 \rightarrow ^3\text{F}_4$ ($\sim 451 \text{ nm}$), $^1\text{G}_4 \rightarrow ^3\text{H}_6$ ($\sim 476 \text{ nm}$).

In $\text{Yb}^{3+}\text{-Tm}^{3+}$ co-doped systems, different process may result in up-conversion. Fig. 3 shows the energy level diagrams of Tm^{3+} and Yb^{3+} . The pump light excites only the Yb^{3+} ions, and three successive energy transfers from Yb^{3+} to Tm^{3+} populate $^3\text{H}_5$, ($^3\text{F}_3$, $^3\text{F}_2$), and $^1\text{G}_4$ [1]. Though

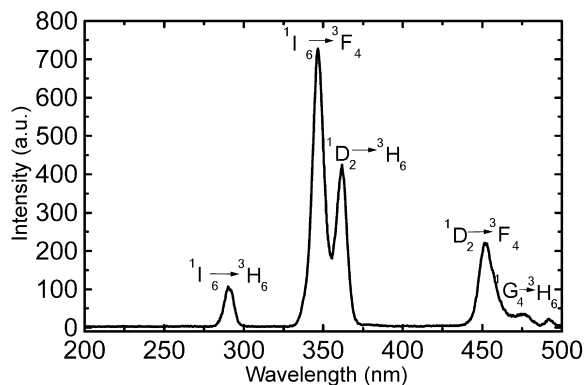


Fig. 2. Room temperature up-conversion emission spectra of AlF_3 : $10\%\text{Yb}^{3+}$, $0.2\%\text{Tm}^{3+}$ particles (at the pump power = 600 mW).

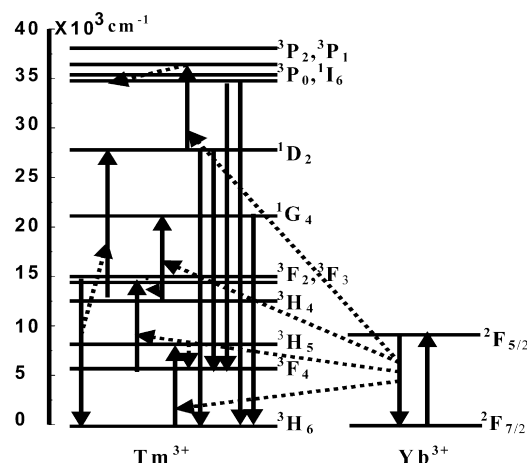


Fig. 3. Energy-level diagram of Tm^{3+} and Yb^{3+} .

the $\text{Tm}^{3+}\text{-Tm}^{3+}$ interaction is weak in the sample with low Tm^{3+} concentration, owing to the large energy mismatch ($\sim 3516 \text{ cm}^{-1}$) in the transfer $^2\text{F}_{5/2} \rightarrow ^2\text{F}_{7/2}$ (Yb^{3+}): $^1\text{G}_4 \rightarrow ^1\text{D}_2$ (Tm^{3+}), the process $^3\text{F}_2 \rightarrow ^3\text{H}_6$ (Tm^{3+}): $^3\text{H}_4 \rightarrow ^1\text{D}_2$ (Tm^{3+}) may alternatively play the most important role in populating $^1\text{D}_2$ [10]. Thereafter, the state $^1\text{I}_6$ can be populated by $^2\text{F}_{5/2} \rightarrow ^2\text{F}_{7/2}$ (Yb^{3+}): $^1\text{D}_2 \rightarrow ^1\text{I}_6$ (Tm^{3+}).

For unsaturated up-conversion, emission intensity, I_s , is proportional to I^n , where I is the intensity of the excitation light and the integer n is the number of photons absorbed per up-converted photon emitted [11]. In order to clarify the UV up-conversion mechanism, intensity dependence of the up-converted emissions are shown in Fig. 4 (the particles): at low pump power, $n = 4.91$ for the emission at 347 nm, $n = 4.02$ for the emission at 361 nm, $n = 4.18$ for the emission at 451 nm, and $n = 3.04$ for the emission at 476 nm. At high pump power, $n = 2.76$, 2.47, 2.65, 1.3 for the emissions at 347, 361, 451, 476 nm, respectively. The results show that the population of the states $^1\text{I}_6$, $^1\text{D}_2$ and $^1\text{G}_4$ come from a five-photon, four-photon and three-photon up-conversion process, which confirms that ET: $^3\text{F}_2 \rightarrow ^3\text{H}_6$

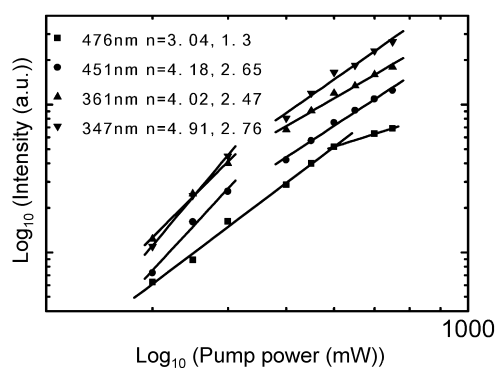


Fig. 4. Intensity dependence of the up-converted fluorescence on the pump power in AlF_3 : $10\%\text{Yb}^{3+}$, $0.2\%\text{Tm}^{3+}$ particles.

Table 1

The intensity ratio $\sigma = I_{347\text{ nm}}/I_{361\text{ nm}}$ at different pump power

The pump power (mW)	300	350	400	500	550	600	650	700	750
σ	0.89	0.97	1.12	1.19	1.31	1.38	1.37	1.44	1.48

(Tm^{3+}): $^3\text{H}_4 \rightarrow ^1\text{D}_2$ (Tm^{3+}) (four-photon up-conversion process) is the fundamental process to populate the state $^1\text{D}_2$. Thereafter, the state $^1\text{I}_6$ can be populated by $^2\text{F}_{5/2} \rightarrow ^2\text{F}_{7/2}$ (Yb^{3+}): $^1\text{D}_2 \rightarrow ^1\text{I}_6$ (Tm^{3+}) (five-photon up-conversion process).

The relative intensities of the UV up-conversion emissions versus excitation powers can be investigated simultaneously by recording the emission spectra at 300–400 nm at different excitation powers, as shown in Fig. 5. The results in Fig. 5 were listed in Table 1, which indicates that the intensity ratio σ for 347 nm emission from $^1\text{I}_6 \rightarrow ^3\text{F}_4$ transition to 361 nm from $^1\text{D}_2 \rightarrow ^3\text{H}_6$ transition is augmented with increasing the 978 nm-excitation power. Based on the UV up-conversion mechanism as shown in Fig. 3, we would use the rate equation to explain it.

Let γ represents the radiative transition rate from the $^1\text{D}_2$ level, N_0 and N the population of the states $^2\text{F}_{5/2}$ (Yb^{3+}) and $^1\text{D}_2$ (Tm^{3+}), X the rate of ET $^2\text{F}_{5/2} \rightarrow ^2\text{F}_{7/2}$ (Yb^{3+}): $^1\text{D}_2 \rightarrow ^1\text{I}_6$ (Tm^{3+}), β the branching ratio of the $^1\text{D}_2 \rightarrow ^3\text{H}_6$ transition, F the pump power and Ω the absorption cross-section of Yb^{3+} . At steady state, from the rate equation [6], we have

$$\sigma = \frac{NN_0XF\Omega}{\gamma\beta} \quad (1)$$

From the formula (1), the intensity ratio σ for 347 nm emission from $^1\text{I}_6 \rightarrow ^3\text{F}_4$ transition to 361 nm from $^1\text{D}_2 \rightarrow ^3\text{H}_6$ transition would be proportional to the pump power F . So increasing of F would make σ higher.

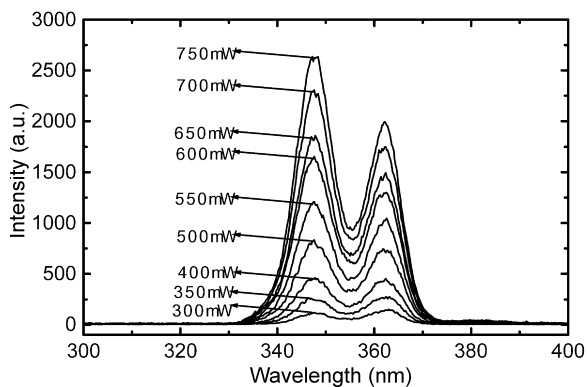


Fig. 5. Part of the emission spectra from $^1\text{I}_6 \rightarrow ^3\text{F}_4$ and $^1\text{D}_2 \rightarrow ^3\text{H}_6$ transitions measured at different excitation powers showing the emission intensity ratios between the two transitions with excitation at 978 nm.

3. Conclusion

In conclusion, $\text{AlF}_3: 0.2\%\text{Tm}^{3+}, 10\%\text{Yb}^{3+}$ particles were prepared by PLA. Under a 978 nm laser diode excitation, intense ultraviolet and weak blue up-conversion emissions were observed in the $\text{AlF}_3: 0.2\%\text{Tm}^{3+}, 10\%\text{Yb}^{3+}$ particles. Intensity dependence of the up-conversion emissions on the pump power that the population of the states $^1\text{I}_6$, $^1\text{D}_2$ and $^1\text{G}_4$ may come from a five-photon, four-photon and three-photon energy transfer up-conversion process.

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