



Spectral tuning and energy transfer in a potential fluorescent lamp phosphor $\text{BaMg}_2\text{Al}_6\text{Si}_9\text{O}_{30}:\text{Eu}^{2+}$

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ABSTRACT

$\text{BaMg}_2\text{Al}_6\text{Si}_9\text{O}_{30}:\text{Eu}^{2+}$ phosphors are synthesized by the solid-state reaction method and their photoluminescence (PL) properties are investigated. The ultraviolet emission originates from $\text{Eu}^{2+}(\text{I})$ substituting for Ba^{2+} sites, whereas the blue emission is attributed to $\text{Eu}^{2+}(\text{II})$ substituting for Mg^{2+} sites. With increasing Eu^{2+} doping concentrations, the blue emission band shifts to long wavelength and the PL intensity ratio of blue to ultraviolet emission increases. Energy transfer between the two different Eu^{2+} ions is analyzed by photoluminescence excitation and emission spectra, and lifetimes. Results indicate that the emission spectra can be tuned by changing Eu^{2+} contents. We have also demonstrated that $\text{BaMg}_2\text{Al}_6\text{Si}_9\text{O}_{30}:\text{Eu}^{2+}$ phosphor is a kind of potential phosphor for fluorescent lamps.

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1. Introduction

Fluorescent lamps (FLs), after a half century of their practical use, are still the major light sources because of their high efficiency and long life. Although the basics of commercial triphosphors—the blue-emitting $\text{BaMgAl}_{10}\text{O}_{17}:\text{Eu}^{2+}$ [1], the red-emitting $\text{Y}_2\text{O}_3:\text{Eu}^{3+}$ [2], and the green-emitting $\text{LaPO}_4:\text{Ce}^{3+}, \text{Tb}^{3+}$ [3] were well established, the development of new phosphors continues because of the importance of phosphor efficiency required for different applications as well as for the production cost. In recent years, researchers have extensively investigated various rare-earth ions doped luminescent materials for the excitation wavelength of 254 nm [4–8]. However, single-phased full-color emitting phosphors for FL application were rarely reported.

In our previous work, we synthesized a series of $\text{BaMg}_2\text{Al}_6\text{Si}_9\text{O}_{30}:\text{Eu}^{2+}$, Tb^{3+} , and Mn^{2+} phosphors suitable for UV excitations [9]. But we had never systematically investigated the Eu^{2+} contents effect on the emission spectra of $\text{BaMg}_2\text{Al}_6\text{Si}_9\text{O}_{30}:\text{Eu}^{2+}$. In this work, not only the luminescence properties but also the energy transfers in $\text{BaMg}_2\text{Al}_6\text{Si}_9\text{O}_{30}:\text{Eu}^{2+}$ with various Eu^{2+} concentrations, are studied. These results demonstrate that the emission spectra can be tuned with changing concentrations of Eu^{2+} ions. By using the optimized Eu^{2+} content, white light emitting

$\text{BMAS}:\text{Eu}^{2+}$, Tb^{3+} , and Mn^{2+} phosphors can be generated with higher color stability upon 254 nm excitation.

2. Experimental

The $\text{Ba}_{1-x}\text{Mg}_2\text{Al}_6\text{Si}_9\text{O}_{30}$ (BMAS): $x\text{Eu}^{2+}$ phosphors were synthesized by a high-temperature solid-state reaction. The constituent oxides or carbonates BaCO_3 (99.9%), Al_2O_3 (99.9%), SiO_2 (99.9%), MgO (99.9%), Eu_2O_3 (99.99%), and MnCO_3 (99.99%) were employed as the raw materials, which were mixed homogeneously by an agate mortar and pestled for 30 min, placed in a crucible with a lid, and then sintered in a tubular furnace at 1300 °C for 4 h in reductive atmosphere (10% H_2 +90% N_2 mixed flowing gas).

Powder X-ray diffraction (XRD) data were collected using Cu K α radiation ($\lambda=1.54056$ Å) on a Bruker D8 Advance diffractometer equipped with a linear position-sensitive detector (PSD-50m, M. Braun), operating at 40 kV and 40 mA with a step size of 0.02° (2 θ) in the range 10–70°. The measurements of photoluminescence (PL) and photoluminescence excitation (PLE) spectra were performed by using a Hitachi F4500 spectrometer equipped with a 150 W xenon lamp under a working voltage of 700 V. The excitation and emission slits were both set at 2.5 nm. In fluorescence lifetime measurements, the third harmonic (266 nm) of an Nd-doped yttrium aluminum garnet pulsed laser (Spectra-Physics, GCR 130) was used as the excitation source, and the signals were detected with a Tektronix digital oscilloscope (TDS 3052).

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3. Result and discussion

Fig. 1 shows the XRD patterns of BMAS:Eu²⁺ with various Eu²⁺ concentrations. The XRD patterns of all phosphors are similar to JCPDS 83-0740 [10] of the BMAS host structure. Results indicate that doped Eu²⁺ ions do not cause any significant change in the host structure.

Fig. 2(a) shows the PL spectra ($\lambda_{\text{ex}}=330$ nm) of Eu²⁺-activated BMAS as a function of Eu²⁺ concentration ranging from 0.005 to 0.1 mol. The PL spectra of BMAS consist of two distinct bands located at 376 and 450 nm under the excitation wavelength of 330 nm. The band emission is due to the $4f^65d^1 \rightarrow 4f^7$ transition of Eu²⁺ ion. As discussed in our previous work [9,11], BaMg₂Al₆Si₉O₃₀ crystallizes into a hexagonal structure with space group P6/mcc and has two types of independent cation sites, 12-fold coordinated Ba²⁺ site and 6-fold coordinated Mg²⁺ site. The PL spectra of Eu²⁺ in BMAS exhibit two PL bands, indicating the two cation sites. The average interatomic length between Eu²⁺ and oxygen ($d_{\text{Eu-O}}$) is 3.111 Å for Eu²⁺ occupying the Ba²⁺ site and 2.138 Å for occupying the Mg²⁺ site. Therefore we speculate that the band at 376 nm is assigned to Eu²⁺(I) occupying Ba²⁺ with weak crystal field, and the other one at 450 nm corresponds to Eu²⁺(II) occupying Mg²⁺ with strong crystal field. By changing the x value from 0.005 to 0.1, the PL intensity ratio of blue (I_B) to ultraviolet (I_U) emission increases. As shown in Fig. 2(b), this phenomenon can be explained by the energy transfer from Eu²⁺(I) to Eu²⁺(II). As we know, the distance between Eu²⁺ ions is reduced with increasing Eu²⁺ contents, and then the energy transfer from Eu²⁺(I) to Eu²⁺(II) happens. Fig. 3 presents the emission and excitation spectra of BMAS:Eu²⁺. The shape of the PLE spectrum monitored at 376 nm (a) is remarkably different from that monitored at 450 nm (b). It further confirms that the two emission bands are originated from the Eu²⁺ ions occupying two different cation sites. Obviously, there is a significant spectral overlap between the Eu²⁺(I) PL and Eu²⁺(II) PLE spectra in the range 350–420 nm, which implies the possibility of energy transfer from Eu²⁺(I) to Eu²⁺(II). The energy transfer between two different Eu²⁺ ions would result in the enhancement of the blue (I_B) relative to ultraviolet (I_U) band in these phases. In addition, the Eu²⁺(II) emission bands shift to long wavelength by about 26 nm with increasing Eu²⁺ concentrations, as presented in Fig. 2. The redshift can be ascribed to the energy transfer among Eu²⁺(II) ions. In fact, as Eu²⁺ concentrations increase, the distance among Eu²⁺(II) ions becomes shorter, which consequently increases the probability of energy transfer among

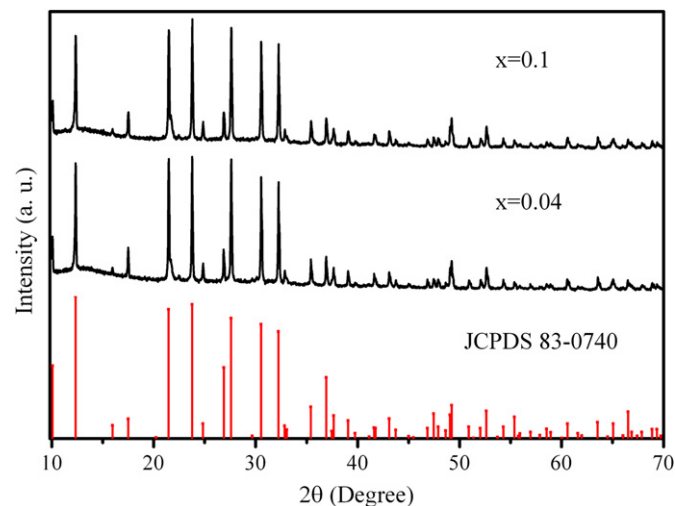


Fig. 1. XRD patterns of BMAS:Eu²⁺ with various Eu²⁺ concentrations.

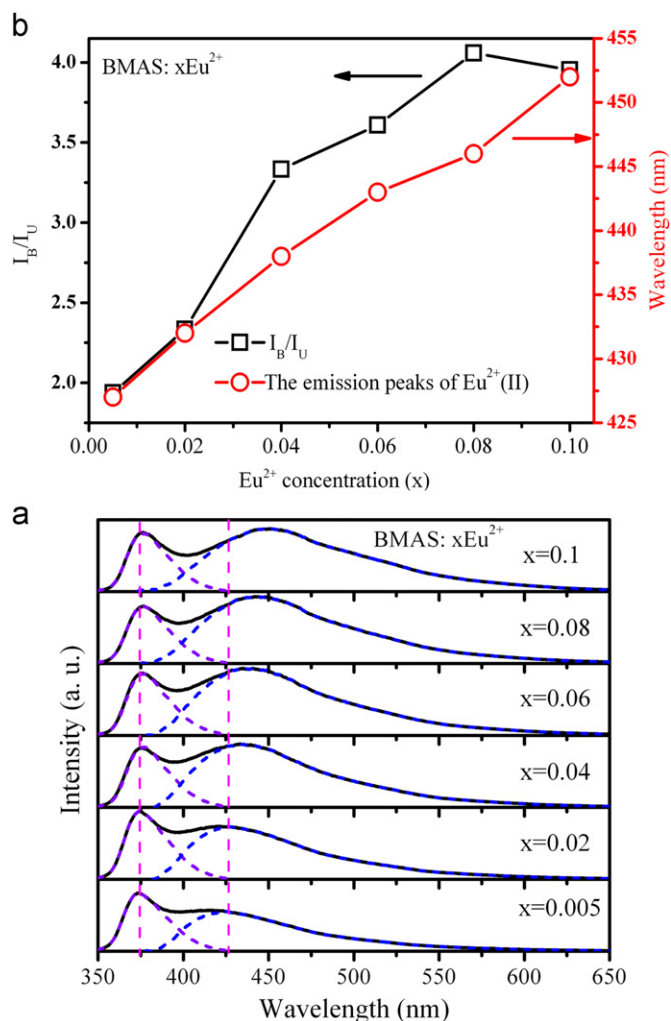


Fig. 2. (a) PL spectra of the Eu²⁺-activated BMAS as a function of Eu²⁺ concentration ranging from 0.005 to 0.1 mol ($\lambda_{\text{ex}}=330$ nm) and (b) concentration dependence of the PL intensity ratio of blue (I_B) to ultraviolet (I_U) emission and the emission peaks of Eu²⁺(II).

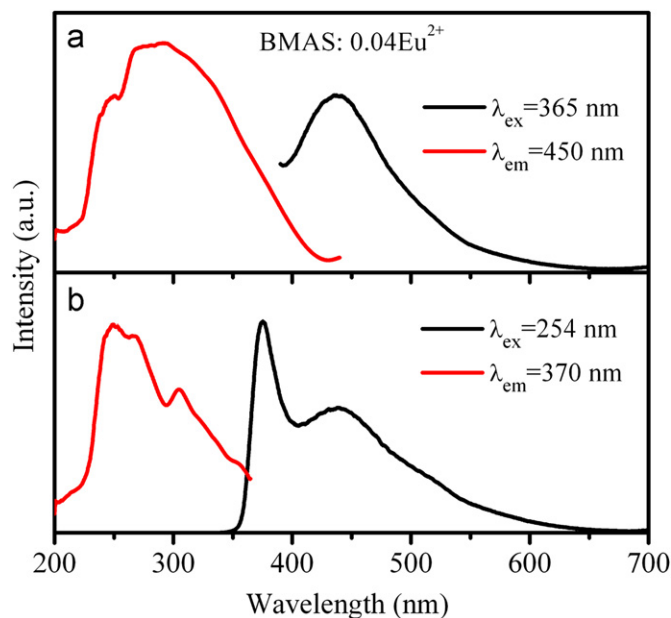


Fig. 3. Emission and excitation spectra of BMAS:Eu²⁺ sample.

$\text{Eu}^{2+}(\text{II})$ ions. This causes the emission energy of Eu^{2+} ions from the higher 5d levels to the 4f ground state to become lower. Hence, the increase in the probability for nonradiative energy transfer in 5d level causes a shift in emission peak toward longer wavelengths. This behavior was also observed in the case of $\text{CaSi}_2\text{O}_2\text{N}_2:\text{Eu}^{2+}$ and $\text{Ca}_3\text{Si}_2\text{O}_4\text{N}_2:\text{Eu}^{2+}$ [12,13], in which the emission energy from the 5d excited state is transferred to the lower levels, therefore resulting in the redshift of emission spectra. BMAS: Eu^{2+} exhibits a broad and strong PLE band peaking at around 254 nm, indicating that BMAS: Eu^{2+} can be used as a promising phosphor applied for fluorescent lamps.

In order to further understand the process of energy transfer, the fluorescence lifetimes τ_1 for $\text{Eu}^{2+}(\text{I})$ and τ_2 for $\text{Eu}^{2+}(\text{II})$ with different Eu^{2+} concentrations are measured and presented in Fig. 4. The fluorescence of $\text{Eu}^{2+}(\text{I})$ decays faster and tends to be a nonexponential function with increasing Eu^{2+} concentrations, reflecting the characteristics of energy transfer from donor $\text{Eu}^{2+}(\text{I})$ to acceptor $\text{Eu}^{2+}(\text{II})$ [14]. It can be seen that the lifetime of $\text{Eu}^{2+}(\text{II})$ is almost the same. Hence, we show the nonexistence of $\text{Eu}^{2+}(\text{II})$ concentration quenching within the range of Eu^{2+} concentrations of interest in this work.

Based on energy transfer, the dependences of the emission intensity ratios of $\text{Eu}^{2+}(\text{II})$ to $\text{Eu}^{2+}(\text{I})$ on Eu^{2+} concentration can be calculated by the measured fluorescence lifetimes of $\text{Eu}^{2+}(\tau_1)$ and $\text{Eu}^{2+}(\tau_2)$. Under steady excitation, the absorbing population of $\text{Eu}^{2+}(\text{I})$ is equal to the radiative population of the level. As for $\text{Eu}^{2+}(\text{II})$, the radiative population of the level is equal to the sum of the absorbing population and the population transfer from $\text{Eu}^{2+}(\text{I})$. So the rate equations describing the energy transfer from $\text{Eu}^{2+}(\text{I})$ to $\text{Eu}^{2+}(\text{II})$ can be written as follows:

$$f_1 = n_1/\tau_1 \quad (1)$$

$$f_2 + Wn_1 = n_2/\tau_2 \quad (2)$$

where f_1 and f_2 are the absorbing populations of $\text{Eu}^{2+}(\text{I})$ and $\text{Eu}^{2+}(\text{II})$, and n_1 and n_2 are the populations of $\text{Eu}^{2+}(\text{I})$ and $\text{Eu}^{2+}(\text{II})$ ions, respectively. The energy transfer rates W can be obtained by $W = 1/\tau_1 - 1/\tau_0$, since the fluorescence intensity of a specific level is proportional to the product of the population and the radiative rate of the level, using Eqs. (1) and (2), the dependences of the intensity ratio of the blue–ultraviolet emissions on Eu^{2+} concentrations can be written as

$$\frac{I_B}{I_U} \propto \frac{n_2\gamma_2}{n_1\gamma_1} = \frac{\gamma_2}{\gamma_1} \left(\frac{\tau_2 f_2}{\tau_1 f_1} + \tau_2 W \right) = \frac{\gamma_2}{\gamma_1} \left(\frac{\tau_2}{\tau_1} \left(\frac{f_2}{f_1} + 1 \right) - \frac{\tau_2}{\tau_0} \right) \quad (3)$$

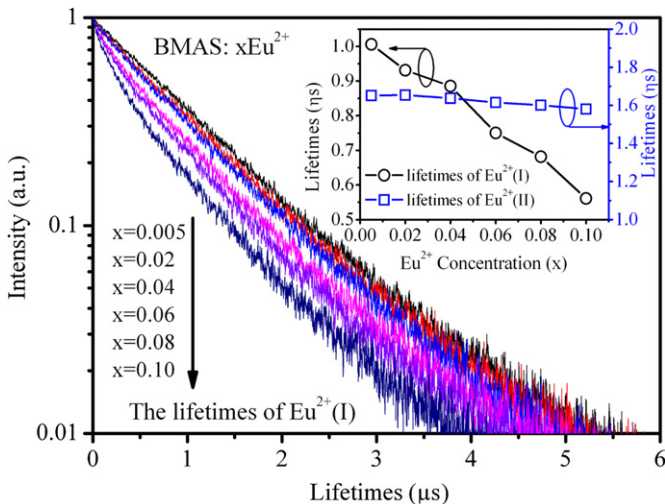


Fig. 4. Fluorescence decay curves of $\text{Eu}^{2+}(\text{I})$ with different Eu^{2+} concentrations. Inset: the calculated average lifetimes of $\text{Eu}^{2+}(\text{I})$ and $\text{Eu}^{2+}(\text{II})$.

where γ_1 and γ_2 are radiative rates of $\text{Eu}^{2+}(\text{I})$ and $\text{Eu}^{2+}(\text{II})$, respectively, which are independent of Eu^{2+} concentrations. τ_2 has been presented in the Fig. 4 inset. In Eq. (3), the γ_2/γ_1 and f_2/f_1 can be considered to be constants. Fig. 5 shows the intensity ratios I_B/I_U vs $1/\tau_1$ at various Eu^{2+} concentrations. It can be seen that the I_B/I_U goes linearly with $1/\tau_1$ when Eu^{2+} content is lower than 0.1, demonstrating that the energy transfer between $\text{Eu}^{2+}(\text{I})$ and $\text{Eu}^{2+}(\text{II})$ plays the main role on the I_B/I_U . The dots at $x=0.04$ and 0.1 that deviate from straightness might be attributed to large variation of f_2/f_1 at these positions. In other words, when Eu^{2+} concentration is 0.04 mol, the distribution of Eu^{2+} at different sites may change significantly, leading to the change of f_2/f_1 ; when Eu^{2+} concentration is 0.1 mol, the saturated absorption population of $\text{Eu}^{2+}(\text{II})$ to $\text{Eu}^{2+}(\text{I})$ may vary remarkably at high Eu^{2+} concentrations, resulting in the variation of f_2/f_1 .

Fig. 6 shows PL spectra of white light emitting BMAS:0.08 Eu^{2+} , 0.08 Tb^{3+} , and 0.16 Mn^{2+} phosphor upon 254 nm excitation. The present phosphor shows four emission bands, ultraviolet band of 376 nm, blue band of 450 nm, green band of 542 nm and red band of 610 nm. Although a significant part of the energy at 376 nm

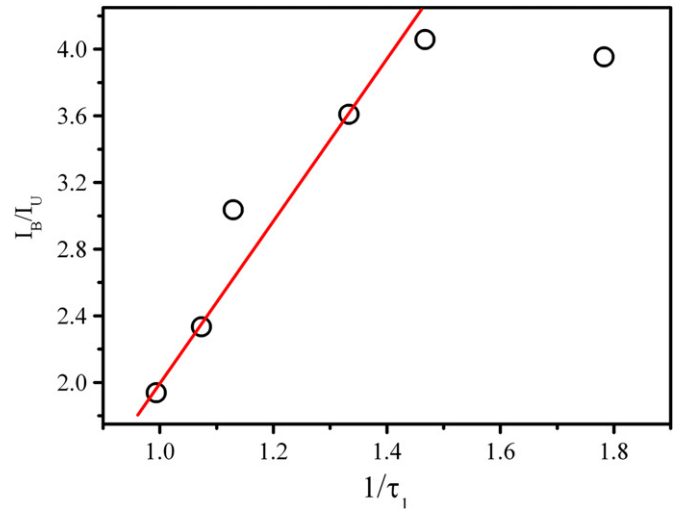


Fig. 5. Intensity ratios I_B/I_U vs $1/\tau_1$ at various Eu^{2+} concentrations.

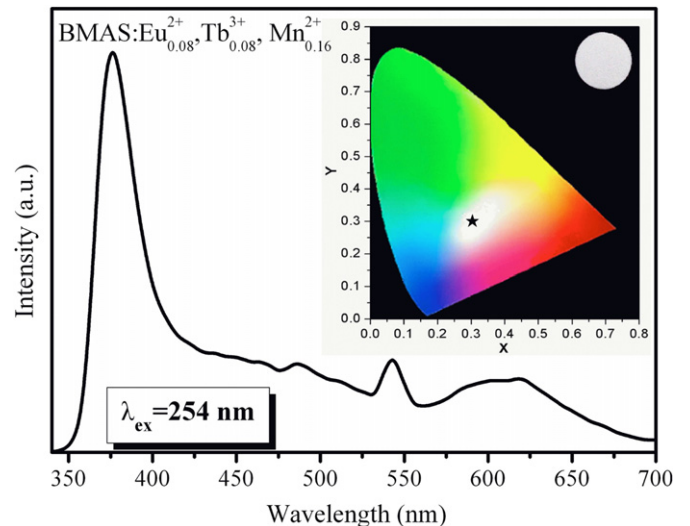


Fig. 6. PL spectra of white light emitting BMAS:0.08 Eu^{2+} , 0.08 Tb^{3+} , and 0.16 Mn^{2+} phosphor upon 254 nm excitation. Inset: the corresponding CIE color coordinates and photographs of this phosphor.

would lead to low quantum efficiency, the full-color emitting BMAS:Eu²⁺, Tb³⁺, and Mn²⁺ phosphors may be applied as potential phosphors for FLs. The emitting color point of a white light sample is illustrated in the CIE chromaticity diagram, as shown in the Fig. 6 inset. The corresponding CIE color coordinates of the phosphor are (0.29, 0.29). By using the optimized Eu²⁺ content at $x=0.08$, white light emitting BMAS:0.08 Eu²⁺, 0.08Tb³⁺, and 0.16Mn²⁺ phosphor can be generated upon 254 nm excitation. Our results indicate that BMAS:Eu²⁺, Tb³⁺, and Mn²⁺ are great potential candidates for using single phosphor converted white light source for recent FLs.

4. Conclusion

In conclusion, BaMg₂Al₆Si₉O₃₀:Eu²⁺ phosphors are synthesized by the solid-state reaction method and the energy transfer among Eu²⁺ ions is investigated. The energy transfer leads to the following results: (1) with increasing Eu²⁺ doping content, the blue emission bands shift gradually to longer wavelength and the PL intensity ratio of blue (I_B) to ultraviolet (I_U) emission increases; (2) the ratio of blue emission to ultraviolet emission by the experiment is analyzed by theoretical calculation based on energy transfer and lifetime measurements. By using the optimized Eu²⁺ content, the white light emitting BMAS:0.08Eu²⁺, 0.08Tb³⁺, and 0.16Mn²⁺ phosphor can be realized upon 254 nm excitation. The present letter indicates that our samples can be promising candidate phosphors for FLs.

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